CM2023/10/16 Monday, October 16, 2023 9:45 AM

Balana

Tensor fields: At each spatial location, we deal with a tensor

Examples of uses of tensor fields

Jaw



what are
$$\nabla_{2}$$
 ∇_{2} , ∇_{x} is $2D_{1}2D_{2}^{2}$
glad div cool
Heat equal: $q = -k\nabla I$
grad
curl $\hat{D} - \nabla xH + F = 0$ from Mexinell's eqns

Gradient m-th order tensor field. Tis an 1 = $\int_{i_1,\dots,i_m} \mathcal{C}_{i_1} \otimes \mathcal{C}_{i_2} \otimes \cdots \otimes \mathcal{C}_{i_m}$ Ma Ro where $T_{i_1, \dots, i_{m,j}} = \frac{\partial T_{i_1, \dots, i_{m,j}}}{\partial X_{i_j}}$ Ke to tern) or 1 eı 3CI for an orthograped excerdinate system, fixed in space

Since definition (1) is coordinate-dependent (using components of T in a given coordinate system) we need to demonstrate that it's a tensor $\chi_a = \sqrt[N]{2}$

Example, a 2nd order tensor: $= \frac{\partial \left(\frac{1}{12} \left(\chi_{1}, \chi_{2}, \chi_{3} \right) \right)}{\partial \chi_{3}} \quad in \quad (e_{1}, e_{2}, e_{3})$ $= \partial T_{mn} \left(\dot{x_1} , \dot{x_2} , \dot{x_3} \right) (in e_{1,e_{2}}, e_{3})$ ¥ X $(PT)'_{mnp} = \bigcup_{mi}$ 1 need nj not depending on x d Qmi Qui Tij (XuXerX) Qmi Qnj dTij (KuXerX) 9 Jun note = Qui Q $\frac{\partial}{\partial ij} (x_1, x_2, y)$ $Q_{m}; Q_{n};$ $= Q_{m_{r}} Q_{\eta_{i}}$ 0RK) Tij = <u>DQrk</u> 3ir $\times_{\mathcal{K}} = Q_{\mathcal{K}\mathcal{K}} \times_{\Gamma}' \rightarrow$ note QrK form coordinate Ŀ NK ้ว

We demonstrated that the gradient of a second order tensor is a 3rd order tensor (components transform as 3rd order tensor). Same applies to any other tensor order.

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Interpretation of the gradient operator:

Interpretation of the gradient operator: (Y, J, Uz) is the Cortesian coordinate system Guven: change of book from My to ytby ((y1, y2, y2) Question : $\sum \phi = \phi(y + \Delta y) - \phi(y)$ total increment 20 Dy $(y_1 + by_1, y_1 + by_2, y_3 + Ay_3) = (y_1, y_2, y_3) + (y_1 + y_2 + y_3) + (y_1 + y_2 + y_3) + (y_1 + y_2 + y_3) + (y_2 + y_3 + y_3) + (y_3 + y_3) + (y_3$ Et : Lust increment of t $+\left(\frac{1}{2}\frac{\partial \Phi}{\partial y_i \partial y_j}\Delta y_i \Delta y_i\right)$ S'to 2nd movement, $= S \varphi + \delta^2 \varphi + \delta^3 \varphi - \dots = \begin{bmatrix} \partial \varphi & \partial \varphi & \partial \varphi \\ \partial \varphi & \partial \chi & \partial \chi & \partial \chi \end{bmatrix}$ L\$= \$14+by)-\$(4) L total veral: Hessian δф $+ \frac{1}{2} \begin{bmatrix} Dy_{1} & Dy_{2} & Dy_{3} \end{bmatrix} \begin{pmatrix} \phi_{11} & \phi_{12} & \phi_{13} \\ \phi_{11} & \phi_{322} & \phi_{13} \\ \phi_{31} & \phi_{32} & \phi_{133} \end{pmatrix} \begin{bmatrix} by_{1} \\ Dy_{2} \\ Dy_{3} \\ D$ + H. O.T higher order $\phi_{ij} = \frac{\partial \phi}{\partial y_i \partial y_i}$ terms $\Delta \phi = \delta \phi * \delta' \phi - - \pi \delta \phi$ Δy ->0 as 50 alon Jector



We can do a Taylor's expansion of a vector and observe:

$$V = V_{1} \left(\frac{1}{2} \sqrt{\frac{1}{2}} \sqrt{\frac{1}{2}} \sqrt{\frac{1}{2}} + \frac{1}{2} \sqrt{\frac{2}{3}} \right)$$

$$M = \left(\frac{1}{2} \sqrt{\frac{1}{2}} \sqrt{\frac{1}{2}} + \frac{1}{2} \sqrt{\frac{2}{3}} \right)$$

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$$M = \left(\frac{1}{2} \sqrt{\frac{1}{2}} \sqrt{\frac{1}{2}} \sqrt{\frac{1}{2}} \sqrt{\frac{1}{2}} + \frac{1}{2} \sqrt{\frac{1}{2}} \sqrt{\frac{1}{2}} \sqrt{\frac{1}{2}} \sqrt{\frac{1}{2}} \right)$$

$$M = \left(\frac{1}{2} \sqrt{\frac{1}{2}} \sqrt{\frac{1}$$

Orthonormal curvilinear coordinate systems:

$$\begin{cases} \mathcal{Y}_{1} = \mathcal{Y}_{1}(X_{1}, X_{2}) & \text{def of} \\ \mathcal{Y}_{2} = \mathcal{Y}_{2}(X_{1}, X_{2}) & \text{ourvilinear coordinate} \\ \text{system 8} \\ \text{relation between $\mathbf{x}_{2}\mathbf{y}$ & \mathbf{y}_{1} = \mathbf{0} \end{cases}$$

J' V

P

Example: Polor coordinate system:
$$X_1 = Y$$

 $\begin{cases} y_1 = x_1 \ G \ x_2 \\ y_2 = x_1 \ Sin x_2 \end{cases}$

How do we calculate the gradient in the polar coordinate system?

$$\vec{F} = f \vec{e}_{T}$$

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$$\vec{F} = f \vec{e}_{T}$$

$$\vec{F} = d (re_{T}) = (dr)e_{T} + T(de_{T})$$

$$\vec{F} = d (re_{T}) = (dr)e_{T}$$

$$\vec{F} = d (re_{T})e_{T}$$

$$\vec{F} = d (re_$$

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$$(ar) (cr) = r (arr)$$

$$= (dr)e_{1} + r (d\theta e_{0})$$

$$= (dr)e_{1} + r (d\theta e_{0})$$

$$= (dr)e_{1} + r (d\theta e_{0}) e_{0}$$

$$= (dr)e_{1} + r (d\theta e_{0}) e_{0}$$

$$= (dr)e_{1} + r (d\theta e_{0}) e_{0}$$

$$= (dr)e_{1} + (r d\theta)e_{0}$$

$$= (dr)e_{1}$$

Demonstration of gradient equation for polar coordinate, for a vector

Demonstration of gradient equation for polar coordinate, for a vector

$$V = V_{\Gamma} \ell_{\Gamma} + V_{0} \ell_{0} \implies dv = d(v_{0}e_{\Gamma}) + d(V_{0}e_{0}) = (dv_{\Gamma})e_{\Gamma} + V_{\Gamma}(de_{\Gamma}) + (W_{0})e_{0} + V_{0}(de_{0}) = (v_{\Gamma}r_{0}d_{\Gamma} + v_{\Gamma,0}d_{0})e_{\Gamma} + v_{\Gamma}(e_{0}d_{0}) + (V_{0,r}d_{\Gamma} + V_{0,0}d_{0})e_{0} + V_{0}(-e_{\Gamma}d_{0}) = (V_{\Gamma}r_{0}d_{\Gamma} + (V_{0,r}d_{\Gamma} + (V_{0,r}d_{\Gamma} + (V_{0,r}d_{\Gamma} + (V_{0,r}d_{\Gamma}))e_{0})e_{0}, \text{ Node } dF_{r} = dr$$

$$= V_{0}e_{\Gamma} + (V_{0}e_{\Gamma} + (d_{0})e_{0})e_{\Gamma} + (V_{0,r}d_{\Gamma} + (V_{0,r}d_{\Gamma} + (V_{0,r}d_{\Gamma}))e_{0})e_{0} + (V_{0,r}d_{\Gamma})e_{0} + (V_{0,r}d_{\Gamma})e_{0})e_{0}$$

$$= V_{r}r_{r} + (d_{0}e_{0})e_{0} = (V_{r}r_{r} + (V_{0,r}d_{\Gamma})e_{0})e_{0} + (V_{0,r}d_{\Gamma})e_{0} + (V_{0,r}d_{\Gamma})e_{0})e_{0}$$

$$= (dv_{1}e_{0})e_{0} = (V_{1}r_{r} + (V_{0,r}d_{\Gamma})e_{0})e_{0} + (V_{0,r}d_{\Gamma})e_{0})e_{0}$$

$$= (dv_{1}e_{0})e_{0} = (V_{1}r_{r} + (V_{0,r}d_{\Gamma})e_{0})e_{0} + (V_{0,r}d_{\Gamma})e_{0} + (V_{0,r}d_{\Gamma})e_{0})e_{0}$$

$$= (dv_{1}e_{0})e_{0} = (V_{1}r_{r} + (V_{1,0}e_{0})e_{0})e_{0}$$

$$= (dv_{1}e_{0})e_{0} = (V_{1}r_{r} + (V_{1,0}e_{0})e_{0})e_{0} + (V_{1}e_{0})e_{0} + (V_{1}e_{0})e_{0}$$

How de ale adapted divergence
$$2$$
 curl : T is might other
 $V_{0}T = T_{1} - C_{m} \int_{0}^{10} e_{1} \otimes \cdots \otimes e_{1} m_{-1}$
main order
last two entries are centracted
 $M = I$ $V_{0}T = T_{0}i = trace (VT)$
Tolken obtined for vector
 $avril V = ax (portT)$
Examples in Corkesian cordinate $\int_{1}^{10} \frac{1}{2} \frac{1}{2}$

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$$\begin{aligned}
 & U_{3,1} & V_{3,2} & Y_{3,3} \\
 & \nabla_{0}V &= t(0 CQ (\nabla V) &= V_{0,1} + V_{2,2} + V_{3,3} = V_{1,1} \\
 & avrl V &= avrl s(kew \nabla V) &= \nabla X V = brl \\
 & J & J \\
 & J &$$

Martinec_Zdenek_Charles_U_Prague_Martinec-ContinuumMechanics.pdf Please read appendix C