CM2021/09/28 Tuesday, September 28, 2021 4:29 PM

Cayley-Hamilton equation: Characteristic equation of the tensor:

$$-\delta^{3} + I_{1} \delta^{2} - I_{2} \delta + I_{3} = 0$$

$$I_{1} = tr(S) = S_{1} + S_{22} + S_{33}$$

$$I_{2} = \frac{1}{2} \left[(tr(S))^{2} - tr(S^{2}) \right] = \frac{1}{2} (S_{11} S_{11} - S_{11} S_{11})$$

$$I_{3} = det S_{1,1}$$



Cayley-Hamilton theorem: For a diagonizable matrix, it too satisfies its characteristic equation. For these symmetric 3x3 matrices

$$-6^{3}+5,6^{2}-5,6+5,1=0$$

$$-S^{3}+5,8^{2}-5,8+5,2+5,2=0$$

$$C^{0}$$

Representation of symmetric tensors

$$(1) 3e_{1}^{r} = 0 e_{1}^{r} e_{1}^{r}$$

$$(2) 3e_{1}^{r} = 0 e_{1}^{r} e_{1}^{r}$$

$$(3) 3e_{1}^{r} = 0 e_{1}^{r} e_{1}^{r}$$

$$(3) 3e_{1}^{r} = 0 e_{1}^{r} e_{1}^{r}$$

$$(3) 3e_{1}^{r} = 0 e_{1}^{r} e_{1}^{r}$$

$$(4) 3e_{1}^{r} = 0 e_{1}^{r} e_{1}^{r}$$

$$(5) 6e_{1}^{r} e_{1}^{r}$$

$$(6) e_{1}^{r} e_{1}^{r}$$



another or the normal coordinate system ____ 1e') is we ran to coordinate promotor making $Q : \begin{bmatrix} -e^{\pi} \\ e^{\pi} \\ 0 \\ e^{\pi} \end{bmatrix}$ pet. t ~ (e, e, e,) Tij = Qin Qjn Tm. Simlarly Sty: Que Qin Sm What are the components of S in ki, en es sy strong Weston: Tij = e; Ten components of Time eves-Real $S_{ij} = e_i^{\bullet} \cdot \left(Se_j^{\bullet}\right)$ Similarly $= \underbrace{e_{i}}_{i} \cdot \underbrace{(6)}_{i} \underbrace{e_{j}}_{j}$ equil juno summation 6; (e[•]...e[•].) $\begin{bmatrix} 3 \end{bmatrix}^{\bullet} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}^{\bullet}$ 5 Si Sj no summah S[°], = S, 6, -6, - J [S]= diagonal (G, 62, 63) ×3 Jet S= Si e & e = 2 6; 4 80 = di eise! Numerical example: The figure shows the eigenvalues and vectors of S,

5-10

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Provide the components of S in e1, e2 and e*1, e*2 systems

The figure shows the eigenvalues and vectors of S,

$d_{1} = t0$ $(-\frac{1}{12}, \frac{1}{12}) = d_{1}^{2}$ $d_{1} = d_{2}^{2}$ $d_{2} = (\frac{1}{12}, \frac{1}{12})$ Provide the components of S in e1, e2 and e*1, e*2 systems * <1 $=5\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\right]+\left(0\left[\frac{1}{2}\right]\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\right]\right]\right] +\left(0\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\right]\right]\right] +\left(0\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\right]\right]\right] +\left(0\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\right]\right]\right] +\left(0\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\right]\right] +\left(0\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\right]\right]\right] +\left(0\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\right]\right]\right] +\left(0\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\right]\right]\right] +\left(0\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\right]\right]\right]\right] +\left(0\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\right]\right]\right]\right] +\left(0\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\right]\right]\right]\right] +\left(0\left[\frac{1}{2}\left[\frac{1}{2}\right]\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\right]\right]\right]\right]\right] +\left(0\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\left[\frac{1}{2}\right]\right]\right]\right]\right] +\left(0\left[\frac{1}{2}\left[$ $= 5 \begin{bmatrix} \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} \end{bmatrix} + 10 \begin{bmatrix} \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & \frac{1}{2} \end{bmatrix} =$



Second example:

S = I & ei & ei det 9 = b162 63 11ace S= 01 +62+62 S = 5 = 4 400 17 bot 8 701

1.8.5



()nd.

$$L_{n}S = \sum_{i} L_{n}S_{i} e_{i} e_{i} e_{i} e_{i} \left[L_{n}S \right]^{*} = \begin{bmatrix} L_{n}S_{i} \\ L_{n}S_{j} \end{bmatrix} = \begin{bmatrix} L_{n}S_{i} \\ L_{n}S_{j} \end{bmatrix}$$



Skew-symmetric part of the tensor T does not contribute to u.Tu

Many times because of this, we directly work with symmetric tensors in forming such products

A positive definite tensor T satisfies the following properties:



A positive or semi-positive tensor only satisfies the first condition

Va u. Tu>0

Examples

$$T = \begin{bmatrix} 5 & 0 \\ 0 & 0 \end{bmatrix} \qquad \text{not positive definite} \\ \text{bd it's possitive} \\ \text{bd it's possitive} \\ \text{to i) } \begin{bmatrix} 5 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = G \qquad \text{bd} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \neq 0$$

$$T = \begin{bmatrix} 5 & 0 \\ 0 & 3 \end{bmatrix} \qquad 0.Tu = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} \begin{bmatrix} 5 \\ 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 0$$

$$\begin{array}{rcl}
\text{u.T}_{0} = 0 &\equiv 5v_{1}^{2} + 3v_{2}^{2} = 0 &\equiv \begin{bmatrix} v_{1} \\ v_{2} \end{bmatrix} \cdot \begin{bmatrix} 0 \\ 0 \end{bmatrix} \\
& \text{pertive definite} \\
\text{T} \in \begin{bmatrix} 5 & 0 \\ 0 & -3 \end{bmatrix} & \text{u.T}v = 5v_{1}^{2} - 3v_{2}^{2} \\
& \text{u.T}v = 5v_{1}^{2} - 3v_{2}^{2} \\
& \text{u.T}v = -3 < 0
\end{array}$$

penark 1 il Tis diagonalizable & symmetric

$$T = \bigcup_{i} \bigwedge_{i} \bigvee_{i} \bigvee_{i}$$

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Side note:

We can define a norm once we have a positive definite matric
$$\| U \|_{F} = \sqrt{U \cdot T U}$$

FYI Used in structural dynamics $M \dot{U} + [K \dot{U} = F$
 $M > 0$ mode shopes ϕ
 $\| f \|_{M} = \sqrt{\phi \cdot M \phi}$

Question:
Is there any easy may to form a symmetric

$$2$$
 positive (or positive def) tensor?
 $F \longrightarrow C = F F > 0$ if def $F \neq 0$
 rd other tensor $symmetric$
 $C^{+} - (F^{+}F)^{+} = F^{+}(F^{+})^{+} = FF - C$ so C is symmetric
positive check $(r, r, T_{b} = 1)$

pointie check

$$v.Cu = v.F^{\dagger}Fu=(F^{\dagger})^{\dagger}u.Fu=$$

 $Fu.Fv \ge 0$
 $v.v$
 $v.v$
 $v.v$
 $v.v$
 $v.v$
 $v.v$
 $v.v$
 $v.cu = 0$ $Fu.Fv=0 \longrightarrow Fv=0 \longrightarrow v=0$
 $(ddF_{v}o) \int f^{\dagger}u.v_{0}$
 $\int u.v_{0} \int f^{\dagger}u.v_{0}$



Theorem 112 (Polar Decomposition Theorem) Let $\mathbf{F} \in \text{Inv } \mathcal{V}$. Then \exists a unique pair of tensors $\mathbf{U}, \mathbf{V} \in \text{Psym}$ and a unique $\mathbf{R} \in \text{Orth } \mathcal{V} \ni$

 $\mathbf{F} = \mathbf{R}\mathbf{U} = \mathbf{V}\mathbf{R}.$

Moreover, det $\mathbf{R} = +1$ or det $\mathbf{R} = -1$, depending as det $\mathbf{F} > 0$ or < 0.

Motivation for the proof
if

$$F = RU \longrightarrow F^{t}F = (RU)^{t}RU = (U^{t}R^{t})RU$$

 $U^{t}(R^{t}R)U = U^{t}U = U^{2}$
 $I \qquad U^{t}Symm$
 $ReOrthV$
 $F^{t}F = U^{2}$
this gives the motivation on them to define U

$$C = F^{t}F \quad \text{this is symp_point}$$
then are can define its symp_point root
$$U:= \sqrt{C} = \sqrt{F^{t}}F \quad \equiv \quad U^{2} = F^{t}F$$
where $V = \sqrt{C} = \sqrt{F^{t}}F \quad \equiv \quad U^{2} = F^{t}F$
where $V = \sqrt{C} = \sqrt{F^{t}}F \quad \equiv \quad U^{2} = F^{t}F$
where $V = \sqrt{C} = \sqrt{F^{t}}F \quad \equiv \quad U^{2} = F^{t}F$
where $V = \sqrt{C} = \sqrt{F^{t}}F \quad \equiv \quad U^{2} = F^{t}F$
where $V = \sqrt{C} = \sqrt{F^{t}}F \quad = \sqrt{C} = \sqrt{C} = \sqrt{F^{t}}F \quad = \sqrt{C} = \sqrt{F^{t}}F \quad = \sqrt{C} = \sqrt{C}$

second para F=RU=VR if this holds V = RUR' = RURT Ri R^t=R⁻¹ is V pas sym. $\chi \cdot V_{\chi} = \chi \cdot RUR^{\dagger}\chi = R^{\dagger}\chi \cdot U(R^{\dagger}\chi) = y \cdot U_{\chi} \ge 0$ $\chi \cdot V_{\chi} = \chi \cdot RUR^{\dagger}\chi = \chi \cdot U(R^{\dagger}\chi) = y \cdot U_{\chi} \ge 0$ $V = \chi \cdot RUR^{\dagger}\chi = \chi \cdot U(R^{\dagger}\chi) = y \cdot U_{\chi} \ge 0$ $\frac{\nabla i}{\nabla^{t}} \frac{pos}{(RUR^{t})^{t}} = (R^{t})^{t} U^{t}R^{t} = RUR^{t} \cdot V$ FE RU = VR

$$V = R \cup R^{\dagger} \rightarrow V^{2} = (R \cup R^{\dagger})(R \cup R^{\dagger})$$

$$= R \cup (R^{\dagger}R) \cup R^{\dagger}$$

$$= R \cup^{2}R^{\dagger} = RF^{\dagger}FR^{\dagger}$$

$$F^{\dagger}F$$

$$Pecall$$

$$F = R \cup R = FU^{-1}$$

$$= F \cup^{-1} \cup^{-1}F^{\dagger} = FF^{\dagger}$$

$$= F \cup^{-1} \cup^{-1}F^{\dagger} = FF^{\dagger}$$

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F: RU = VR stretches

$$V = \sqrt{C}$$
 $C = F^{t}F$
 $V = \sqrt{B}$ $B = FF^{t}$

. .