

$$\begin{cases} e'_i = Q_{ij} e_j \\ v'_i = Q_{ij} v_j \end{cases}$$

$$Q = \begin{bmatrix} e'_1 & | & e'_2 & | & e'_3 \end{bmatrix} [e_1] [e_2] [e_3]$$

in () system $[v'] = Q[v]$

$$e_j = Q_{ij} e'_i$$

$$v_j = Q_{ij} v'_i$$

$$Q = \begin{bmatrix} e_1 \\ e_2 \\ e_3 \end{bmatrix}$$

in () system

$$[e] = Q^T [e']$$

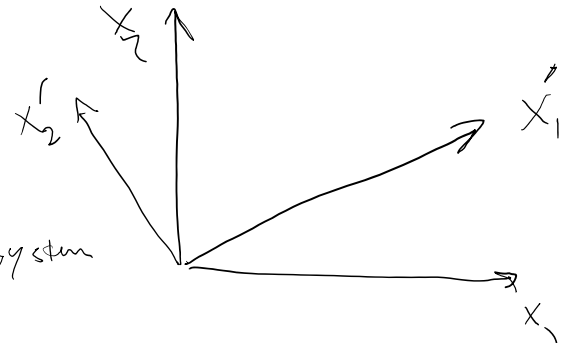
$$[v] = Q^T [v']$$

2D

$$Q = \begin{bmatrix} c & -s \\ +s & c \end{bmatrix} \quad \begin{matrix} c = \cos \theta \\ s = -\sin \theta \end{matrix}$$

Scalars $\lambda \in \mathbb{R}$ is a scalar

$T(x)$
 \downarrow
 temperature is a scalar



Scalars don't change from one coordinate system to another

Example:

Inner product is a scalar

- Part A

$$a = a_i e_i \quad b = b_j e_j$$

$a \cdot b = a_i b_i$

(*)
 $\{e_i\}$ orthonormal coordinate system

want to show that
 IF we had used (*) as definition of
 • then • is a scalar.

$$a \cdot b = (a_i e_i) \cdot (b_j e_j) = a_i b_j (e_i \cdot e_j)$$

$$= \begin{cases} a_i b_j \delta_{ij} = a_i b_i \\ a_i b_j g_{ij} \end{cases}$$

case 1
 orthonormal $\{e_i\}$
 case 2
 not covered here
 general coordinates
 for a coordinate system

metric matrix

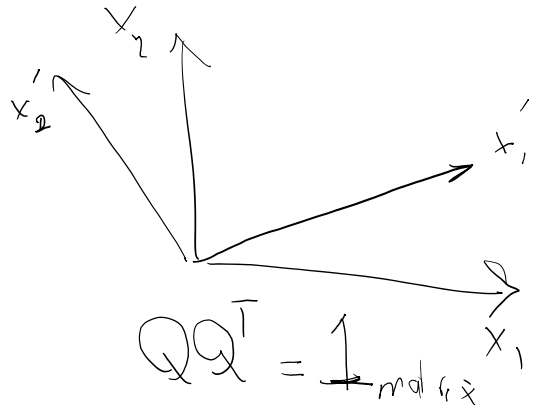
metric matrix for a coordinate system

$$\begin{aligned}
 a \cdot b &= [a_1 \dots a_n] \begin{bmatrix} b_1 \\ \vdots \\ b_n \end{bmatrix} = a_i b_i \quad \text{orthonormal } \{e_i\} \quad g = I \\
 &= [a_1 \dots a_n] \begin{bmatrix} g_{11} & \dots & g_{1n} \\ \vdots & \ddots & \vdots \\ g_{n1} & \dots & g_{nn} \end{bmatrix} \begin{bmatrix} b_1 \\ \vdots \\ b_n \end{bmatrix} = a_i g_{ij} b_j \quad g_{ij} = e_i \cdot e_j \\
 & \quad \text{nonorthonormal } \{e_i\}
 \end{aligned}$$

- Part B

$a_i b_i$ definition is coordinate independent

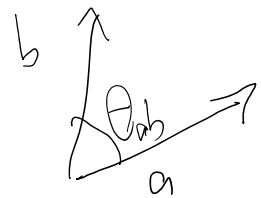
$$\begin{aligned}
 a_i b_i &= Q_{mi} Q_{ni} a'_m b'_n \\
 &= \underbrace{Q_{mi} (Q^T)_{in}}_{(QQ^T)_{mn}} a'_m b'_n \\
 &= \delta_{mn} a'_m b'_n =
 \end{aligned}$$



$$\Rightarrow \boxed{a_i b_i = a'_m b'_m}$$

The reason this works nicely, we defined inner product independent of vector components

$$a \cdot b = |a| |b| \cos \theta_{a,b}$$



A number that is not a scalar

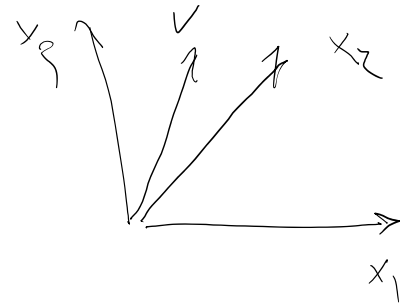
$v \cdot v \cdot e$ _____



A number that is ...

$$v = v_i e_i$$

$$Q_p(v) := \sqrt[p]{\sum_{i=1}^n |v_i|^p}$$

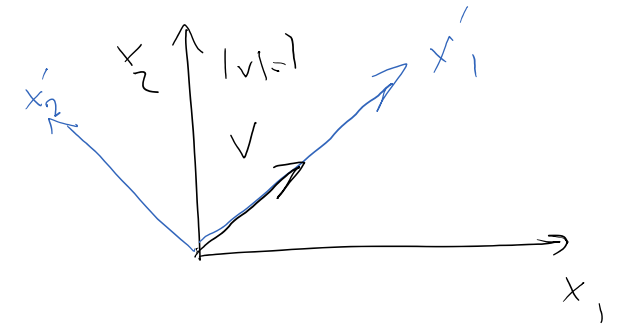


$Q_1(v) = |v_1| + |v_2| + |v_3|$ (Not a scalar!)

$$Q_2(v) = \sqrt{|v_1|^2 + |v_2|^2 + |v_3|^2} = \sqrt{v \cdot v} = |v| \text{ magnitude of } v \text{ scalar}$$

$$Q_1(v) = \frac{\sqrt{2}}{2} + \frac{\sqrt{2}}{2} = \sqrt{2} \text{ in } \{e_1, e_2\}$$

$$Q_1(v) = 1 + 0 = 1$$



Tensors

Coordinate transformation

0th order λ scalar

1st " v vector

~~2nd~~

2nd " A

$$v'_i = Q_{ij} v_j$$

$$A'_{ij} = Q_{im} Q_{jn} A_{mn}$$

4th order e.g. elasticity tensor C

$$C'_{ijkl} = Q_{in} Q_{jm} Q_{kp} Q_{lq} C_{mnpq}$$

2nd order tensors or
Linear operators (functions)

A linear operator L , is a function from vector space V to vector space W that satisfies the following two properties

$$L: V \rightarrow W$$

1 / 1, . . . , 1, Combined

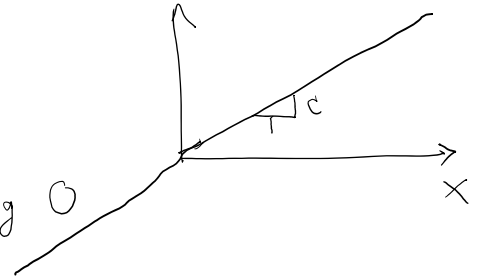
$$L: V \rightarrow W$$

$$\begin{aligned} \forall v_1, v_2 \quad L(v_1 + v_2) &= L(v_1) + L(v_2) \\ \forall v \in V, \lambda \in \mathbb{R} \quad L(\lambda v) &= \lambda L(v) \end{aligned} \left\{ \begin{array}{l} \text{Combined} \\ L(\alpha v_1 + \beta v_2) \\ = \alpha L(v_1) + \beta L(v_2) \end{array} \right. \\ \text{or} \\ L(v_1 + \beta v_2) \\ = L(v_1) + \beta L(v_2) \end{aligned}$$

Example 1D $\mathbb{R} \rightarrow \mathbb{R}$

$$\forall x \quad f(x) = f(x-1) = \underbrace{\alpha}_{\text{slope}} \underbrace{f(1)}_c = cx$$

Linear function in \mathbb{R}^n they are planes passing through 0



For any Linear function $L(0) = 0$

$$0 + L(0) = L(0+0) = L(0) + L(0) \Rightarrow \boxed{L(0) = 0}$$

Example 2

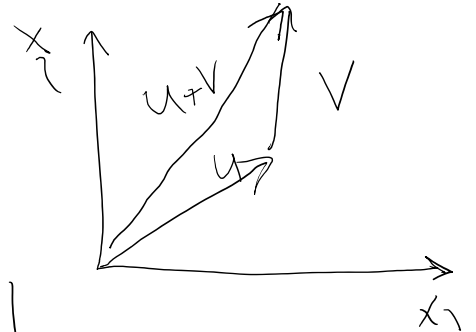
$$L(u) = |u| \quad \mathbb{R}^2 \rightarrow \mathbb{R}$$

is this linear

$$L(\vec{0}) = 0 \quad \text{smiley face}$$

But this is not a linear function

$$L(u+v) = |u+v| \quad \not\leq \quad \underbrace{|u|}_{L(u)} + \underbrace{|v|}_{L(v)}$$



Not a linear function

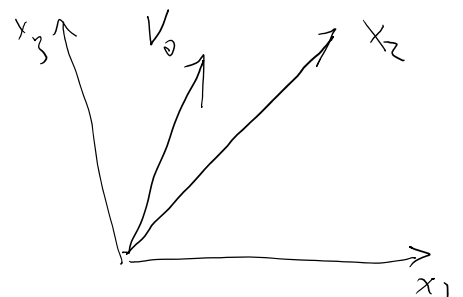
Example 3 $v = \begin{bmatrix} 1 \\ 2 \\ 3 \end{bmatrix}$

$$L(u) = v \cdot u$$

inner product

$$\underbrace{\mathbb{R}^3}_{\mathbb{R}^2} \rightarrow \mathbb{R}$$

$\langle v, \cdot \rangle$



inner product



$$L(u) = [1 \quad 2 \quad 3] \begin{bmatrix} u_1 \\ u_2 \\ u_3 \end{bmatrix} = u_1 + 2u_2 + 3u_3$$

$$L(u + \beta v) = (u + \beta v)_1 + 2(u + \beta v)_2 + 3(u + \beta v)_3 = (u_1 + 2u_2 + 3u_3) + \beta(v_1 + 2v_2 + 3v_3) = L(u) + \beta L(v)$$

In fact all linear operators from $\mathcal{V} \xrightarrow{\text{linear}} \mathbb{R}$ can be written in the form of inner product

$$\forall v \in \mathcal{V} \quad L(v) = v_{\bar{a}} \cdot v$$

There is one remaining point to consider before proceeding on to second-order tensors. That is, the inner product allows us to interpret a (Euclidean) vector as a linear operator that maps a vector into a real number (scalar). In fact, this is the defining property of a *first-order tensor*, and vectors are indeed first order tensors. To illuminate this point, let \mathcal{V} be the set of all vectors in some Euclidean point space \mathcal{E} . Now consider a specific vector $\bar{a} \in \mathcal{V}$, where the overbar indicates that we hold \bar{a} fixed. We can define a function $f_{\bar{a}}$ that maps a vector into a scalar by taking the inner product of \bar{a} and any vector $b \in \mathcal{V}$. That is,

$$f_{\bar{a}}(b) \equiv \bar{a} \cdot b.$$

A review of the properties of the inner product shows that $f_{\bar{a}}$ is indeed a linear operator. In fact, the *Riesz representation theorem* states that every linear function on \mathcal{V} to \mathbb{R} can be represented in this fashion (by varying our choice of the fixed vector \bar{a})! We use a similar approach in the next section to define second-order tensors as a special class of linear operators.

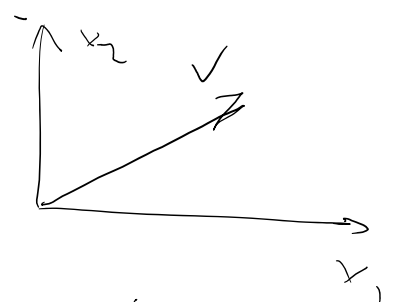
Related to this topic (suitable for a term project)

Concept of vectors, covectors, dual basis, ...

$$V = v^i \underline{e}_i$$

$$L(V) = a \cdot V$$

$$a = a^i e_i$$



e_i 's are orthonormal

$$a = a^i e_i$$

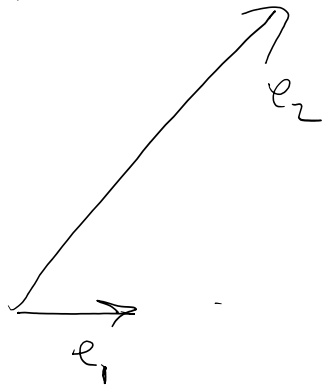
orthonormal

$$L(v) = (v^i e_i) (a^j e_j) = v^i a^j \underbrace{(e_i \cdot e_j)}_{\text{metric matrix}} = \int \begin{matrix} g_{ij} & e_i \text{ orthonormal} \\ g_{ij} & \text{"Not"} \end{matrix}$$

$$L(v) = v^i a^j g_{ij}$$

Dual basis $e^i(e_j) = \delta^i_j$

If the basis is orthonormal e^i is the same as e_i



...

We can represent linear operators $V \rightarrow \mathbb{R}$ by dual basis

$$L = l_i e^i$$

$$L(v) = (l_i e^i) (v^j e_j) = l_i v^j \underbrace{(e^i \cdot e_j)}_{\delta^i_j} = l_i v^i$$

In contrast if we had written L in terms basis

$$L = l^i e_i$$

$$L(v) = l^i v^j g_{ij}$$

$$\delta_{ij} \quad F^i_j$$

$$T_{ij} \quad T^i_j \quad T^{\tilde{ij}} \quad T^i_i$$

For the rest of the course we use orthonormal basis, so indices up or down don't make any difference

Linear functions (operators) themselves form a vector space

space of all linear operators from $V \rightarrow W$

$$S, T, U \in \mathcal{L} : V \rightarrow W$$

$$\lambda, \mu \in \mathbb{R}$$

addition properties

$$\begin{aligned} S + T &= T + S \\ S + (T + U) &= (S + T) + U \\ S + \underline{0} &= S \end{aligned}$$

$$\begin{aligned} S + T &= ? \\ \lambda S &= ? \\ 0 &= ? \end{aligned}$$

scalar product properties

$$\begin{aligned} (\lambda\mu)S &= \lambda(\mu S) \\ \lambda(S + T) &= \lambda S + \lambda T \\ (\lambda + \mu)S &= \lambda S + \mu S \\ 1S &= S \end{aligned}$$

$$(S + T)(u) := S(u) + T(u)$$

$$(\lambda S)(u) := \lambda(S(u))$$

$$0(O_V) = O_W$$

Linear operator zero

You need to show that indeed if $\lambda \in \mathbb{R}, S, T \xrightarrow{\text{linear}} V \rightarrow W$
 so are $S + T$
 λS

$$\begin{aligned}
 (S+T)(u+\beta v) &= (S+T)(u) + \beta(S+T)(v) \\
 &= S(u+\beta v) + T(u+\beta v) && \text{def. } S+T \\
 &= (Su + \beta Sv) + (Tu + \beta Tv) && S, T \text{ are linear} \\
 &= (Su + Tu) + \beta(Sv + Tv) && 3 \text{ properties of vector space } W \\
 &= (S+T)(u) + \beta(S+T)(v) && \text{def } S+T
 \end{aligned}$$

Dyadic product of two vectors:

$$\begin{cases}
 u \cdot v = [u_1 \dots u_n] \begin{bmatrix} v_1 \\ \vdots \\ v_n \end{bmatrix} = u^T v = u_i v_i \\
 u \otimes v = \begin{bmatrix} u_1 \\ \vdots \\ u_n \end{bmatrix} \begin{bmatrix} v_1 & \dots & v_n \end{bmatrix} = uv^T = \begin{bmatrix} u_1 v_1 & u_1 v_2 & \dots & u_1 v_n \\ u_2 v_1 & & & \\ \vdots & & & \\ u_n v_1 & \dots & \dots & u_n v_n \end{bmatrix}
 \end{cases}$$

motivation

Definition of dyadic product:

for vectors $u, v \in V$ we define dyadic product as an operator from $V \rightarrow V$ defined as

$$\forall u \in V \quad (u \otimes v) \left(\underbrace{w}_{\text{vector in } V} \right) = u \underbrace{(v \cdot w)}_{\text{scalar}}$$

we need to show that $u \otimes v$ is in fact a linear operator

$$\begin{aligned}
 (u \otimes v)(\alpha a + \beta b) &= u \{ v \cdot (\alpha a + \beta b) \} && u, v, a, b \in V \\
 &= u \{ \alpha (v \cdot a) + \beta (v \cdot b) \} && \alpha, \beta \in \mathbb{R} \\
 &&& \text{dist. \& homog. of inner product}
 \end{aligned}$$

$$\begin{aligned}
 &= u \{ \alpha (v \cdot a) + \beta (v \cdot b) \} && \text{dist. \& homog. of inner product} \\
 &= \alpha (u(v \cdot a)) + \beta (u(v \cdot b)) && \text{dist. \& scalar product properties} \\
 &= \alpha (u \otimes v(a)) + \beta (u \otimes v(b)) && \text{of } V
 \end{aligned}$$

$L(\alpha a + \beta b) = \alpha L(a) + \beta L(b)$ $L = u \otimes v$ so $u \otimes v$ is
 in fact a linear operator

$$e_1 = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} \quad e_3 = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

motivation

$$e_1 \otimes e_3 = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} \begin{bmatrix} 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad e_1 \otimes e_3 = (e_3 \otimes e_1)^T$$

$$e_3 \otimes e_1 = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{bmatrix}$$

$u \otimes v \neq v \otimes u$ in fact as we'll see later
 $u \otimes v = (v \otimes u)^T$

Motivation

$$T = \begin{bmatrix} T_{11} & T_{12} & T_{13} \\ T_{21} & T_{22} & T_{23} \\ T_{31} & T_{32} & T_{33} \end{bmatrix} = T_{11} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} + T_{12} \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} + T_{13} \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \\
 + \dots + T_{33} \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$\begin{aligned}
 &= T_{11} e_1 \otimes e_1 + T_{12} e_1 \otimes e_2 + \dots \\
 & \quad T_{33} e_3 \otimes e_3
 \end{aligned}$$

$$T = T_{ij} e_i \otimes e_j$$

are the basis for 2nd order tensors
 same way e_i are basis for vectors

$$v = v_i e_i$$

define $\bar{T}_{ij} = e_i(T e_j)$

show $T = \bar{T}_{ij} e_i \otimes e_j$

Dual basis (~~F~~) $T = \bar{T}_{ij} e^i \otimes e^j = \bar{T}^{ij} e_i \otimes e_j$
 $= \bar{T}_i^j e^i \otimes e_j = \bar{T}^i_j e_i \otimes e^j$

$$J = \sigma_{ij} e^i \otimes e^j$$

$$F = F^i_j e_i \otimes e^j$$