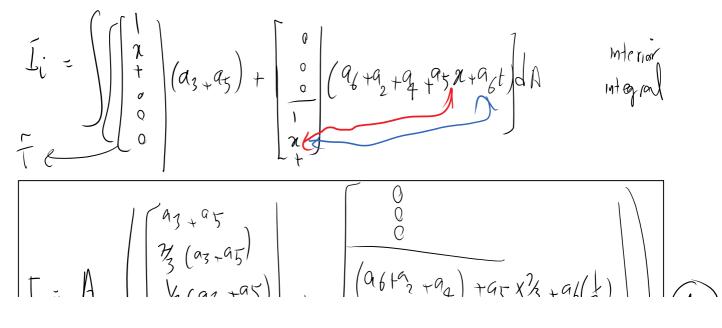
## 2018/03/26

Monday, March 26, 2018 11:38 AM

$$\int \left\{ f\left(a_{3} + a_{5}\right) + \hat{q}\left(a_{6} + a_{2} + a_{4} + a_{5} + a_{6} + a_{6$$



Inflow face:  

$$I_{inf} = \int T[(a_{1} + a_{2} \times + a_{3} \star) + \hat{q}(a_{4} + a_{3} \times + a_{3} \star)] dx$$

$$Je_{i}$$

$$t = 0$$

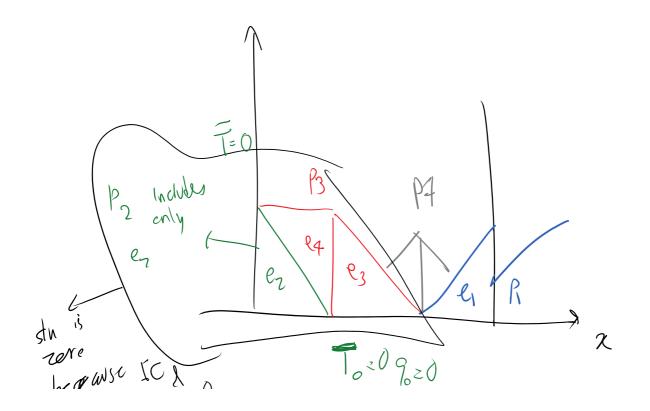
$$f_{or} \quad this simple \quad cose \quad ue \quad directly \quad ntegrate \quad it \quad n \quad \chi \quad hol$$

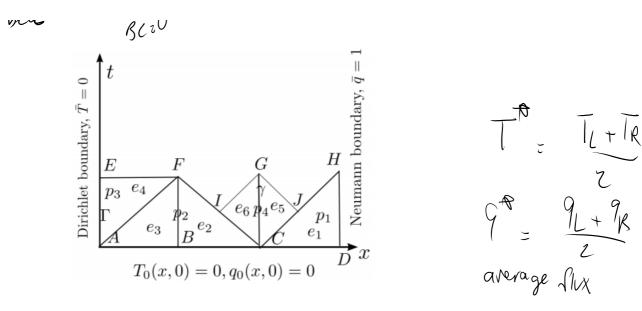
$$f_{or} \quad slant \quad has \quad it's \quad eosier \quad to \quad cse \quad the \quad formula$$

$$\int f_{a} \star \hat{f}_{b} ds \qquad A_{1} \qquad f_{a} \qquad A^{\dagger} \qquad f_{a} \qquad A^{\dagger} \qquad f_{a} \qquad A^{\dagger} \qquad f_{a} \qquad f_{b} \qquad g_{a} \qquad g_{a}$$

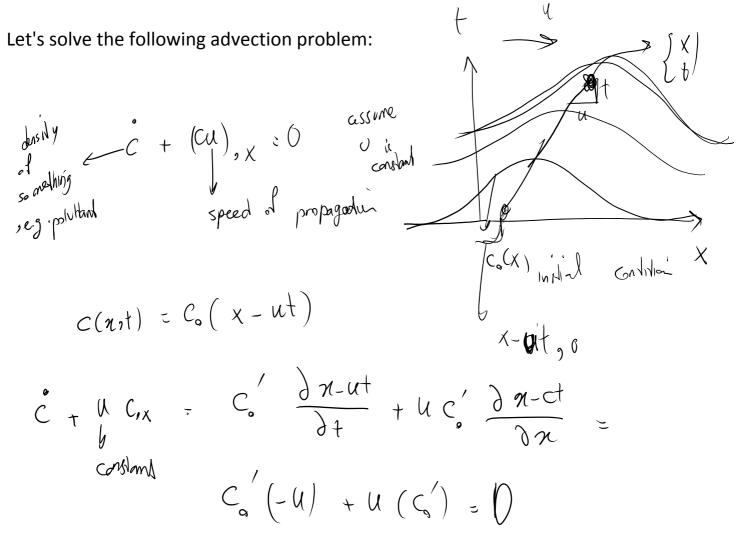
$$f_{a} \text{ this simple rose}$$

$$I_{a} \int \left( \frac{1}{2} + \frac{1}{2} \frac{1}{2} \frac{1}{2} + \frac{1}{2} \frac{1}{2} \frac{1}{2} \frac{1}{2} + \frac{1}{2} \frac{$$

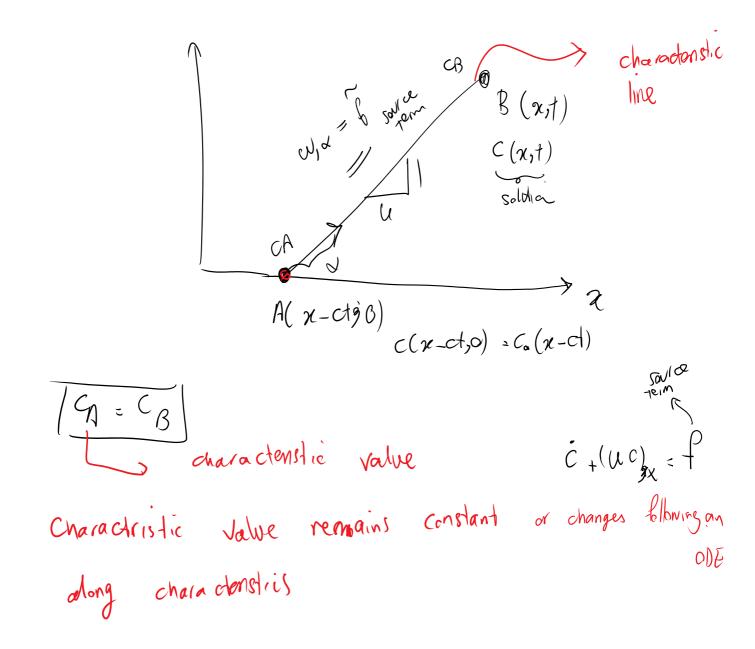




Riemann solutions:



 $\mathbf{a}$ 



The use of characteristics for hyperbolic PDEs:

Along characteristics PDE turns to an ODE which is easy to solve. For no source term case, characteristics in fact remain constant.

How about solid mechanics in 1D:

 $P - \nabla S = \rho b$ 

Quadian of malin

$$\dot{p} - \nabla \cdot d = pb \qquad \text{apodian of makin}$$

$$D \qquad \underline{p} - G_{2X} : pb \qquad \text{two untarrant}$$

$$P_{2} \qquad P_{2} \qquad P_{3} \qquad P_{3}$$

$$I = \begin{bmatrix} 0 \\ p_{1} & p_{1} \end{bmatrix} \begin{bmatrix} -\frac{e}{p} & 0 \end{bmatrix} = \begin{bmatrix} -\frac{e}{p} & 0 \end{bmatrix} = \begin{bmatrix} -\frac{e}{p} & 0 \end{bmatrix}$$

$$spatial flux and dix \qquad source term
$$\begin{array}{c}
q \\
+ & \left(\frac{1}{p_{x}}\right) = S \\
Right way all ariting a liver Garterization line is \\
q \\
+ & \left(\frac{1}{p_{x}}\right)_{2x} + \left(\frac{A_{y}q}{p}\right)_{y} \\
- & \left(\frac{1}{p_{x}}\right)_{2x} + \left(\frac{A_{y}q}{p}\right)_{y} \\
- & \left(\frac{1}{p_{x}}\right)_{x} = S \\
q \\
+ & \left(\frac{1}{p_{x}}\right)_{2x} + \left(\frac{A_{y}q}{p}\right)_{y} \\
- & \left(\frac{1}{p_{x}}\right)_{x} = S \\
- & \left(\frac{1}{p_{x}}\right)_{x} = S$$$$

$$(Uq) + UAU Uq = US$$

$$(U + (UAU) hy = S = new source term Soir US$$

$$(Uht if this is a dugonal matrix)$$

$$UAU^{-1} = D \implies UA = DU$$

$$(IAU^{-1} = DU$$

$$(I$$

$$J_{V} = \frac{1}{2} \int_{V_{0}}^{V_{0}} LEFT eigenvector \# i \text{ of } A \\D_{ii} = eigenvector \# i \text{ of } A \end{bmatrix}$$

$$Note d Right eigen pairs one have
$$A(v_{i})_{n\chi i} = \frac{1}{2} \int_{V_{0}}^{V_{0}} \int_{V_{0}}^{V_{0}} \int_{U_{0}}^{V_{0}} \int_{U_{$$$$

$$(V_i, D_i)$$
 let a genve dor, eigenvalue on  $A \iff$   
 $(V_i, D_i)$  right " --- At

The solution process:  

$$g' + A q_{x} : S$$
  $VA - DV$ , define  $W = Uq$   
 $modinix of$   
 $loft organized organized of  $W = Uq$   
 $modinix of$   
 $loft organized of  $W = Ug$   
 $W + D W_{x} = Sor$   $S_{u} = US$   
 $W + D W_{x} = Sor$   $S_{u} = US$   
 $Solve the uncaupled simple ODES for  $\omega, \dots, \omega_{n}$   
 $W_{i}(x,t)$   
 $\overline{(3)} q = U^{+}\omega \rightarrow q(x,t)$  will be  
known$$$ 

We do this solution process through an example:

$$-c^{2}u^{2} = -cu^{2}$$

$$u^{2} = -cu^{2}$$

$$\left( \begin{array}{ccc} \lambda = -C & U = \begin{bmatrix} C & I \end{bmatrix} \right)$$
  
Similarly for  $\lambda = C$   
$$\left( \begin{array}{ccc} \lambda = C & U = \begin{bmatrix} -C & I \end{bmatrix} \right)$$

So 
$$\int = \left[ \begin{array}{c} U_{1} \\ U_{2} \end{array} \right] \cdot \left[ \begin{array}{c} C \\ -C \end{array} \right] D \cdot \left[ \begin{array}{c} A_{1} \\ 0 \\ A_{2} \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ 0 \\ -C \end{array} \right] D \cdot \left[ \begin{array}{c} A_{1} \\ 0 \\ A_{2} \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ 0 \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C \\ -C \end{array} \right] \cdot \left[ \begin{array}{c} -C$$