

Spacetime Interfacial Damage Model for Elastodynamic Fracture with Riemann Contact Conditions

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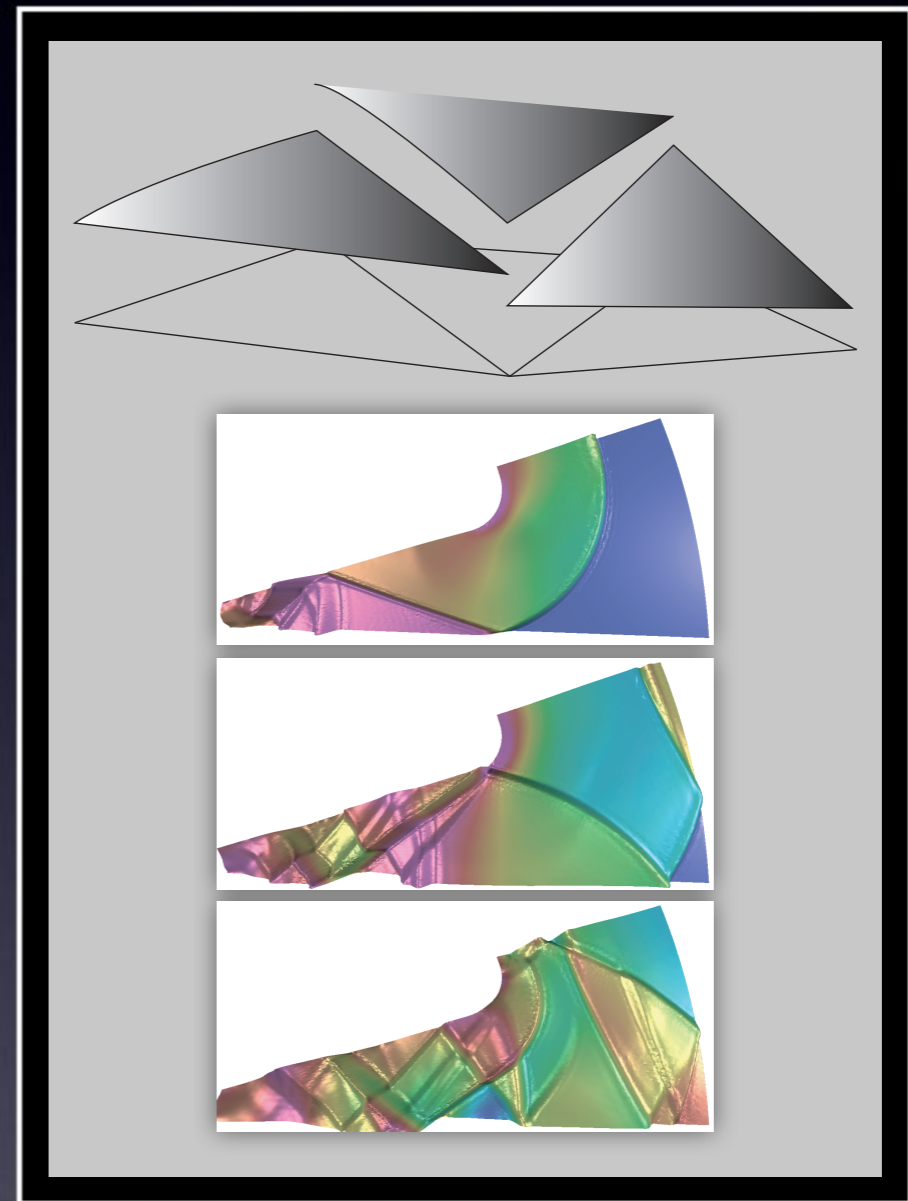
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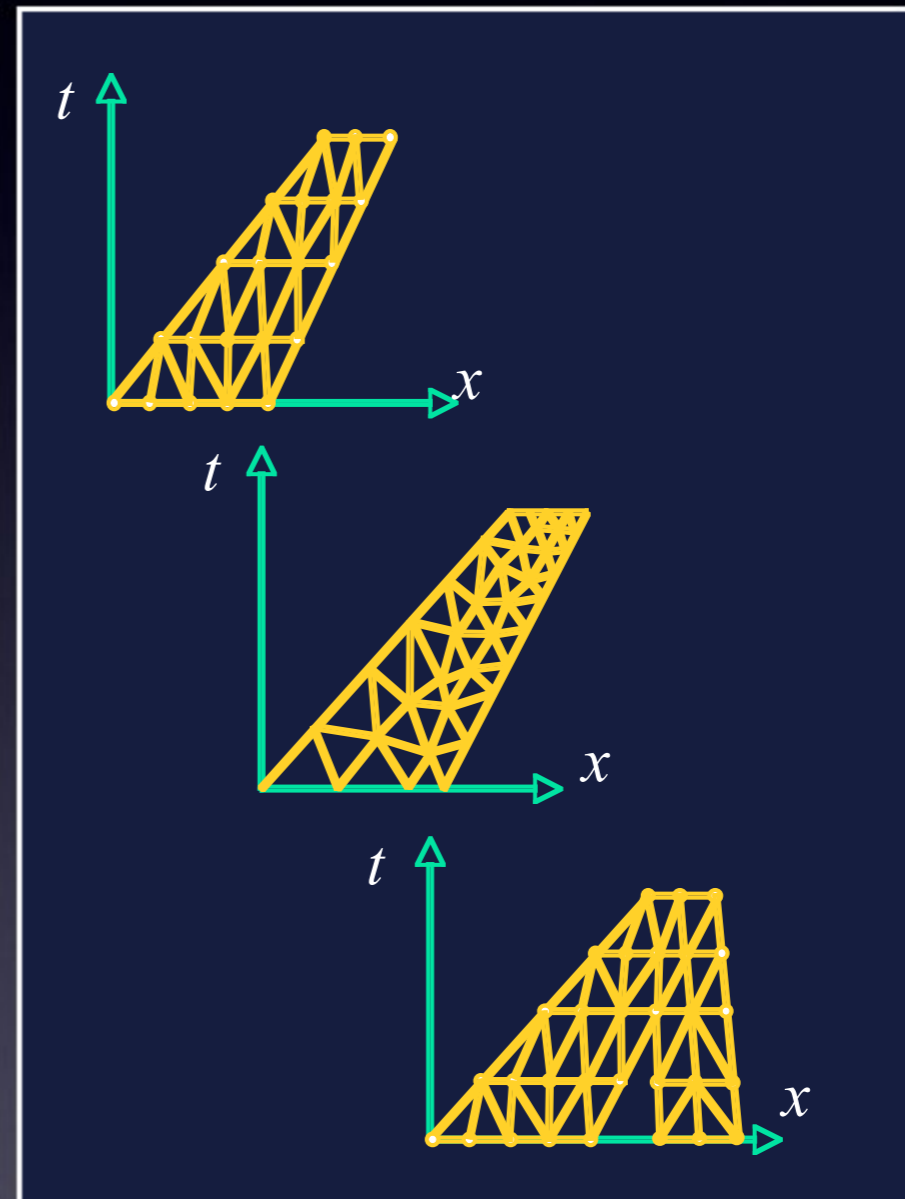
Features of spacetime discontinuous Galerkin finite element methods

- Inter-element discontinuous basis functions
- Weak enforcement of balance/conservation jump conditions in spacetime (e.g., Rankine–Hugoniot conditions for conservation laws)
- Enables exact conservation per element and $O(N)$ complexity for hyperbolic problems

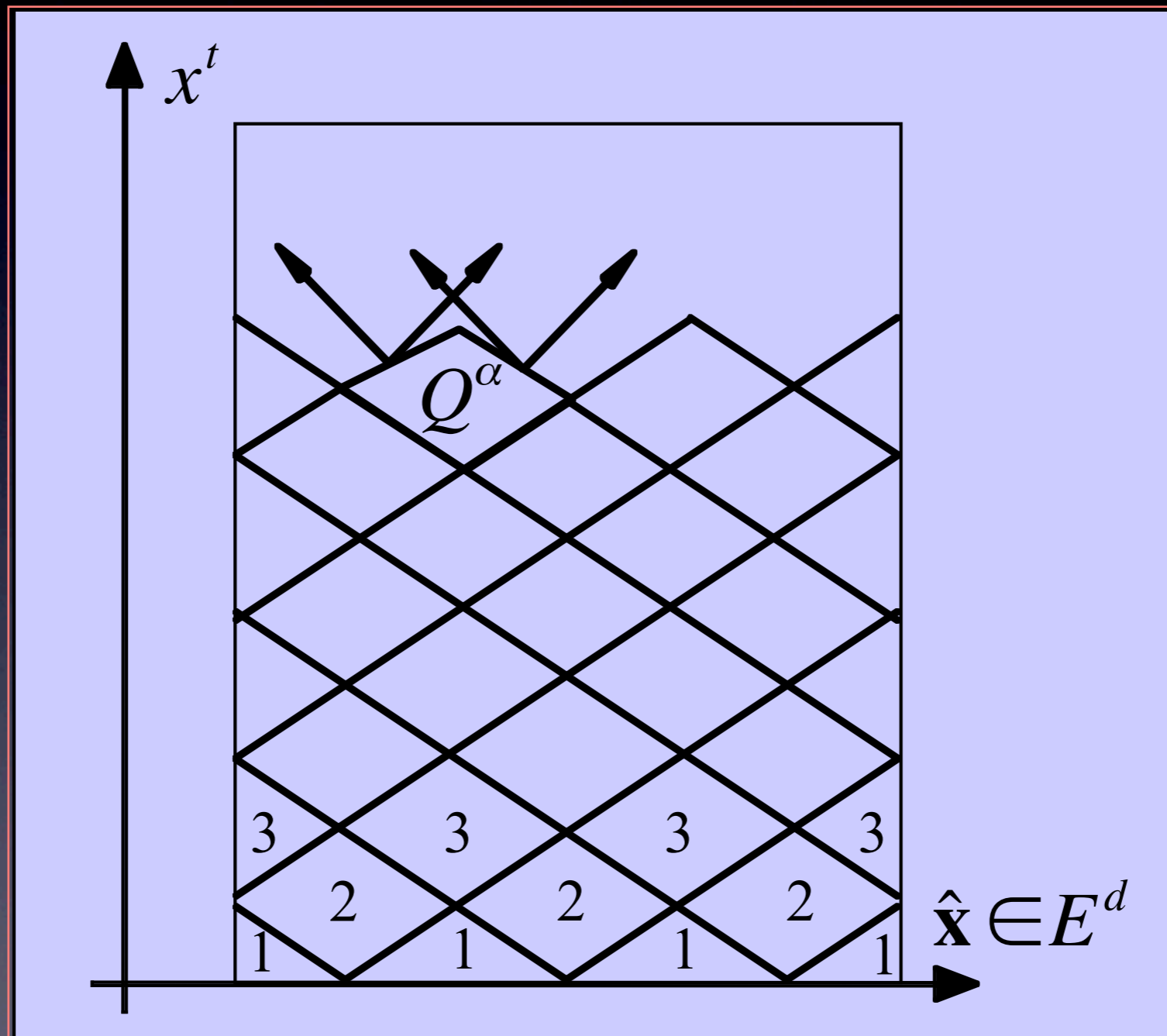


Features of spacetime discontinuous Galerkin finite element methods

- Direct discretization of spacetime
 - Replaces separate temporal integration
 - Unstructured spacetime mesh eliminates tangling in moving-boundary problems
 - Unambiguous numerical framework for initial/boundary conditions (vs. finite volume, finite difference)

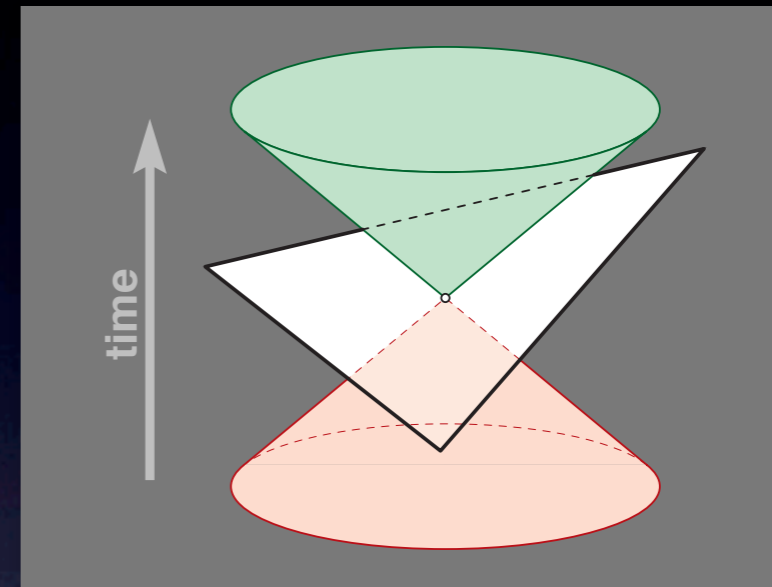


Causal Spacetime Mesh and $O(N)$ Advancing-Front Solution Strategy

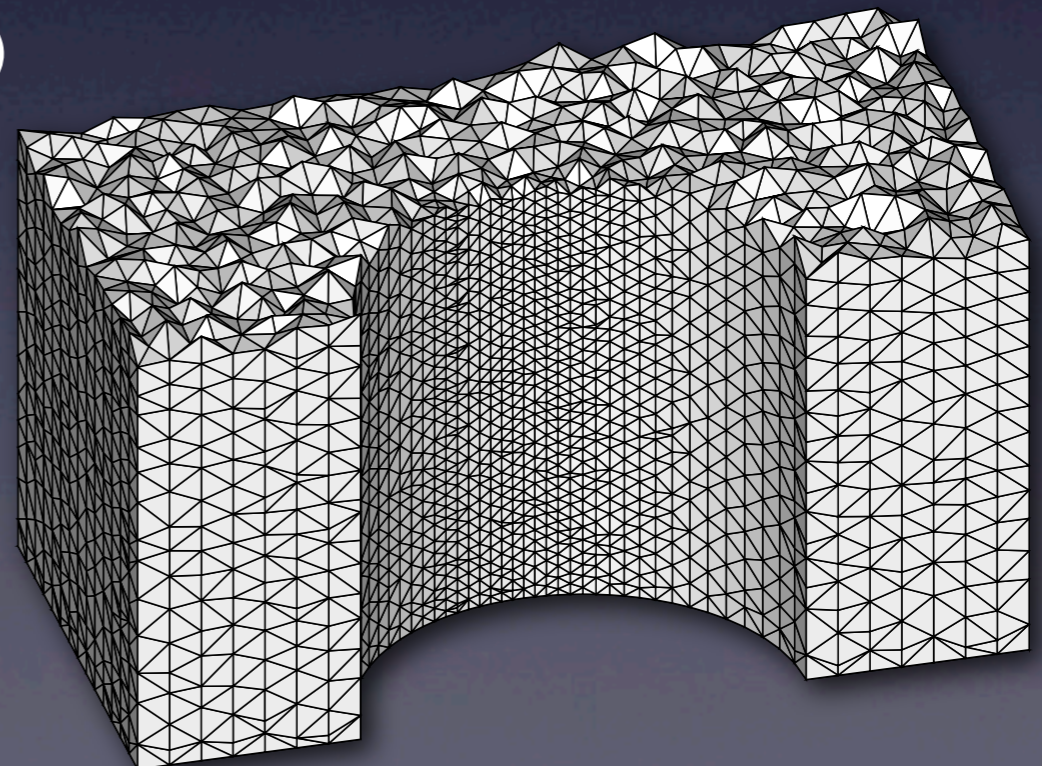
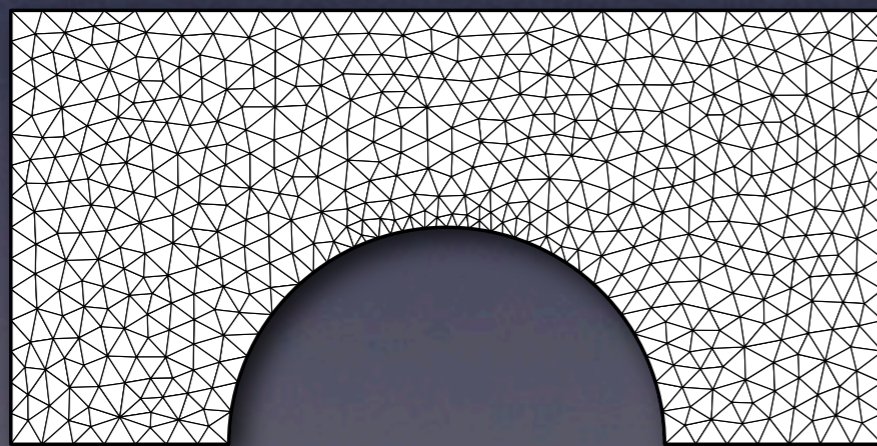


Tent Pitcher: causal spacetime meshing

- Given a space mesh, Tent Pitcher constructs a spacetime mesh such that every facet on sequence of advancing fronts is spacelike (patch height bounded by *causality constraint*)
- Similar to CFL condition, except entirely *local* and not related to stability (required for $O(N)$ solution)

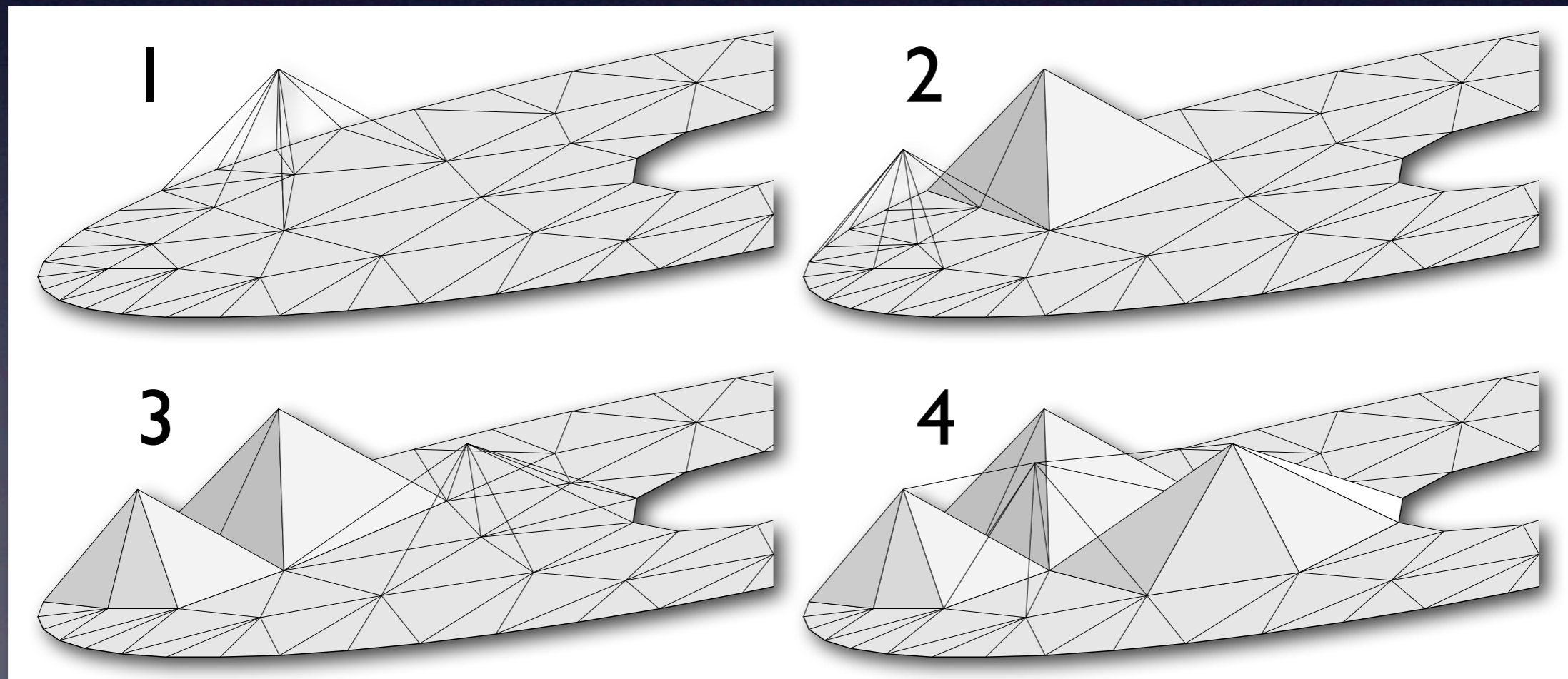


causality constraint



Tent Pitcher: patch-by-patch meshing & solution

- Patches ('tents') of tetrahedra; solve immediately for $O(N)$ method with rich parallel structure
- Maintain "space mesh" as advancing, space-like front with non-uniform time coordinates



Force-like fields

($d, d+1$ -forms with vector coefficients)

- Body force ($(d+1)$ -form): $\mathbf{b} = \mathbf{b}\Omega$
- Spacetime momentum flux (d -form): $\mathbf{M} = \mathbf{p} - \boldsymbol{\sigma}$

$\mathbf{M}|_{\Gamma} =$ linear momentum flux across d -manifold Γ

- Linear momentum density: $\mathbf{p} = \mathbf{p}\star dt$
- Stress/traction: $\boldsymbol{\sigma} = \boldsymbol{\sigma} \wedge \star d\mathbf{x}$
- Stokes Theorem: $\int_{\partial Q} \mathbf{M} = \int_Q d\mathbf{M}$

Momentum Balance

- Integral form of linear momentum balance:

$$\int_{\partial Q} \mathbf{M} = \int_Q \rho \mathbf{b} \quad \forall Q \subset D$$

$$\int_Q (d\mathbf{M} - \rho \mathbf{b}) = \mathbf{0} \quad \forall Q \subset D \quad (\text{Stokes Thm.})$$

- Local form with jump part:

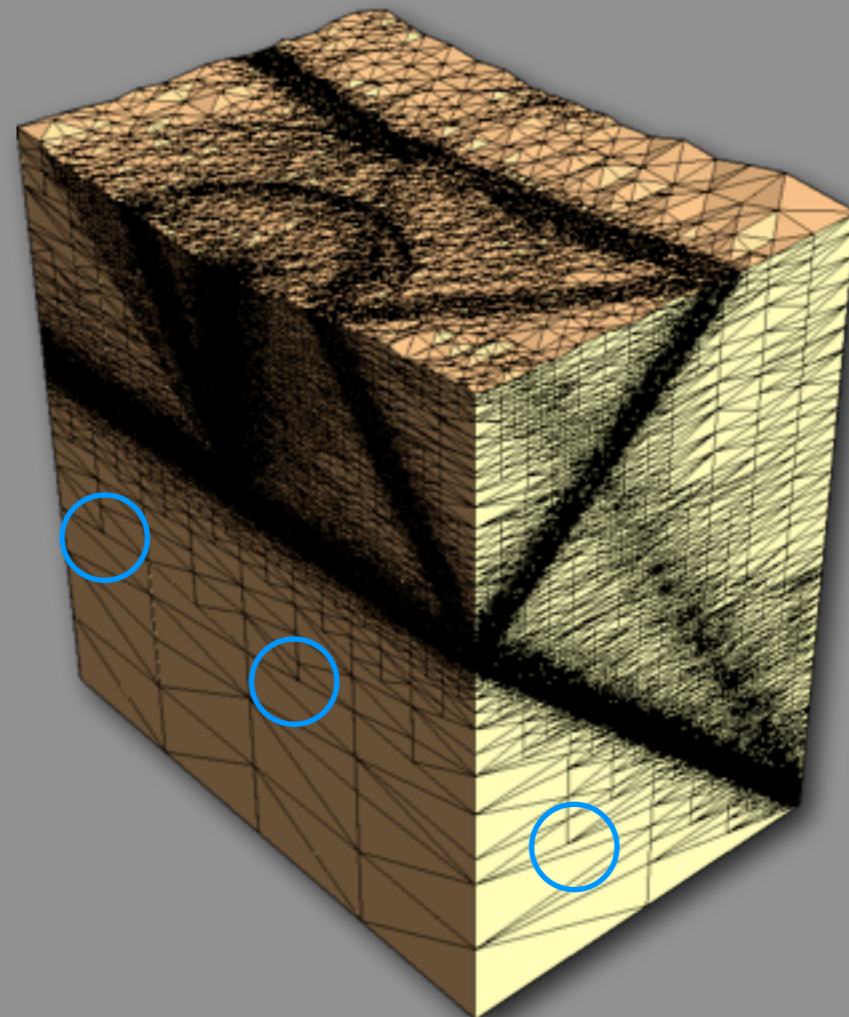
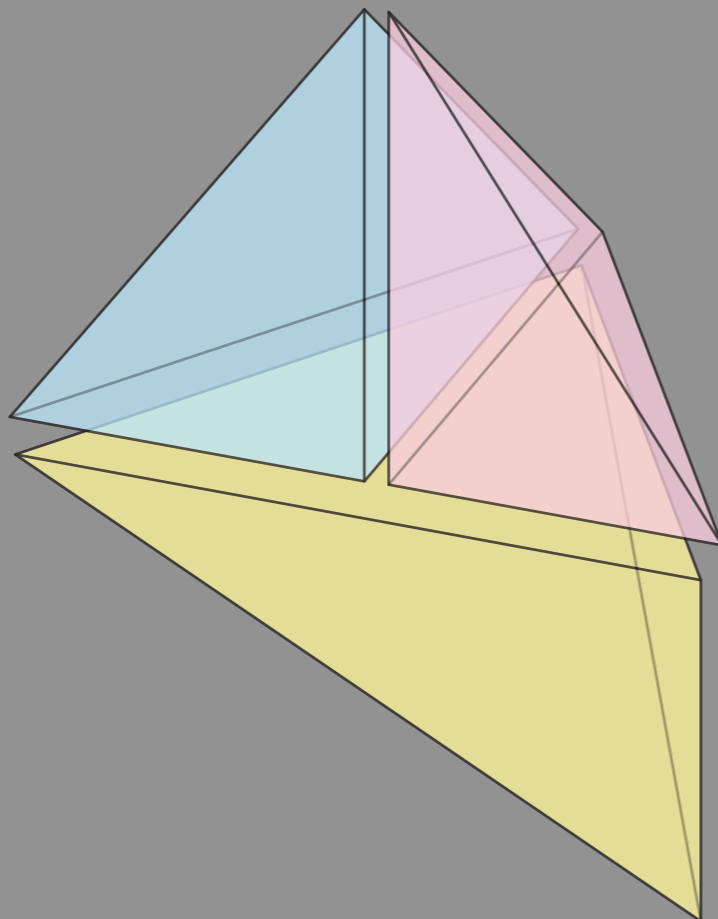
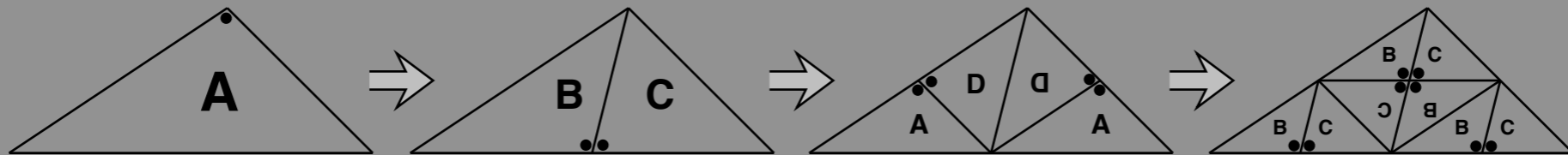
$$(d\mathbf{M} - \rho \mathbf{b})|_{D \setminus \Gamma^J} = \mathbf{0}$$

$$[[\mathbf{M}]]|_{D \cap \Gamma^J} = \mathbf{0} \rightarrow (\mathbf{M}^* - \mathbf{M})|_{Q \cap \Gamma^J} = \mathbf{0}$$

$$\mathbf{M}^* = \text{Riemann or prescribed value}$$

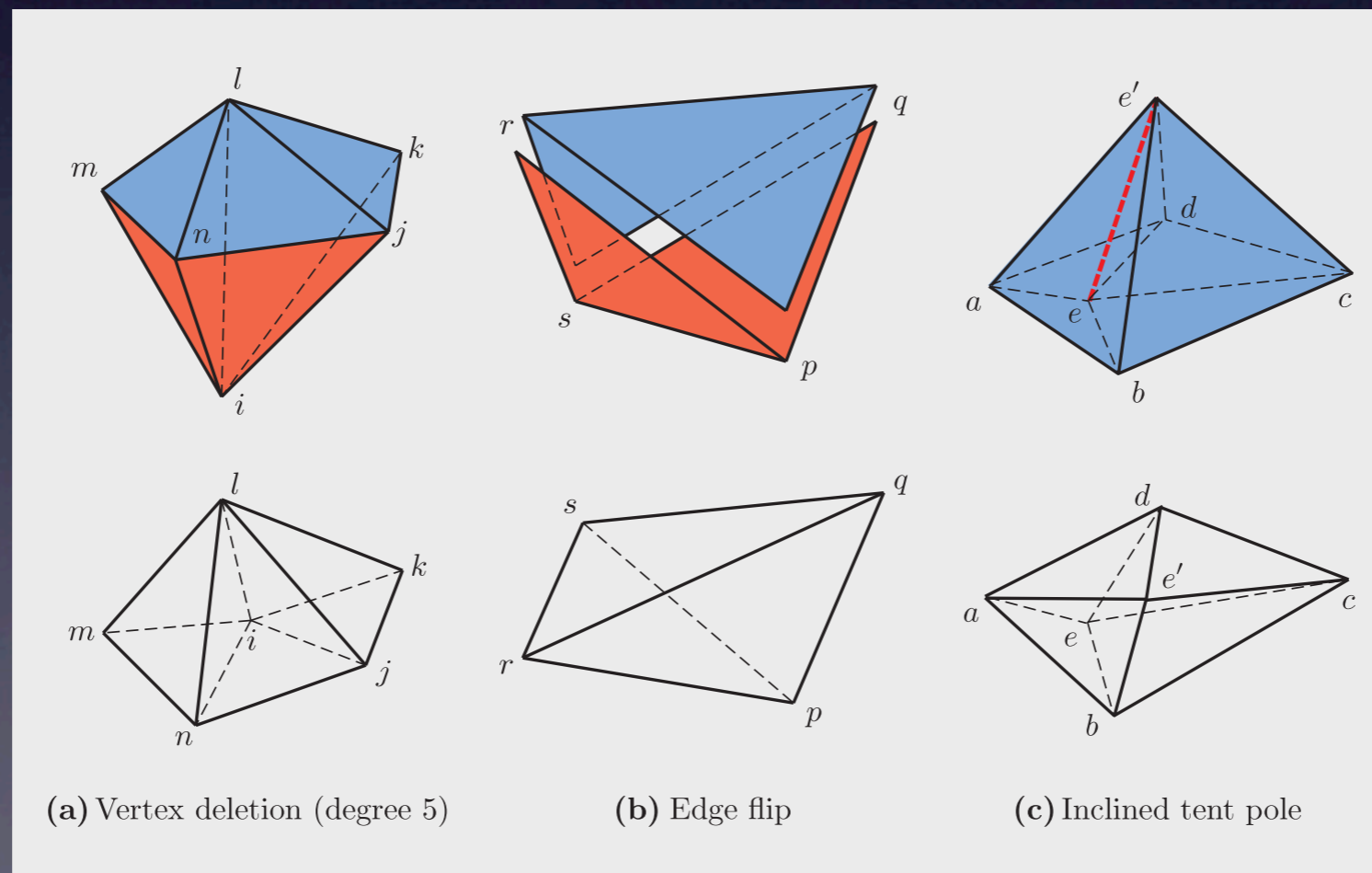
Adaptive refinement

● Refine space mesh using *newest vertex method* to maintain element quality



More adaptive meshing

- New spacetime adaptive meshing operations:
 - Edge flip; vertex deletion; inclined tent poles (ALE) for smoothing, tracking and repositioning
 - Spacetime format eliminates projection error



Energy balance and dissipation error indicator for adaptive meshing

- On every spacetime element Q , SDG solution satisfies:

$$\begin{aligned}\varphi^Q &= \int_Q \dot{\mathbf{u}} \wedge \rho \mathbf{b} + \frac{1}{2} \int_{\partial Q} (\dot{\mathbf{u}}^* \wedge \mathbf{M}^* + \boldsymbol{\varepsilon}^* \wedge \mathbf{iM}^*) \\ &= \frac{1}{2} \int_{\partial Q} \{(\dot{\mathbf{u}}^* - \dot{\mathbf{u}}) \wedge (\mathbf{M}^* - \mathbf{M}) + (\boldsymbol{\varepsilon}^* - \boldsymbol{\varepsilon}) \wedge \mathbf{i}(\mathbf{M}^* - \mathbf{M})\}\end{aligned}$$

Right hand side is non-negative numerical dissipation for Q .

- To control adaptively the numerical dissipation, we use the element-wise error indicator:

$$\varphi^Q \approx \text{tol}_\varphi$$

Cohesive error measure and adaptive error indicator

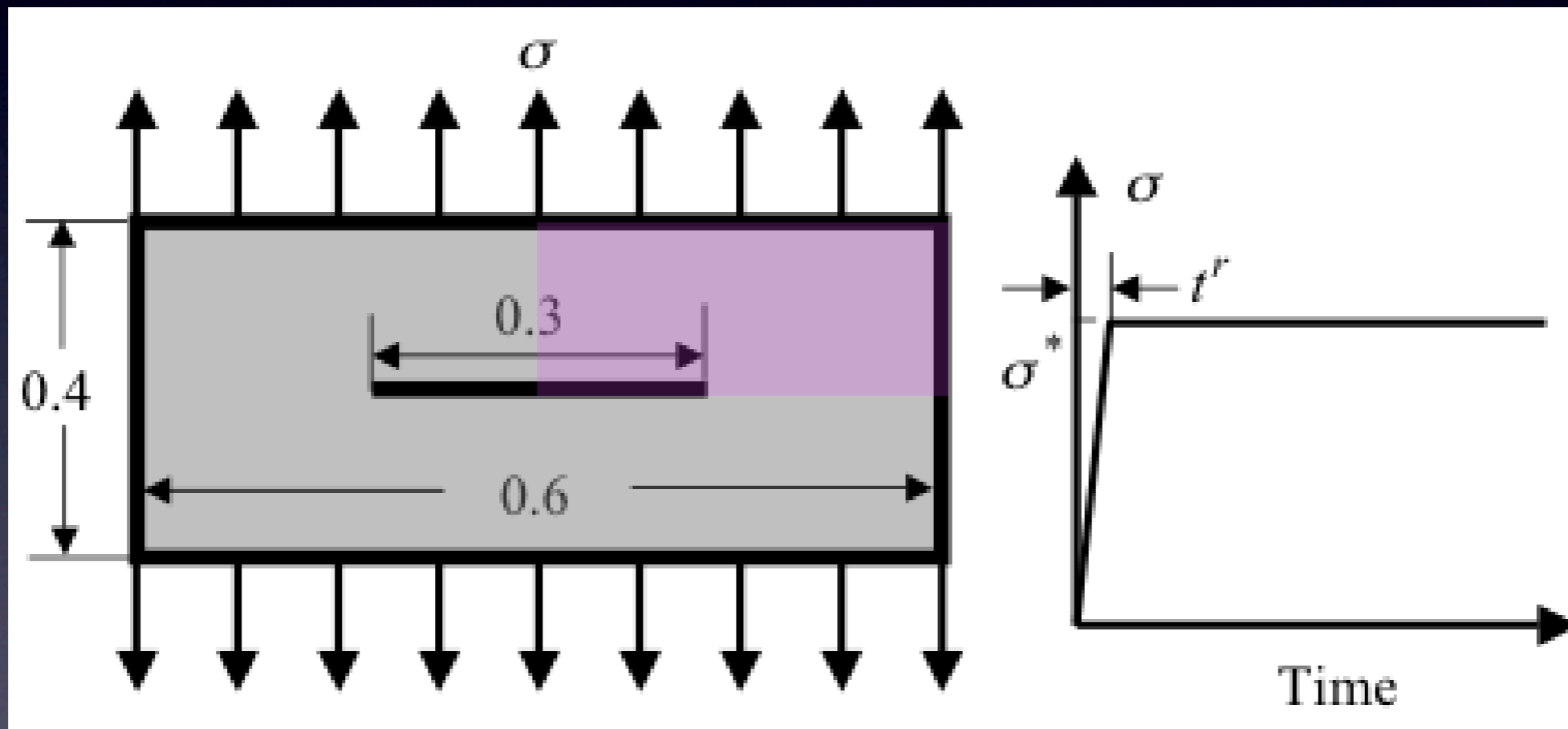
- Cohesive energy error on cohesive surface trajectory, Γ^C :

$$\epsilon_C = \|(\mathbf{t} - \mathbf{t}_{\text{coh}}) \cdot \mathbf{v}\|_{L_1(\Gamma^C)}$$

- Cohesive error indicator for element Q :

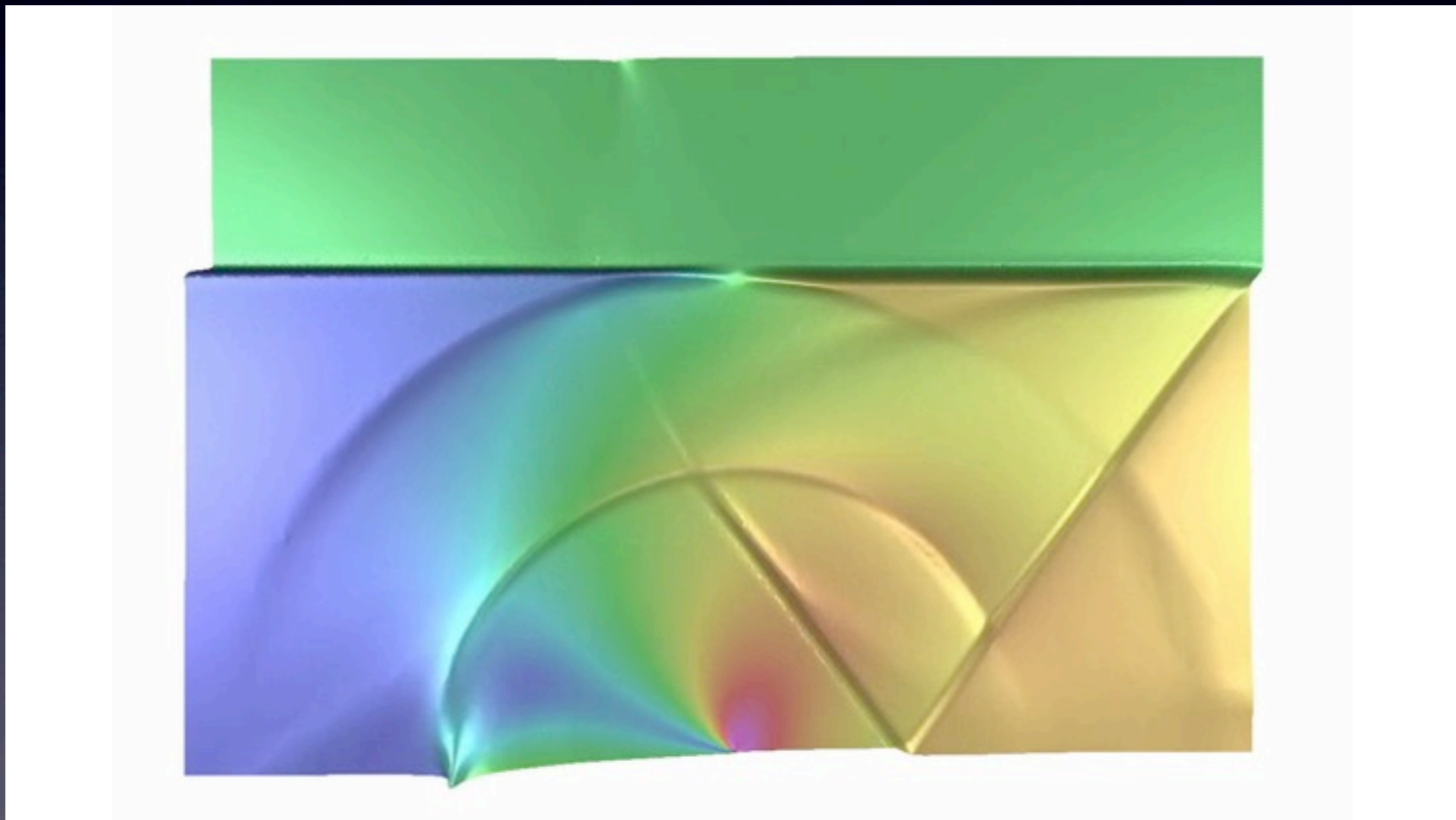
$$\epsilon_C^Q = \|(\mathbf{t} - \mathbf{t}_{\text{coh}}) \cdot \mathbf{v}\|_{L_1(\partial Q_{\text{coh}})} \approx \text{tol}_c$$

Crack-tip Wave Scattering



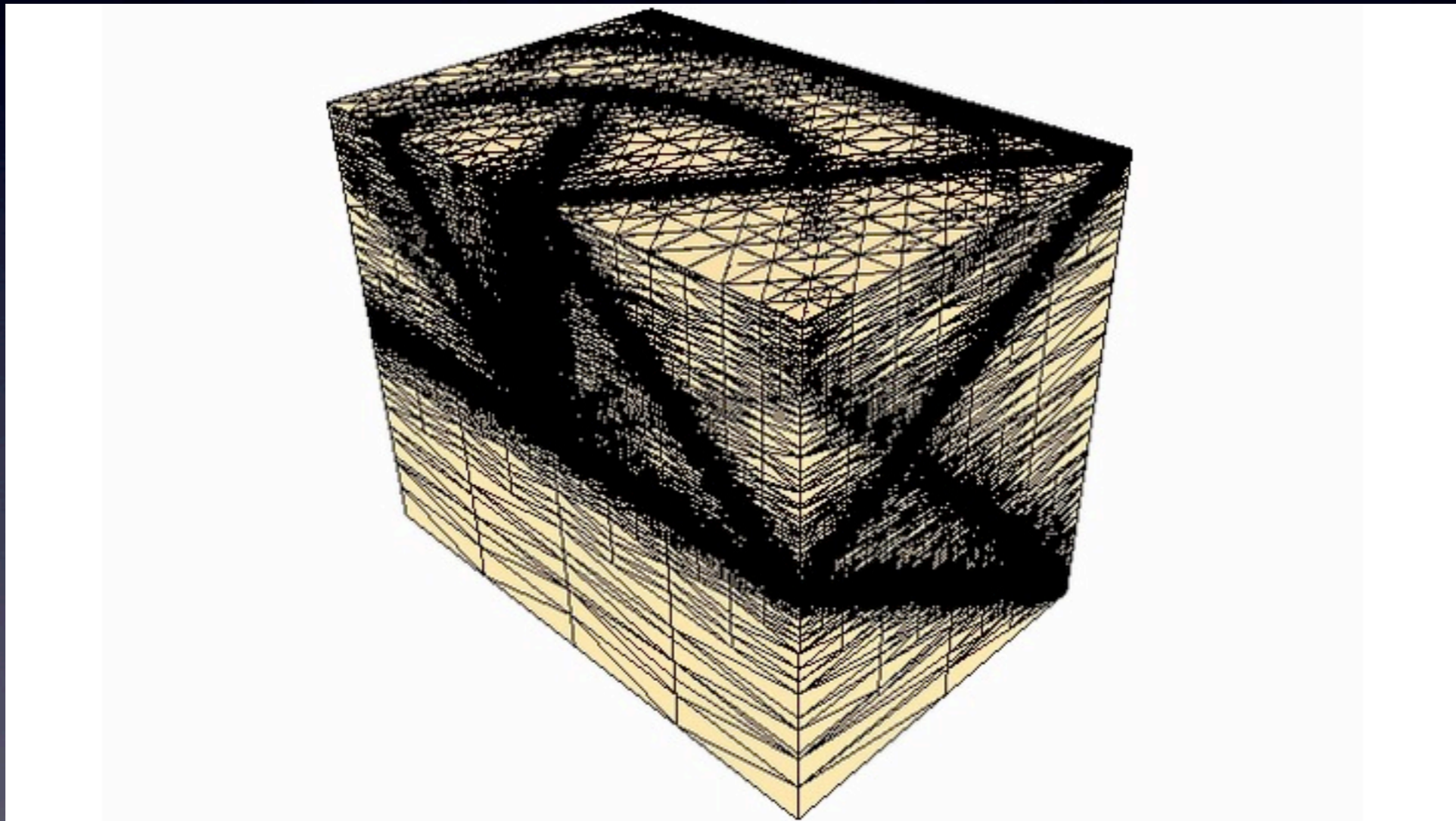
Crack-tip Wave Scattering

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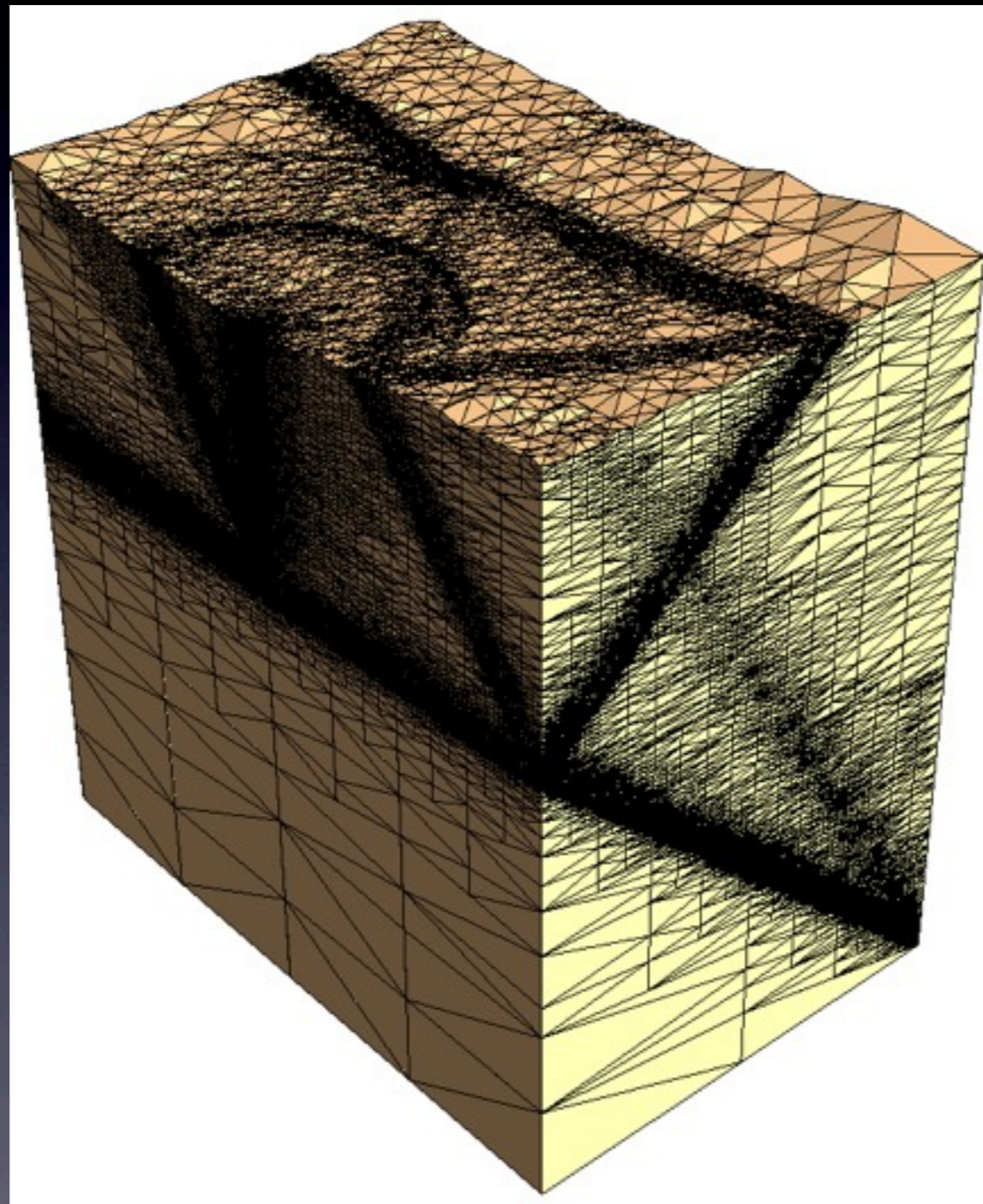


Crack-tip Wave Scattering

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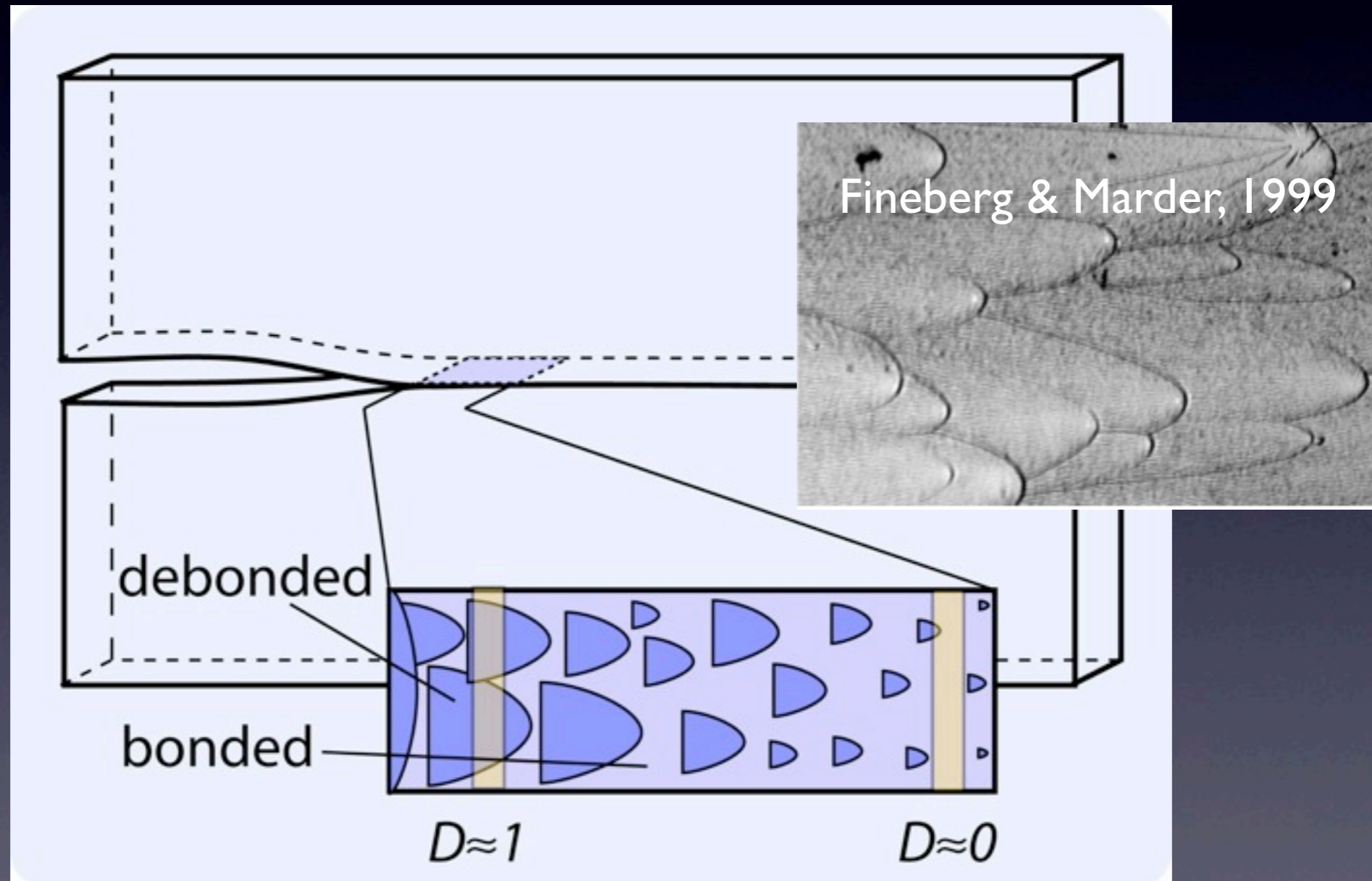


Crack-tip Wave Scattering

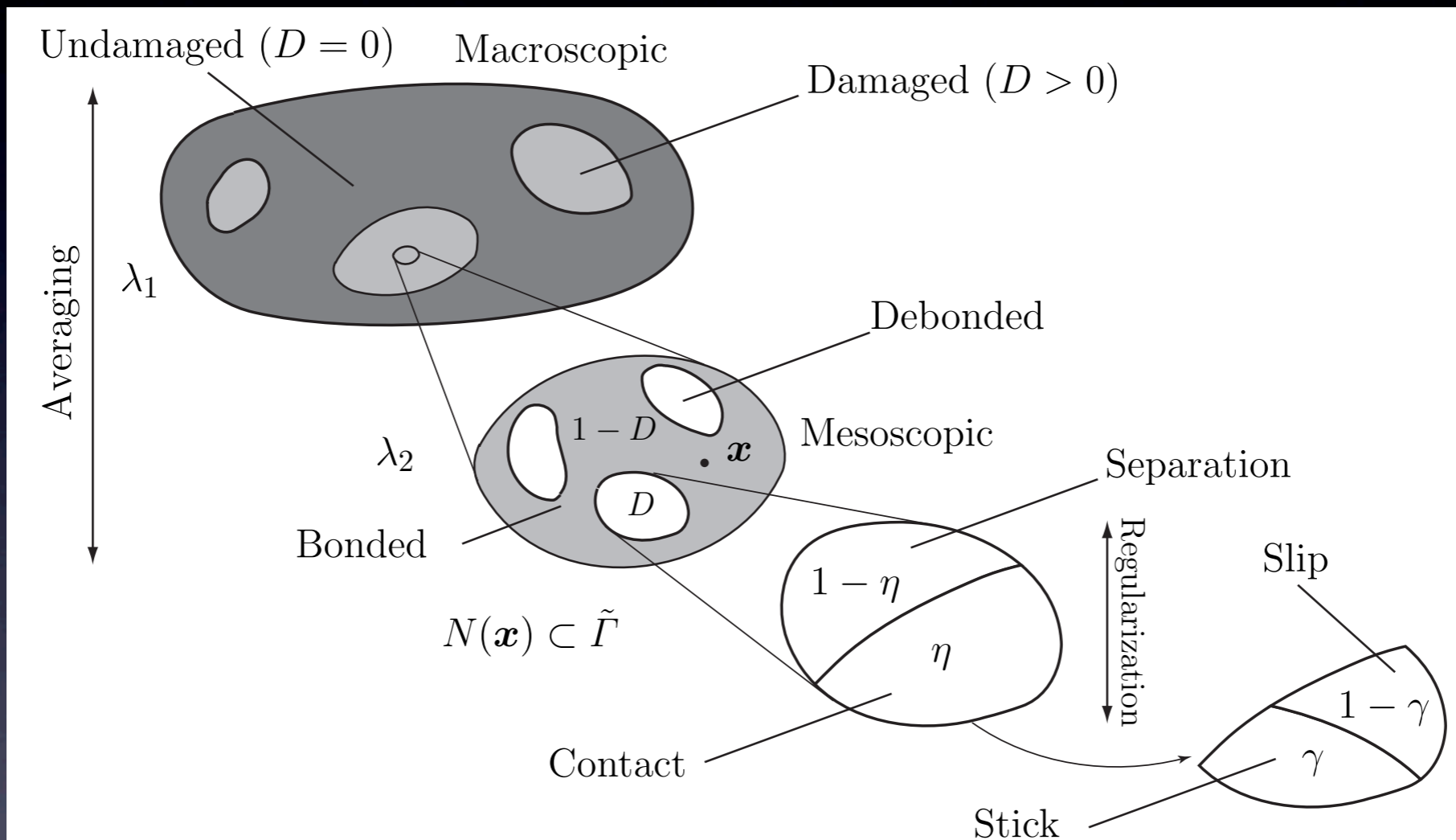


Delay Damage Evolution Cohesive Model

D = Area fraction of debonded surface



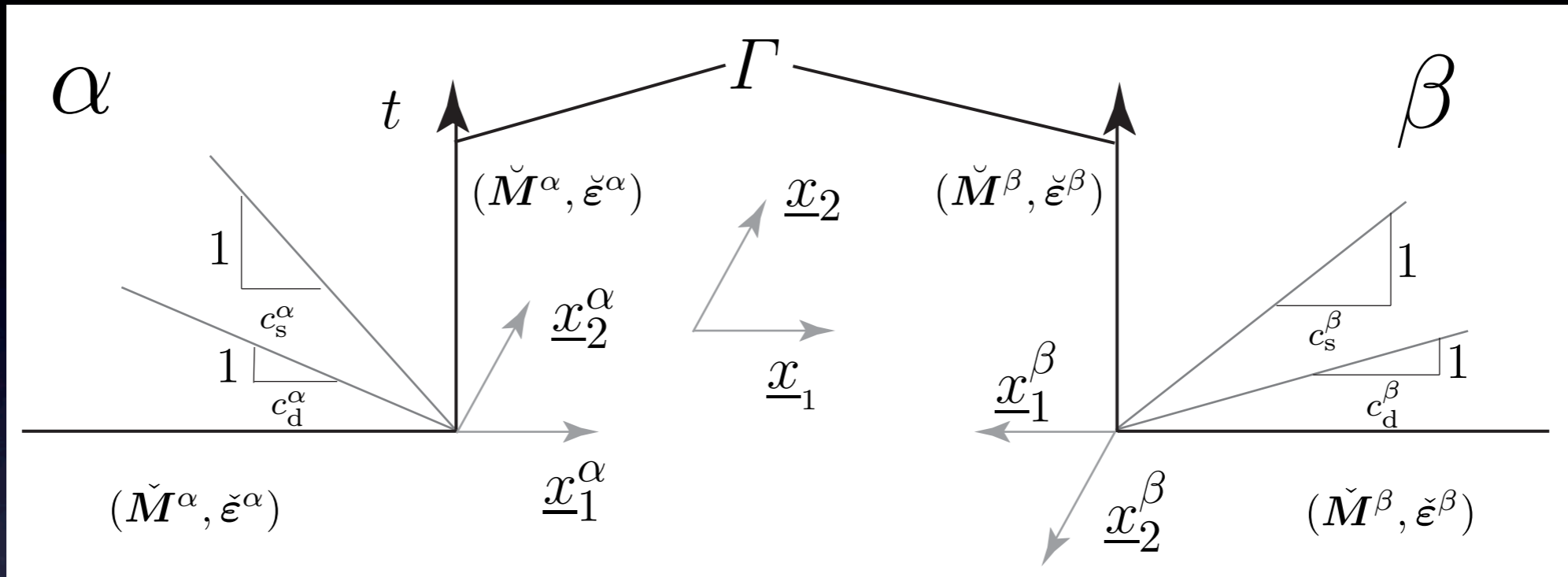
Mesososcopic interface subdivisions



Contact conditions

- Common numerical idealizations of contact
 - ✦ Discrete models - goal is to prevent node penetration
 - ✦ Elastostatic contact conditions
 - ✦ Penalty methods and other inexact formulations
 - $s = f(\Delta \mathbf{u}) = K \Delta \mathbf{u} (\Delta \mathbf{u} < 0)$ but $\Delta \mathbf{u} = 0!$
 - Approach exact solution for large $K \Rightarrow$ divergence
 - ✦ Variational Inequality/Lagrange multiplier methods
 - Lagrange variables increase the problem size
 - typically require an implicit solution scheme

Riemann Contact/Separation Solutions



- Direct solution from momentum balance and compatibility jump conditions
- Matching conditions at material interface
- Preserves characteristic structure of solution
- Weak enforcement of the continuum, dynamic jump conditions converges to correct result

Interface Matching condition

Bonded and contact-stick

- Matching conditions

- Momentum: $[[\check{\mathbf{s}}]] = 0$
- Compatibility: $[[\check{\mathbf{v}}]] = 0$

- Riemann Solution

- $\underline{s}^{jI} = [[\underline{w}^{(j)} / \rho \underline{c}^{(j)}]] / [[(\rho \underline{c}^{(j)})^{-1}]]$
- $\underline{v}_j^I = [[\underline{w}^{(j)}]] / [[(\rho \underline{c}^{(j)})^{-1}]]$

- Contact-stick and bonded solutions are identical

- Contact-stick and separation solutions are physically distinct:

- Contact-stick

- Separation

$$\Delta \mathbf{u} = \mathbf{0}$$

$$\mathbf{s} = ?$$

$$\Delta \mathbf{u} = ?$$

$$\mathbf{s} = \mathbf{0} \quad \text{or} \quad \mathbf{s} = f(\Delta \mathbf{u})(TSL)$$

Interface Matching condition

Contact-slip

- Matching conditions

- Momentum: $[[\check{\mathbf{s}}]] = 0$
- Compatibility: $[[\check{\underline{v}}_1]] = 0$
- Coulomb's law: $\underline{s}^{jII} = k \langle -\underline{s}_I^1 \rangle_+ (e_{\check{\mathbf{v}}})^j \quad j \neq 1$

- Riemann Solution

$$\underline{s}^{jII} = \begin{cases} \underline{s}^{1I} & j = 1 \\ k \langle -\underline{s}_I^1 \rangle_+ (e_{\tau_I})^j & j \neq 1 \end{cases}$$

$$\underline{v}_j^{II} = \begin{cases} \underline{v}_1^I & j = 1, \\ \underline{w}^{(j)} / \rho \underline{c}^{(j)} - \underline{s}^{jII} / \rho \underline{c}^{(j)} = \check{\underline{v}}_j + (\check{\underline{s}}^j - \underline{s}^{jII}) / \rho \underline{c}^{(j)} & j \neq 1 \end{cases}$$

- Slip velocity $[[\check{\mathbf{v}}]] \rightarrow \mathbf{0}$ at stick-slip transition \Rightarrow Discontinuous response from $(e_{\check{\mathbf{v}}})^j$
- many regularizations proposed to solve the problem (Karnopp, Mostaghel, Quinn)
- $\tau_I > 0$ at stick-slip transition \Rightarrow Continuous response from $\tau_I > 0$

Interface Matching condition Separation

- Matching conditions

- Momentum and constitutive equation: $\check{\mathbf{s}} = \mathbf{0}$
- No compatibility condition

- Riemann Solution

- $\mathbf{s}^{III} = \mathbf{0}$

- $\underline{v}_j^{III} = \frac{w^{(j)}}{\rho \underline{c}^{(j)}}$

Mesososcopic subdivision

Damaged area (D)

- General Properties of Damage parameter
 - ◆ Nondecreasing (irreversible process)
 - ◆ Stress Induced (Ravi-Chandar, Yang 1997)
 - ◆ Exact Bonded solution is recovered for $D = 0$
 - ◆ Stress free (crack) surface condition at $D = 1$
- Rate Dependent damage model
 - ◆ Hardening effects at high strain rates (experimental results from Fineberg and Marder 1999)
 - ◆ Introduces a fracture related length scale
 - ◆ Mesh dependency of static damage models (strain softening failure, Bazant and Belytschko 1984)

Damage Evolution

Follow delay damage evolution model of Allix, Feissel and Thévenet:

$$\tilde{\sigma} = \frac{\sqrt{\sigma_N^2 + (\beta\sigma_T)^2}}{\sigma_c} \quad \text{normalized effective stress}$$

$$\bar{D} = f(\tilde{\sigma}) \quad \text{static damage value}$$

$$\dot{D} = \begin{cases} \frac{1}{\tau} [1 - H(\langle f(\tilde{\sigma}) - D \rangle_+)] & \text{if } D < 1 \\ 0 & \text{if } D = 1 \end{cases} \quad \text{damage evolution law}$$

where τ is a *characteristic time scale* for debonding, and $H(0) = 1$ and $\lim_{x \rightarrow \infty} H(x) = 0$. We use

$$H(x) = e^{-ax}$$

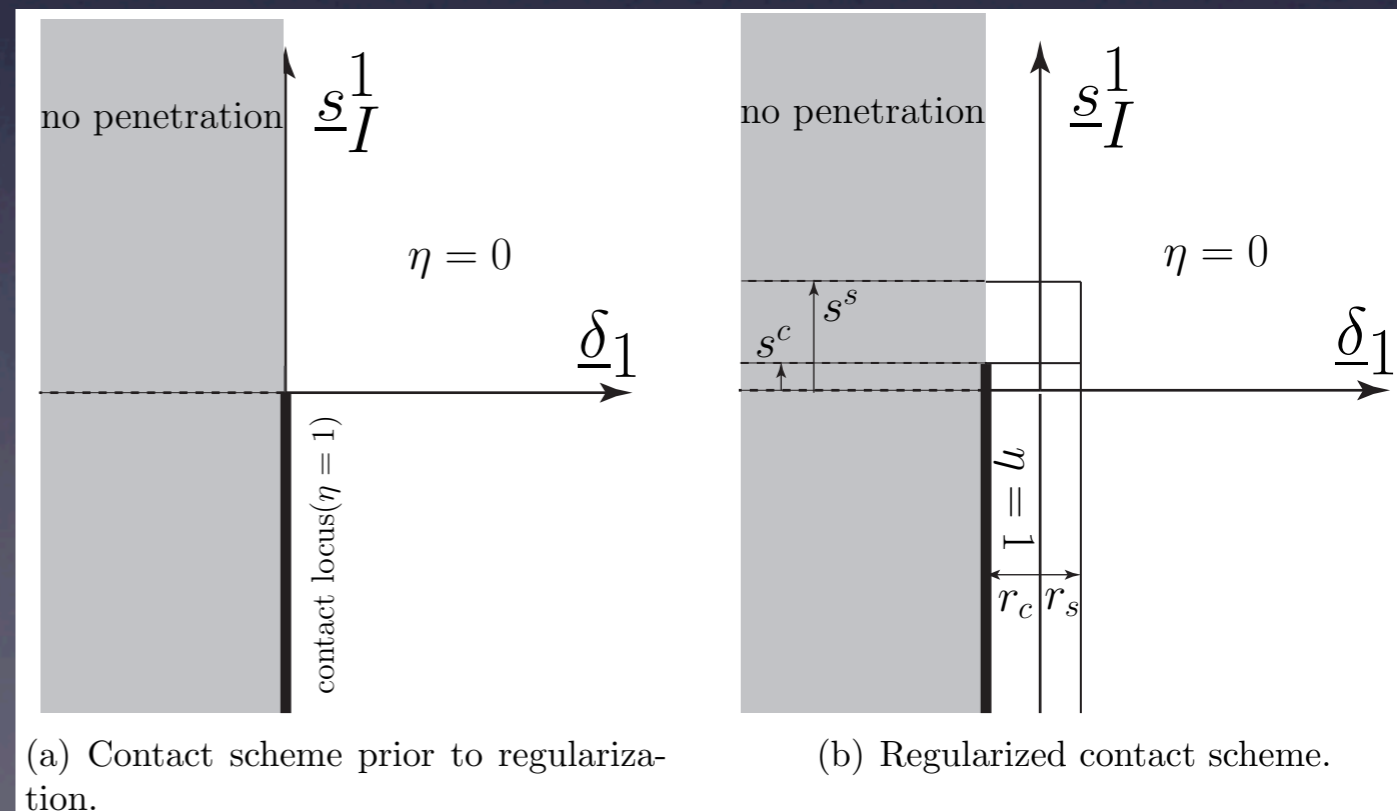
This model implies

$$D = \bar{D} \Rightarrow \dot{D} = 0$$

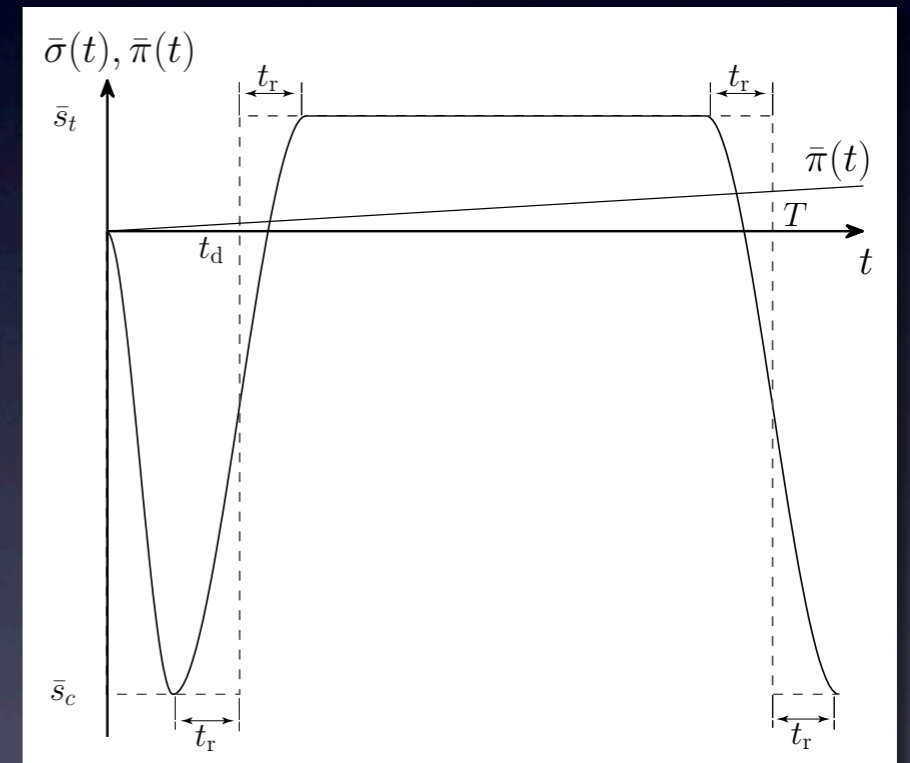
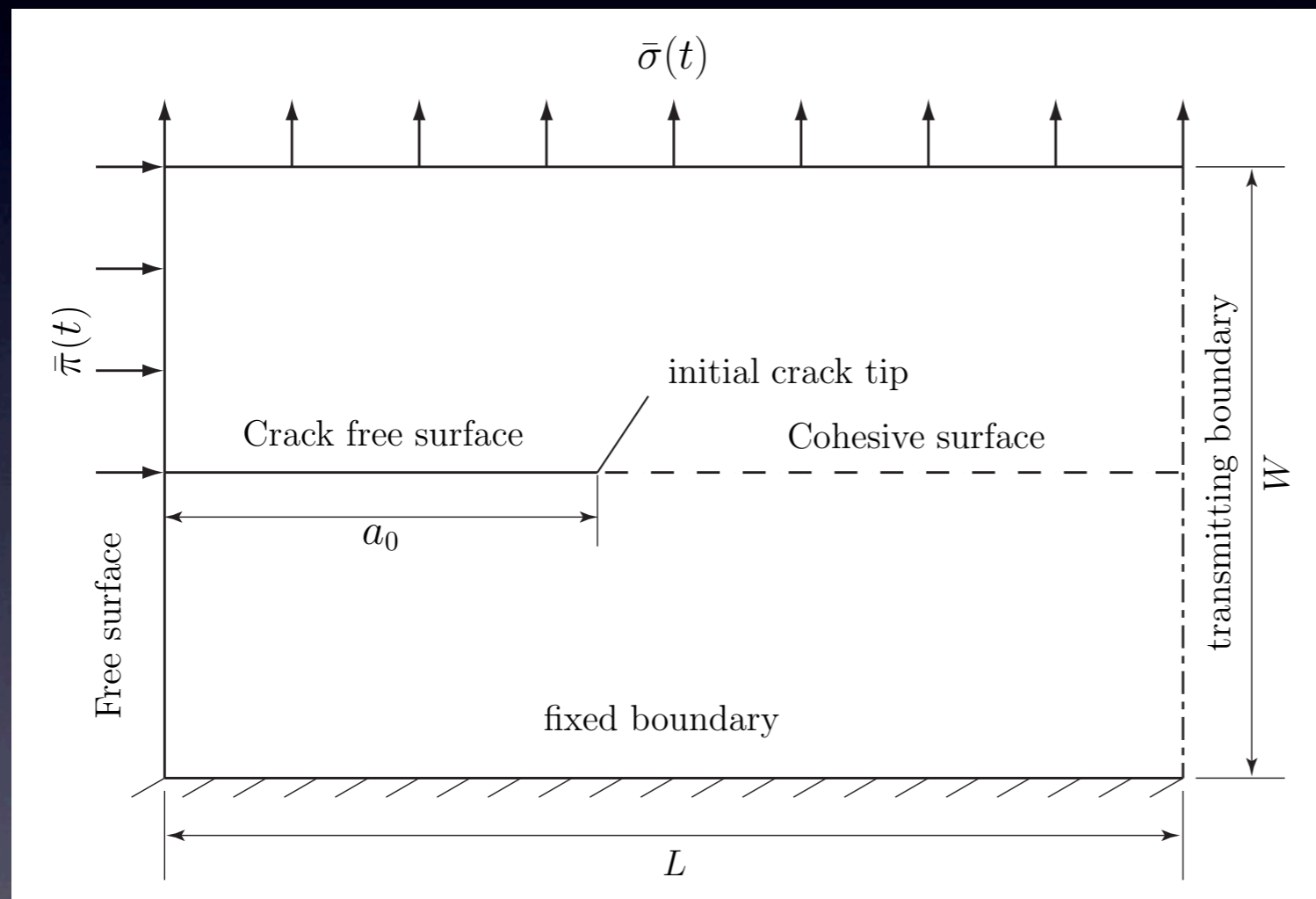
Mesososcopic subdivision

Contact area (η)

- Contact/Separation transition may introduce shocks in tractions and velocities
- Regularization is required for shock capturing schemes
- Regularization is based on **normal contact traction** and **macroscopic traction**
- A minimum separation value of r_c is ensured in the regularized model
- Surface toughness may provide physical regularization
- No regularization is required for stick-slip areas

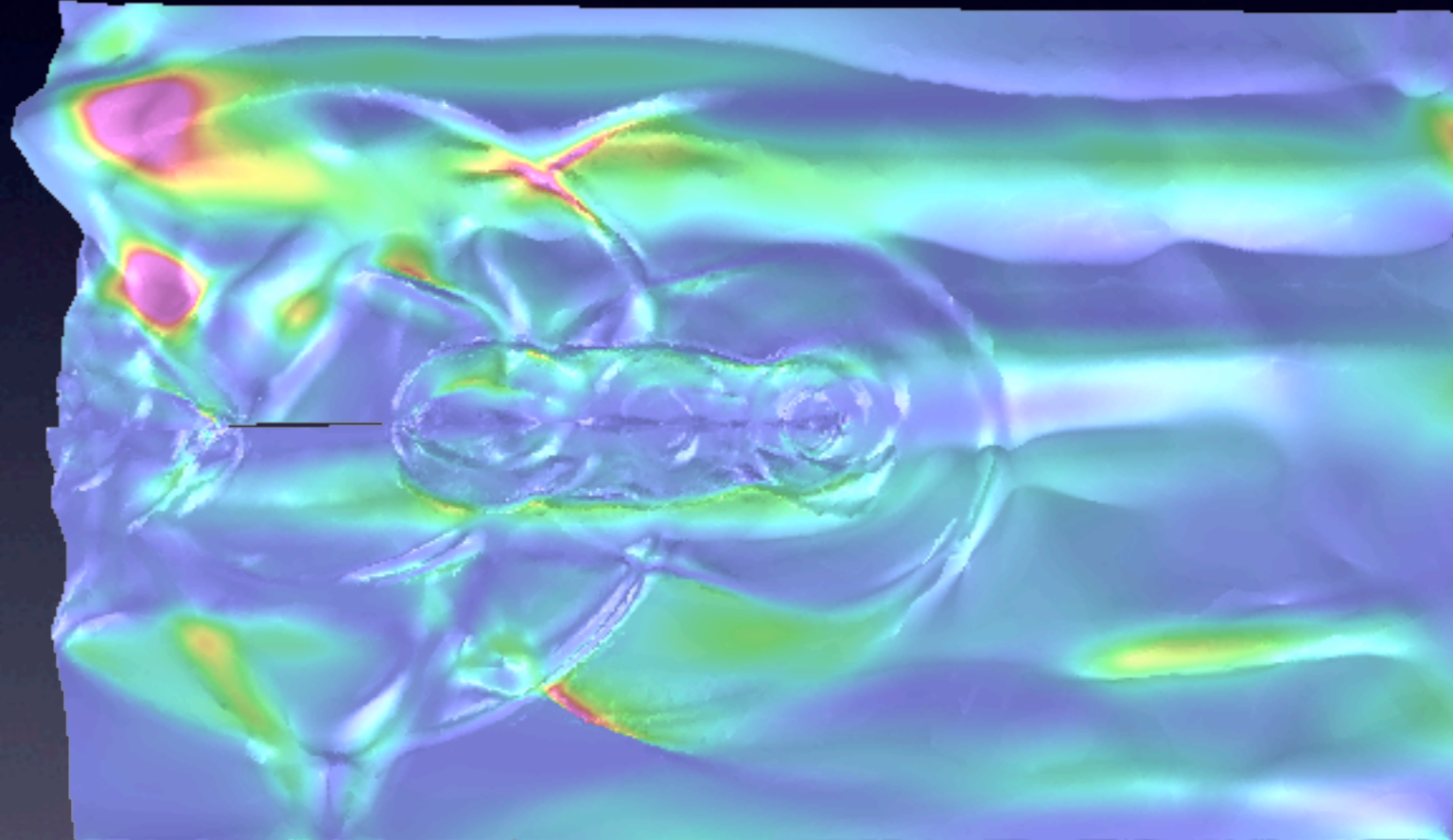


Crack closure example: brittle fracture under cyclic, dynamic loading conditions



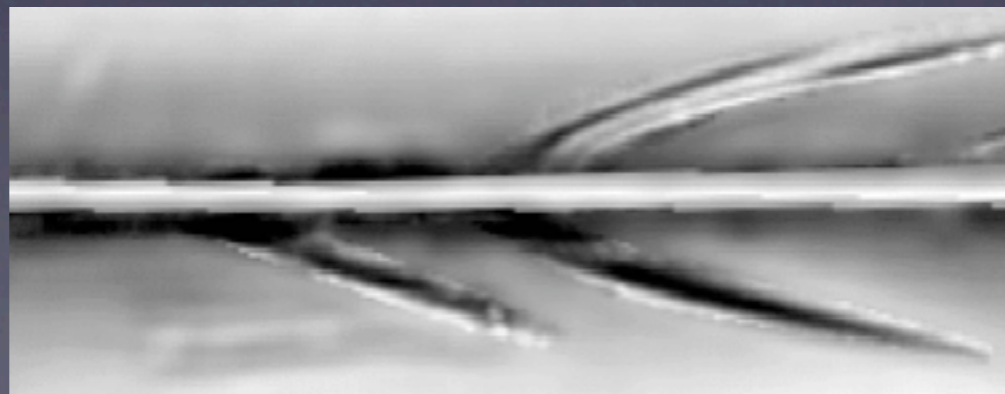
Crack closure example: brittle fracture under cyclic, dynamic loading conditions

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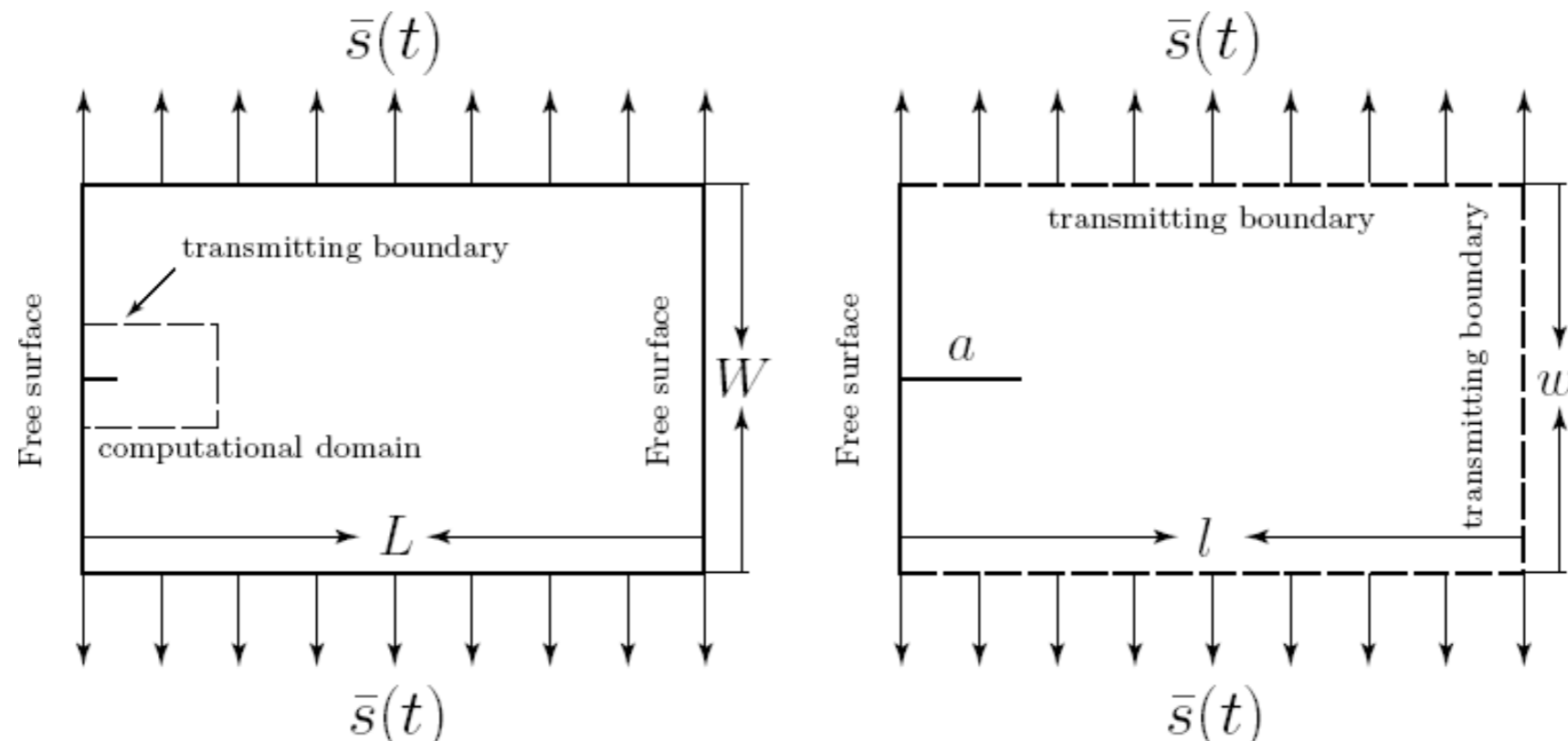


Solution-dependent crack paths

- Inclined tent poles, edge bisection and smoothing extend cohesive zone in any desired direction
- Effective stress criterion for extending cohesive surface
 - ◆ maximum value determines direction for extension
- Nucleation + coalescence = branching
 - ◆ Probabilistic nucleation criterion models random defects
 - ◆ Introducing cohesive surface does not affect behavior until damage accumulates



Solution-dependent crack propagation Geometry and Loading



(a) The experimental setting.

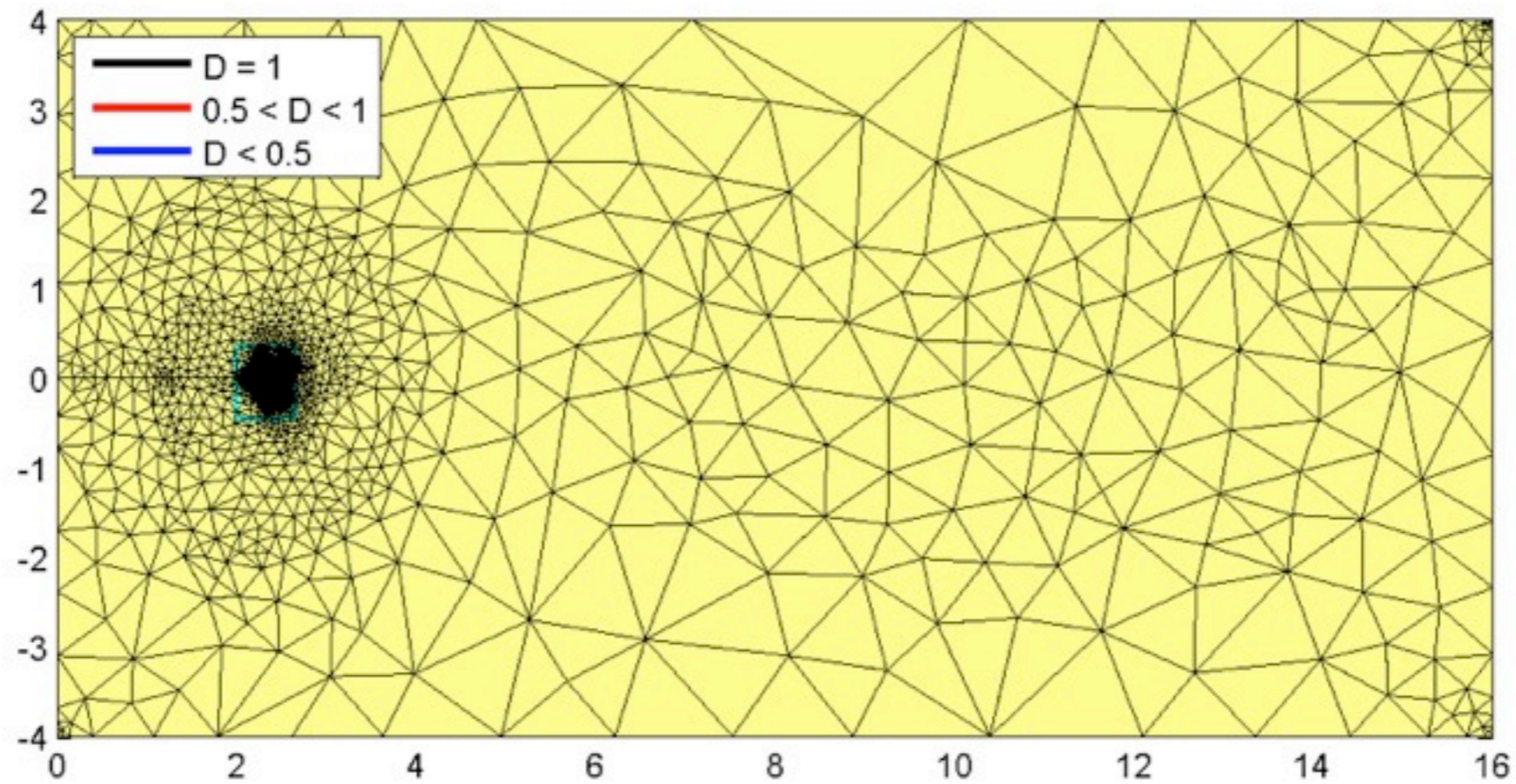
(b) The computational domain.

PMMA:

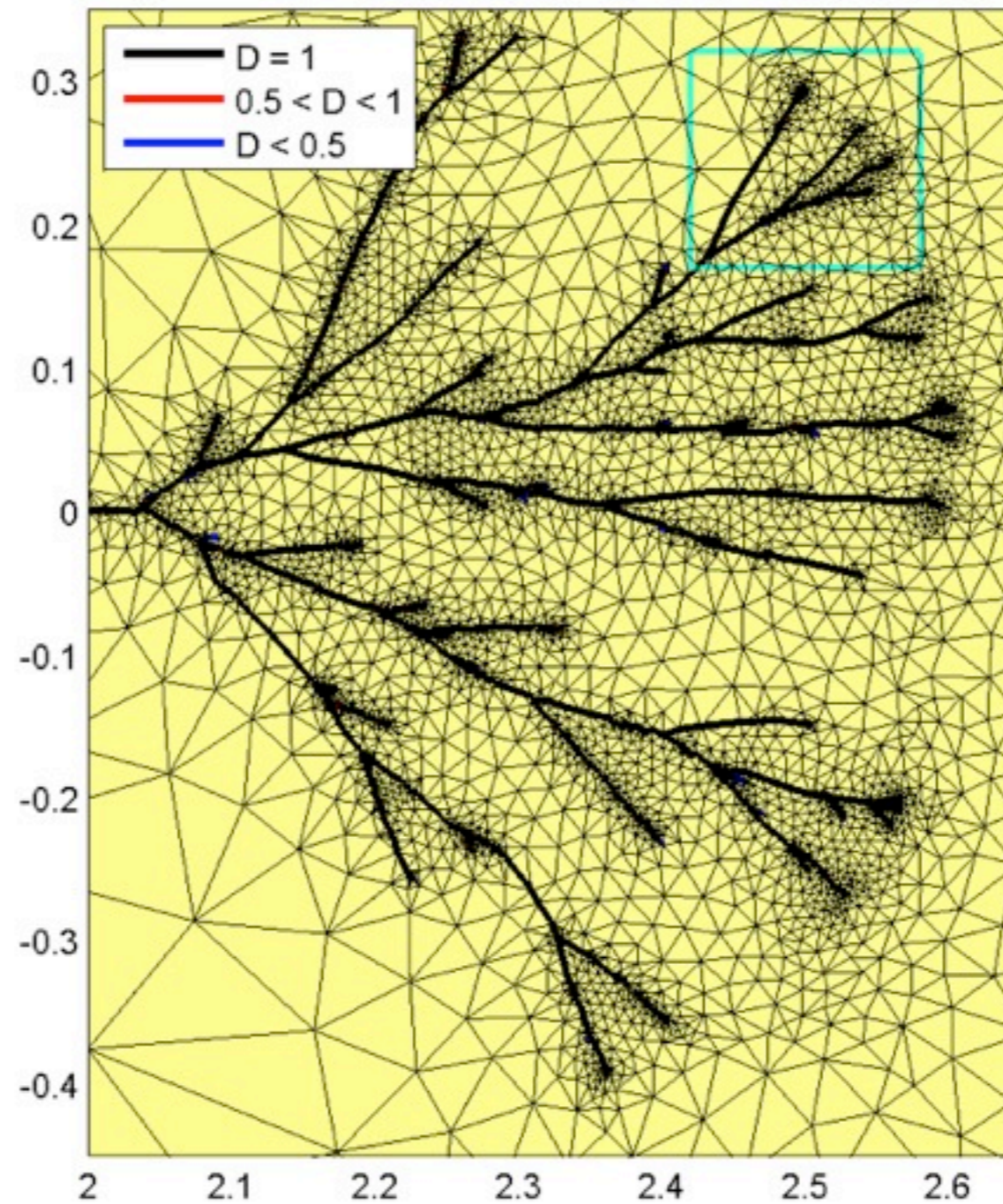
- $E = 3.24GPa.$
- $\rho = 1190 \frac{Kg}{m^3}.$
- $w = 0.020MPa.$
- $\tau = 10^{-2} \mu s.$

based on Sharon and Fineberg experiment (1996)

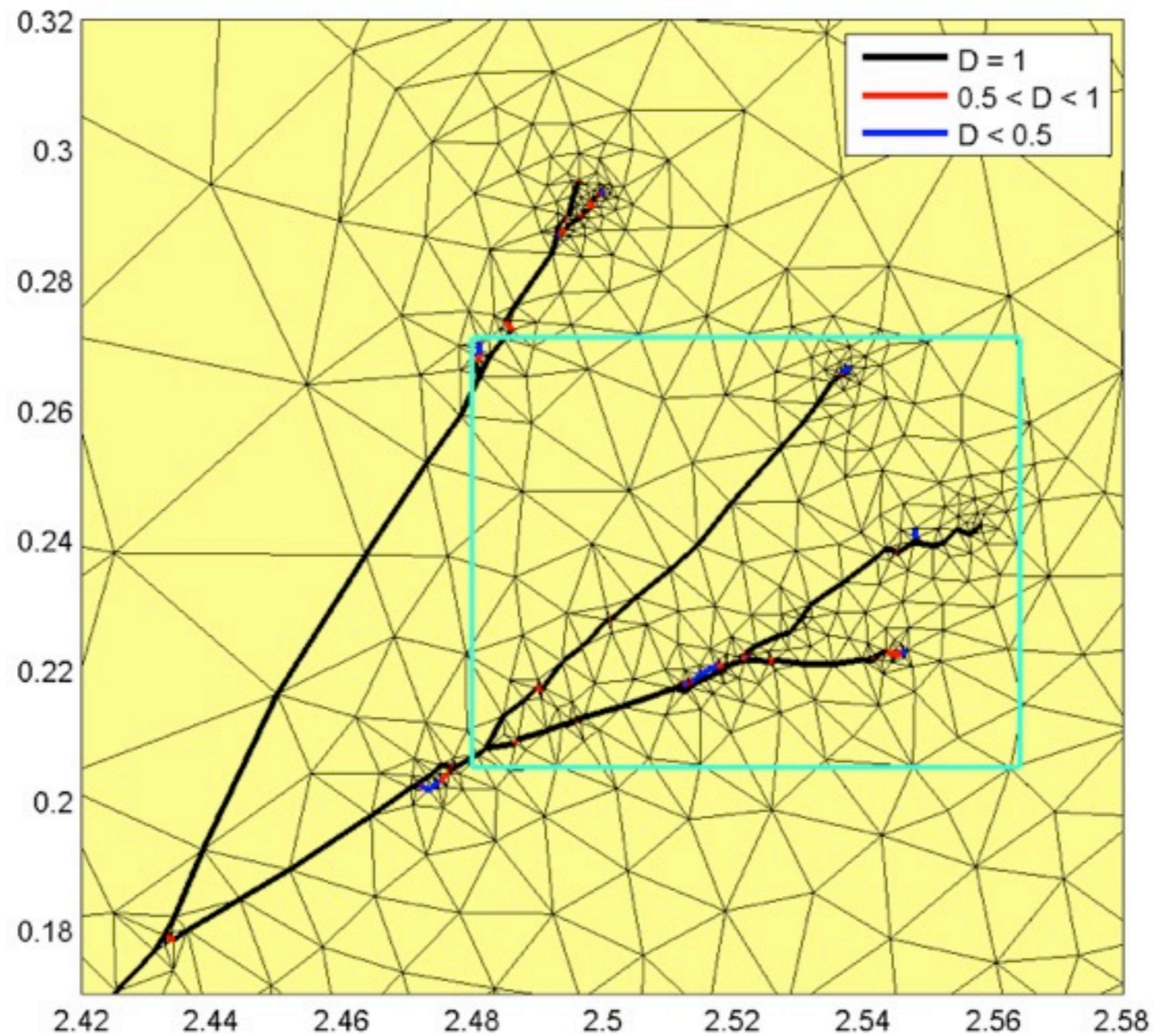
Refinement Details: 1 x zoom



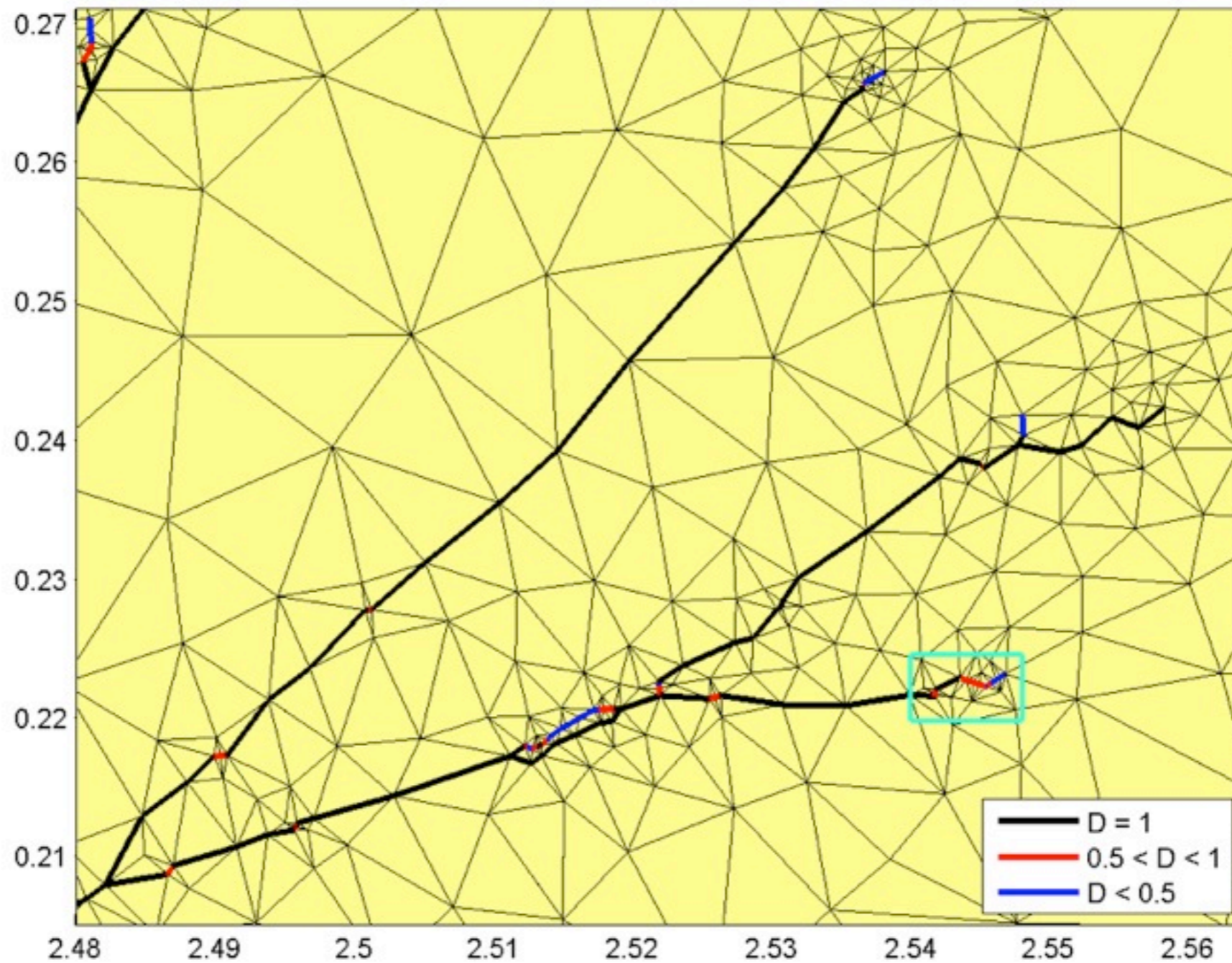
Refinement Details: 25 x zoom



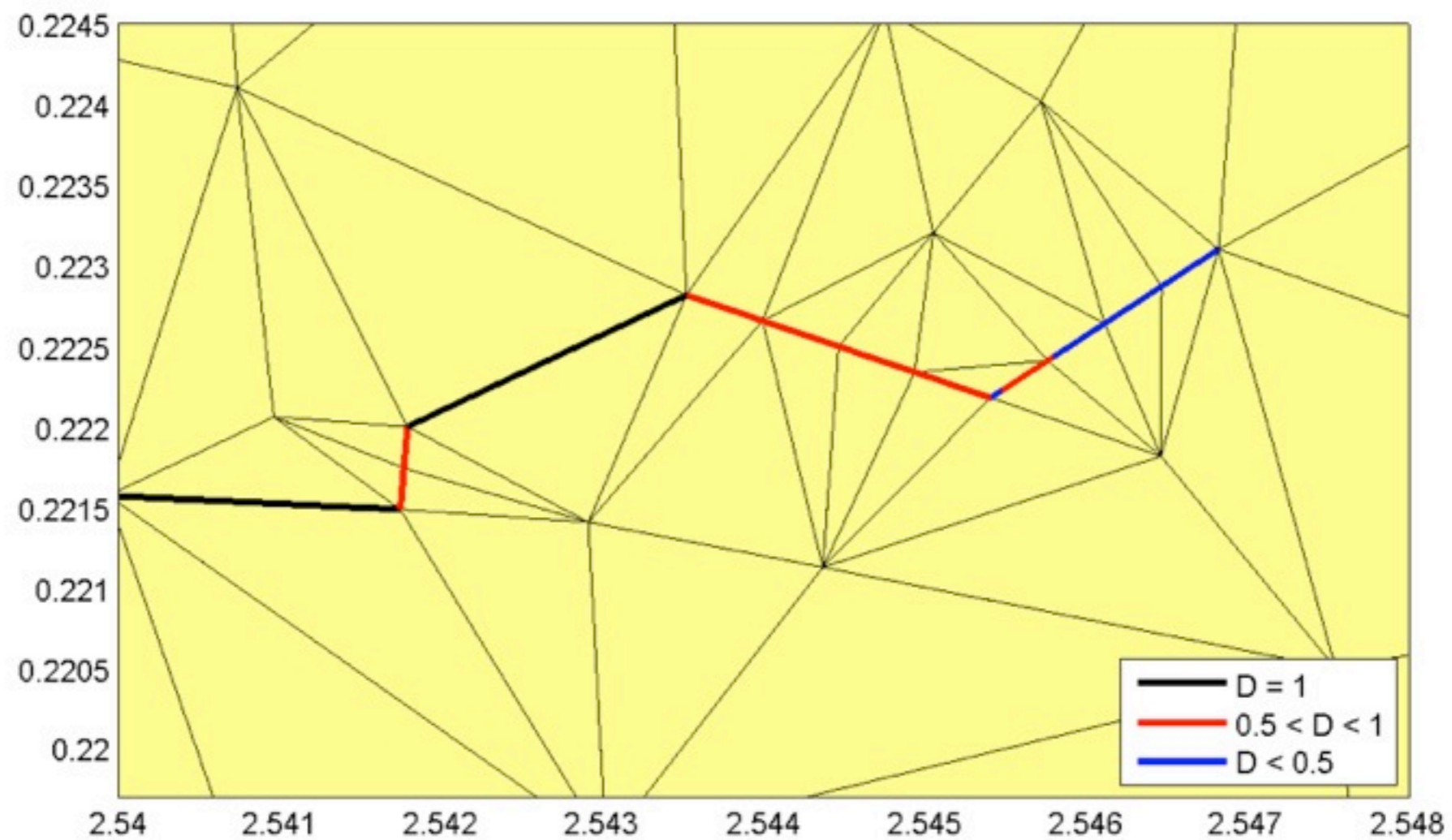
Refinement Details: 100 x zoom



Refinement Details: 200 x zoom

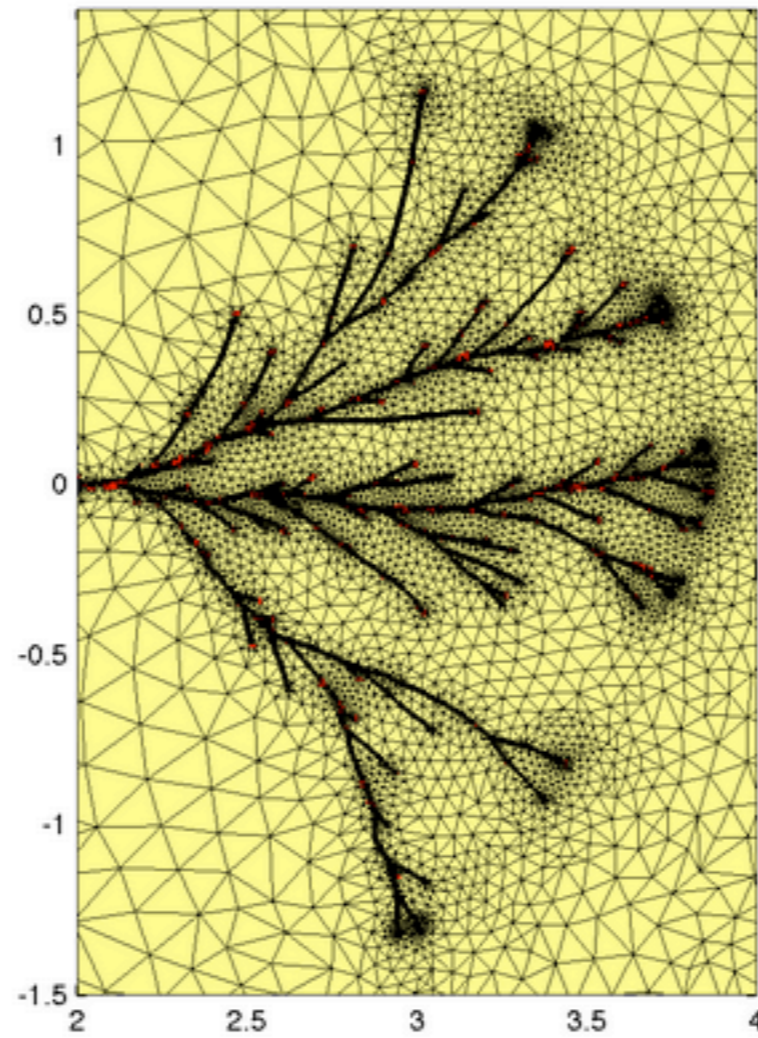


Refinement Details: 2000 x zoom

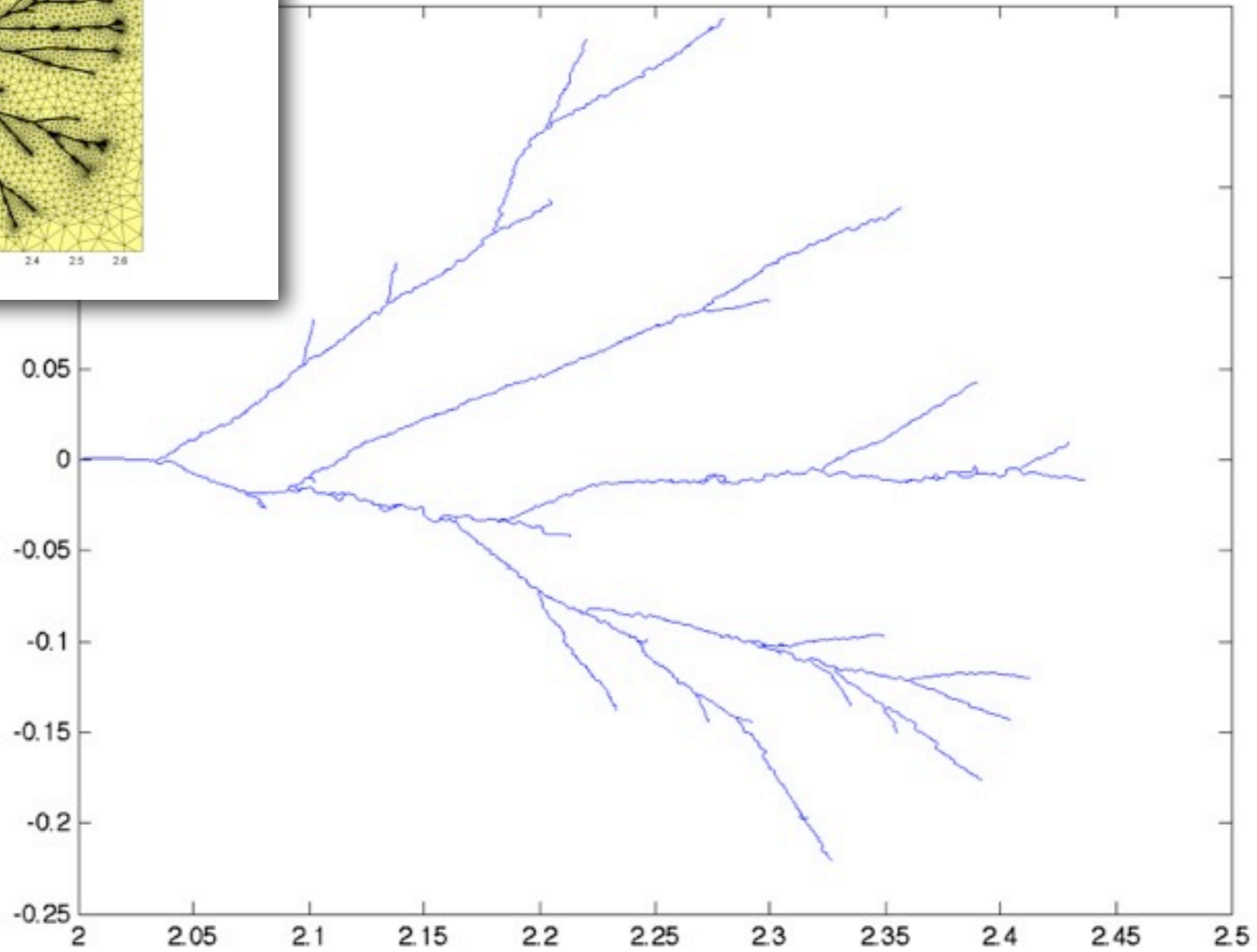
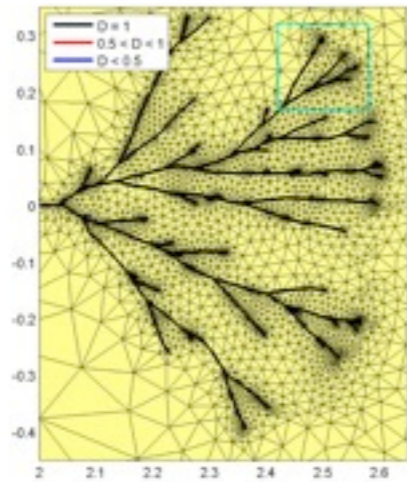


Same problem with crack closure contact model, random defects and nucleation

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Effect of contact conditions on crack path

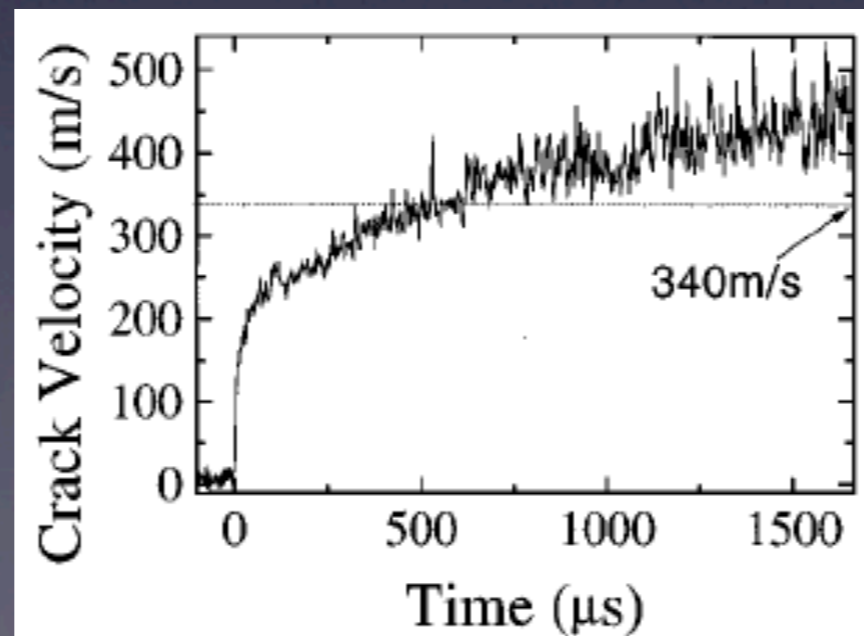


Surface roughening at high crack velocities

- Almost uniform hoop stress distribution at $\dot{a} > 2/3c_R$ (Yoffe, 1951)
- Experimentally observed roughness on crack surface (Smekal, 1953; Kerkhof, 1973; Sharon and Fineberg, 1996)
- Surface roughening and branching instabilities in dynamic fracture; Phenomenological wavy crack path (Gao, 1992)
- Dynamical stability of a propagating crack: Obrezanova (2002) at $\dot{a} > 1/3c_R$ crack may admit one or more oscillatory modes of instability
- Stability of dynamically propagating cracks in brittle materials: Uenishi *et. al.* (2001): Surface roughness induced by increase of crack velocity

Single crack velocity agreement with experiments

- Material degradation and energy loss due to induced heat should be contributed in the model (Sharon and Fineberg 1996)
- Bulk material dissipative and damage models to be incorporated
- Material degradation around the crack tip, the high inertia zone, and subsequent drop in wave speed (Ravi *et. al.*, 2007)
- Surface tension or hardness in computing the limiting velocity, Kerkhof (1997)



Summary

- Robust adaptive model for crack growth
 - ◆ Ensures accurate rendering of cohesive models
 - ◆ Eliminates all mesh-dependence effects
 - ◆ Supports nucleation
 - ◆ Captures multi-scale behavior
 - ◆ Captures the details of crack propagation (Branching, micro branching, surface roughening)
- New dynamic, continuum contact formulation
- Open physical modeling issues
 - ◆ Crack acceleration too rapid (typical of cohesive models)
 - ◆ Branching pattern does not completely match experiment
 - ◆ Some aspects heuristic, but have framework for developing