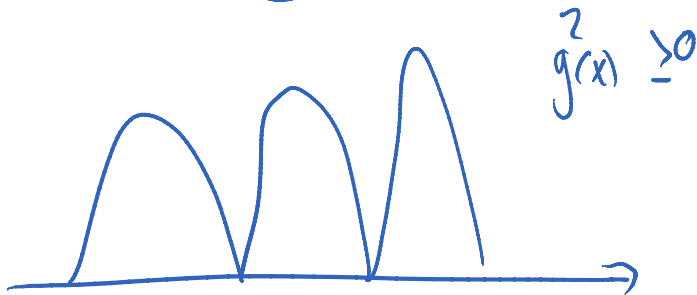
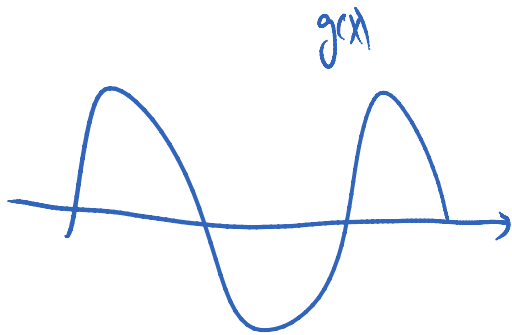
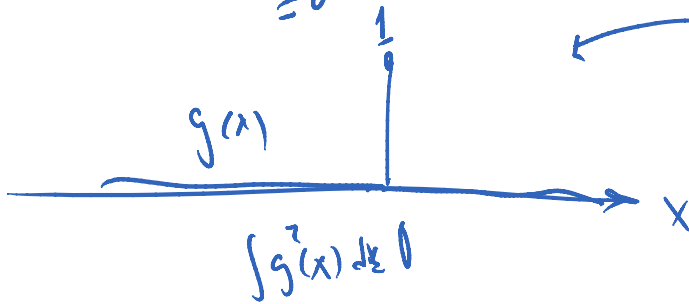


02/04/2014

Tuesday, February 04, 2014
11:58 AM



$$\int \underbrace{g(x)^2}_{\geq 0} dx = 0 \implies g(x) = 0 \text{ a.e.}$$

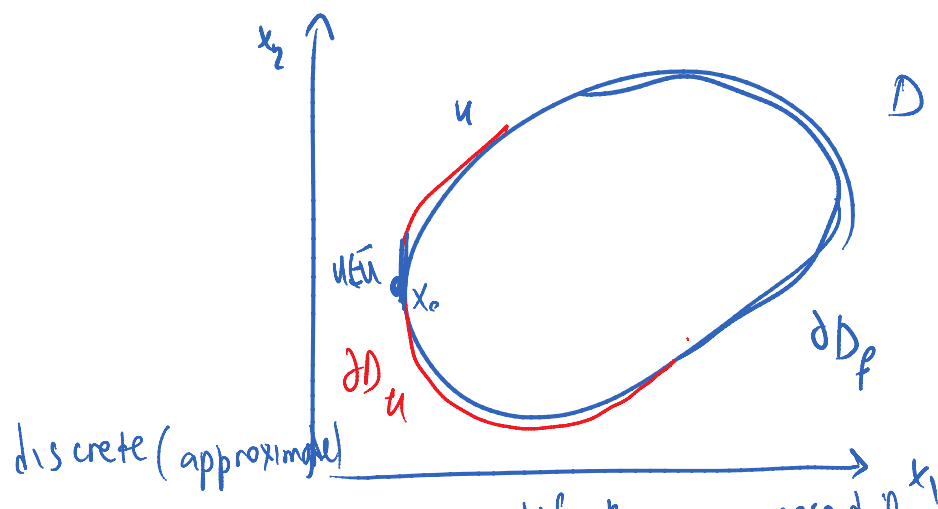
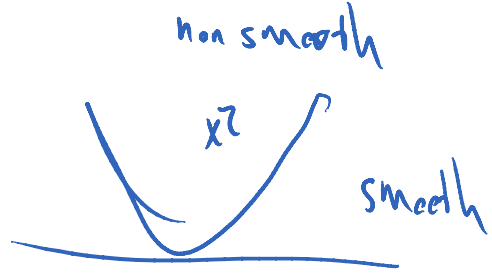
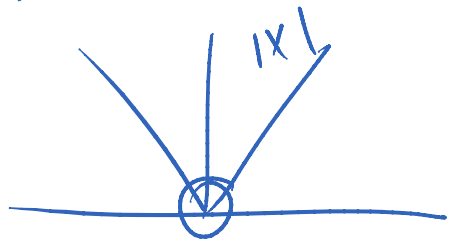


$g(x)$ is continuous

$$\int_D g(x)^2 dx = 0 \iff \forall x \in D: g(x) = 0$$

Why not this?

$$\int_D |g(x)| dx = 0 \iff \forall x \in D \quad g(x) = 0$$

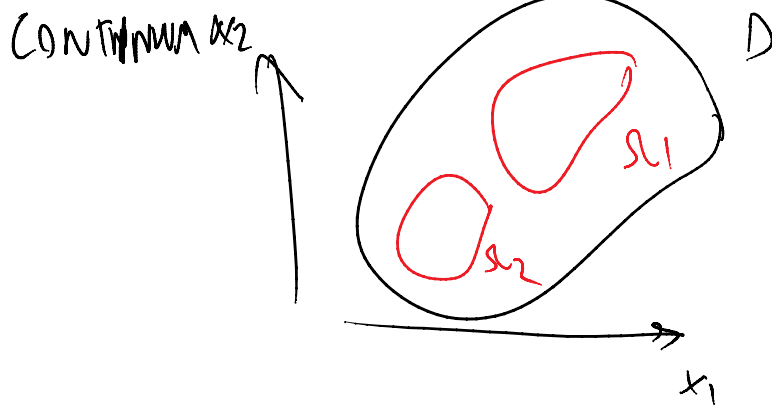


$u^h = \sum_{i=1}^n \underbrace{a_i}_{\substack{\text{n knowns} \\ \downarrow}} \underbrace{\phi_i(x)}_{\substack{\text{satisfy homogeneous essential} \\ \text{BC's } \phi_i(x_0) = 0}} + \underbrace{\phi_p(x)}_{\substack{\text{satisfies essential} \\ \text{boundary} \\ \text{condition } \phi_p(x_0) = \bar{u}(x_0)}}$

trial or test functions

$$u^h(x_0) = \bar{u}$$

1. We started with balance law:



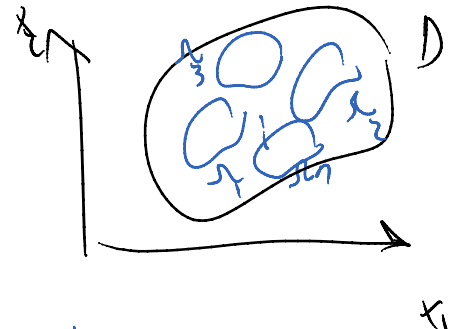
$\forall \Omega \subset D$

$$\int_{\partial \Omega} \vec{F} \cdot d\vec{s} - \int_{\Omega} \vec{r} \cdot dV = 0$$

\downarrow flux

DISCRETE

We choose n subdomains Ω_i $i=1, \dots, n$

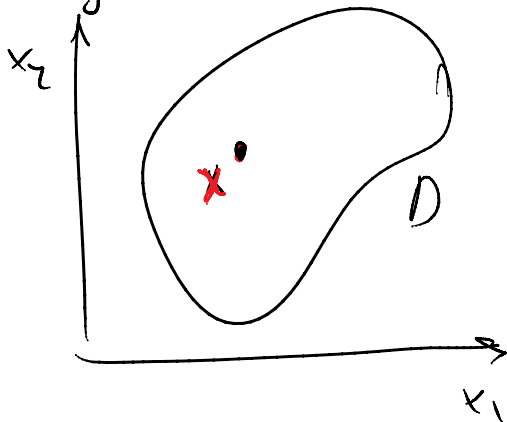


we choose n Ω_i

for $i=1, \dots, n$

$$\int_{\partial \Omega_i} \vec{F} \cdot d\vec{s} - \int_{\Omega_i} \vec{r} \cdot dV = 0$$

2. strong form



$\forall x \in \Omega$

$$\nabla \cdot \vec{F} + r = 0$$

$$L_M(u) - r = 0$$

In a discrete setting we choose n x_i :



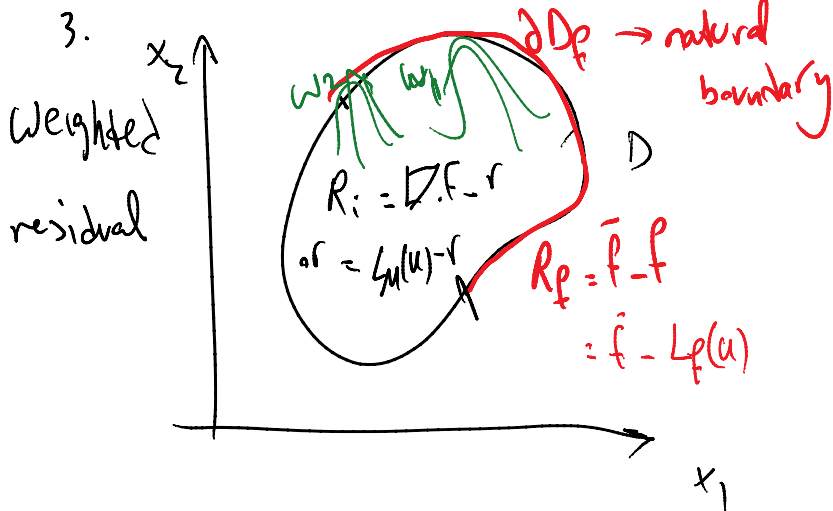
for $i=1, \dots, n$

$$\nabla \cdot \vec{F}(u(x_i)) - r = 0$$

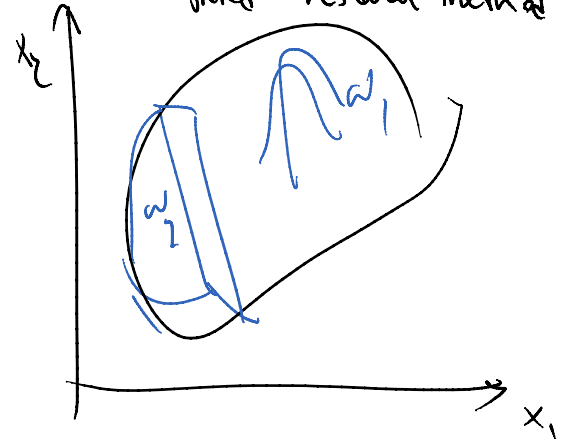
$$L_M(u) - r = 0$$

$$\nabla \cdot f(u(x_i)) - r = 0$$

$$L_M(u)(x_i) - r = 0$$



Discrete: weighted residual method



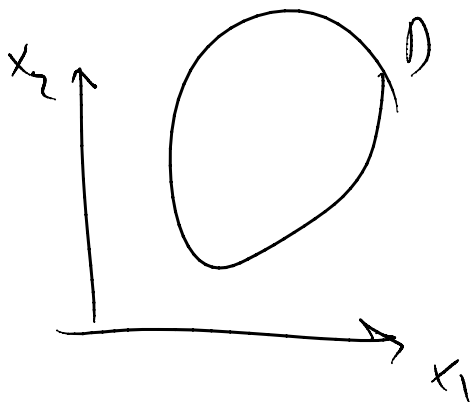
for $i=1, \dots, n$

$$\int_D W_i \cdot R_i \, dV + \int_{\partial D_f} \omega_i \cdot R_f \, dS = 0$$

$\forall w \in W:$

$$\int_D \omega \cdot R_i \, dV + \int_{\partial D_f} \omega \cdot R_f \, dS = 0$$

4. weak statement

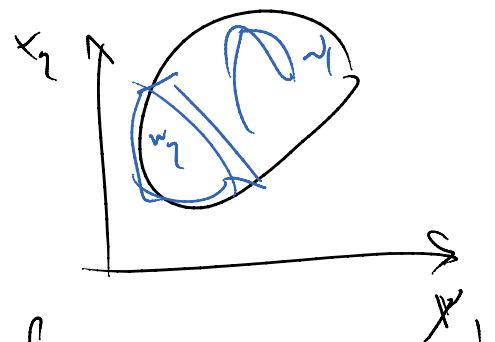


$\forall w$

$$\int L_M^f(w) L_M(u) \, dV =$$

discrete

instead of $\forall w$
we choose n w 's



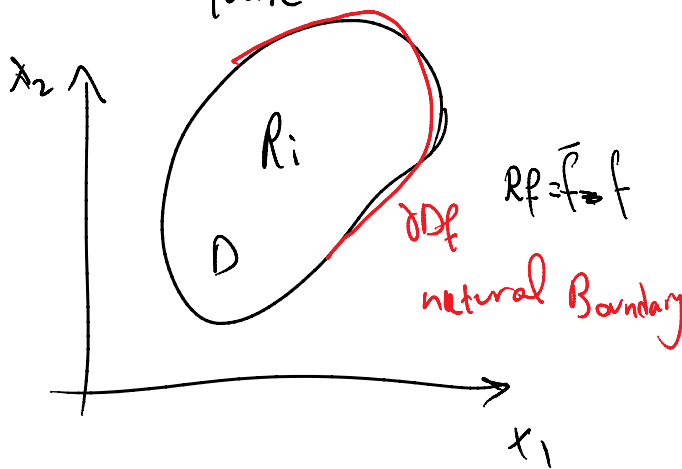
$i=1, \dots, n$

HW

$$\int_D L_m^t(\omega) L_m(u) dv = \int_{\partial D_f} \bar{f} ds$$

$$\int_D L_m^f(\alpha_i) L_m(u) dv + \int_{\partial D_f} \bar{f} ds = 0$$

5. not directly following 4
Least Square



$$R^2 \approx \int_D R_i^2 dv + \int_{\partial D_f} R_f^2 ds = 0$$

for all $\tilde{u} \in V$ (satisfies essential BC)

$$\widetilde{R^2(u)} \leq R(\tilde{u})$$

$$R^2(u^h) = \text{known}$$

$$R^2(\sum \alpha_i \phi_i(x) + \phi_p(x))$$

unknown

$$= R^2(a_1, \dots, a_n)$$

$$R^2(u^h) \leq R^2(\tilde{u}^h)$$

for all $\tilde{u}^h =$

$$\sum \tilde{\alpha}_i \phi_i(x) + \phi_p(x)$$

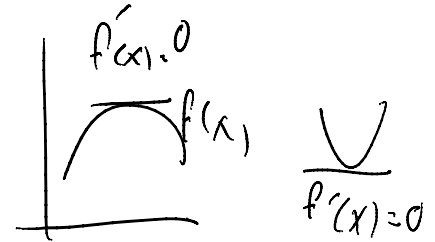
$$R^2(a_1, \dots, a_n)$$

$$R^2(a_1, \dots, a_n)$$

we want to minimize it

Q

So gradient of R^2 should be zero

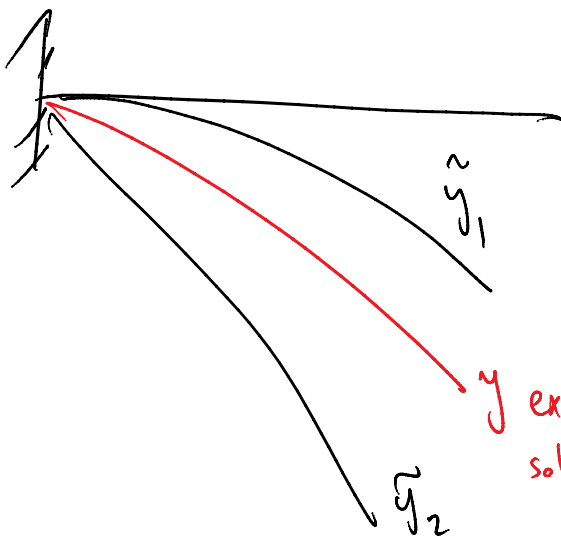


in discrete setting

$$\frac{\partial R^2(a_1, \dots, a_n)}{\partial a_i} = 0 \quad \forall i=1, \dots, n$$

6. (not following 1-5, it's an independent statement)

Energy method



for $y, y_2 \in V$ } satisfies essential

$$u^h(x) = \sum_{\text{unknowns}} (\tilde{a}_i) \phi_i(x) + \phi_p(x)$$

for discrete solution

$$u^h(x) = \sum (\tilde{a}_i) \phi_i(x) + \phi_p(x)$$

example

$$\mathcal{I}(u) = \int_0^L EA (u')^2 dx - \int_0^L q u(x)$$

$$\mathcal{I}(a_1, \dots, a_n)$$

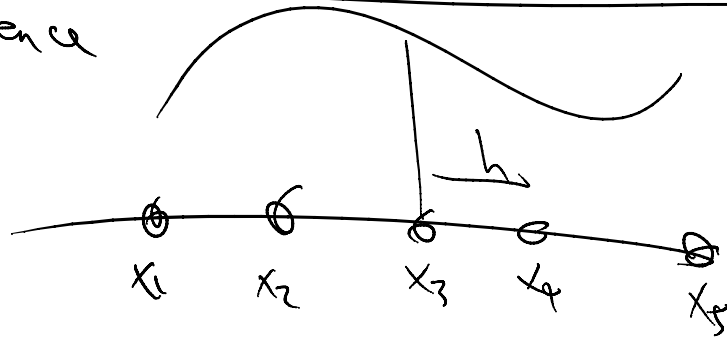
For $J, U = V$ [essential bc]

$$\mathcal{T}(a_1, \dots, a_n)$$

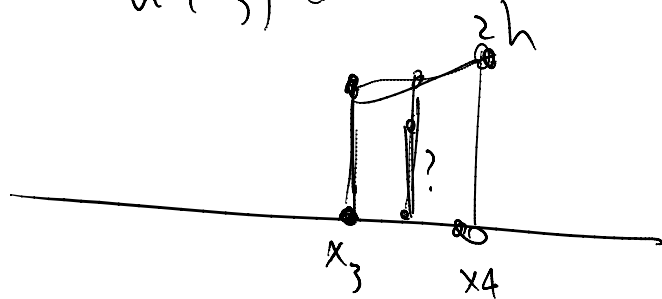
$$\mathcal{T}(y) \leq \mathcal{T}(\tilde{y})$$

$$\frac{\partial \mathcal{T}}{\partial a_i} = 0 \text{ for } i=1, \dots, n$$

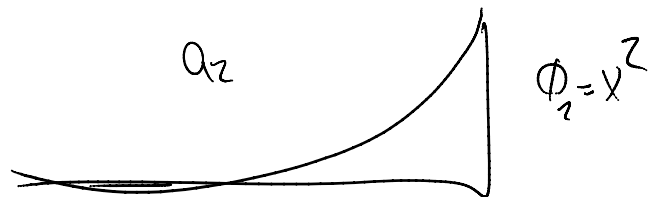
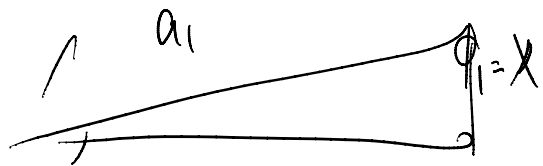
Finite difference

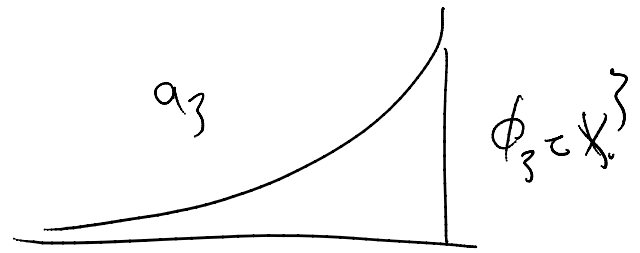


$$u'(x_3) = \frac{u(x_4) - u(x_2)}{2h}$$

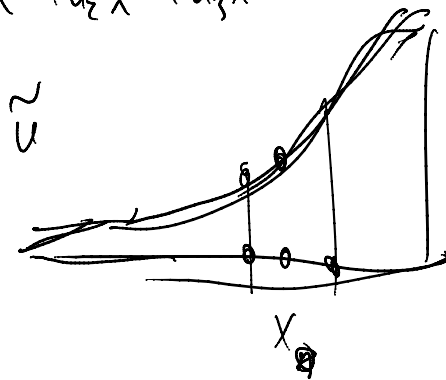


3:





$$u^h = a_1 \phi_1(x) + a_2 \phi_2(x) + a_3 \phi_3(x) = a_1 x + a_2 x^2 + a_3 x^3$$



$$u^h(x_0) = a_1 \phi_1(x_0) + a_2 \phi_2(x_0) + a_3 \phi_3(x_0)$$

as soon as we have a_i 's
we can "directly" evaluate the
function everywhere