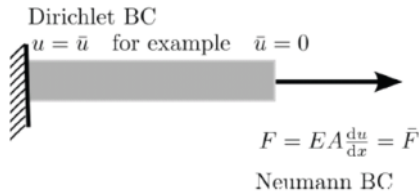


Bar problem



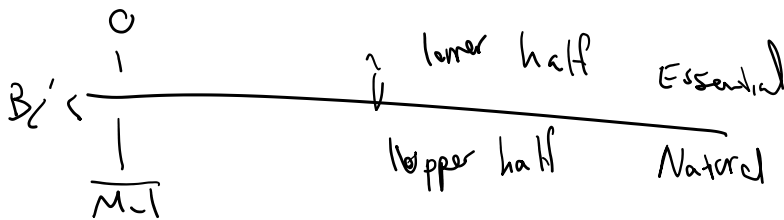
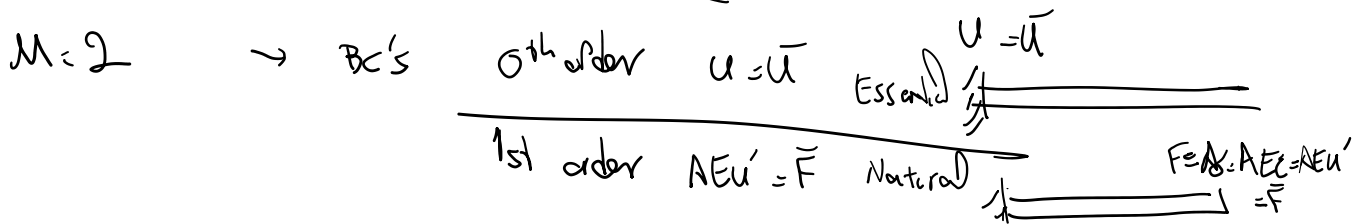
Operator	Sample	1D bar	operator order
$L_{2m}(u) = r$	$\frac{d}{dx} (EA \frac{du}{dx}) = -q$	$L_{2m} = \frac{d}{dx} (EA \frac{d(\cdot)}{dx})$	$m = 1 (M = 2)$
$L_u(u) = \bar{u}$	$u = \bar{u}$	$L_u = (\cdot)$	$M_u = 0$
$L_f(u) = \bar{f}$	$EA \frac{du}{dx} = \bar{f}$	$L_f = EA \frac{d(\cdot)}{dx}$	$M_f = 1$

$(EAU')' - q = 0$

order of DE $\leftarrow M = 2$

$m = \frac{M}{2} = 1$

BC's are order 0 to $M-1$



Slides 32 to 34 provide the formulation of the beam problem. I'll cover it later, but it's good to read it and see how the essential and natural BCs are divided

WRS and Weak statement:

Weighted Residual Statement (WRS)

$R_i = (EAU')' + q$ in \mathcal{D}

$R_u = \bar{u} - u$ essential BC residual on \mathcal{D}_u

$R_f = \bar{F} - F = \bar{F} - EAU'$ on \mathcal{D}_f

DE: $(EAU')' + q = 0$

$F = EAU' = \bar{F}$

Multiply by weights

$\int_{\mathcal{D}} w_i R_i dV$

some function of w eg w'

$$\int_D w R_i dV + \int_{\partial\Omega} f(w) R_u dS + \int_{\partial\Omega} w R_p dS = 0$$

weight funcn

$R_u = \bar{u} - u$

$\frac{\partial u}{\partial n} \Big|_{x=0}$ 1D case, this problem $f(w) R_u \Big|_{x=0}$

$\frac{\partial p}{\partial n} \Big|_{x=L}$ 1D, this problem $w R_p \Big|_{x=L} = 0$

In continuous FEM (this course), I'll show that we don't have the middle term (integral of residual on essential BC)

$R_u = \bar{u} - u$

$R_i = (EAu')' + q$

$R_p = \bar{F} - EAu$

$$\int_0^L w R_i dx + w R_p \Big|_{x=L}$$

$$= \int_0^L w ((EAu')' + q) dx + w (\bar{F} - EAu) \Big|_{x=L} = 0$$

0th order 2nd order

$$\int_0^L UV' dx = \int_0^L (U'V - UV') dx$$

$$= UV \Big|_0^L - \int_0^L U'V dx$$

$$\int_0^L w (EAu')' dx$$

$U \quad V \rightarrow V = EAu'$

$$= w EAu' \Big|_0^L - \int_0^L (w' EAu') dx$$

$$\int_0^L w (EAu')' dx = w EAu' \Big|_{x=L} - w EAu' \Big|_{x=0} - \int_0^L w' EAu' dx$$

$$\int_0^L w(EAu') dx + \int_0^L wq dx + w(\bar{F} - EAu)|_{x=L} = 0$$

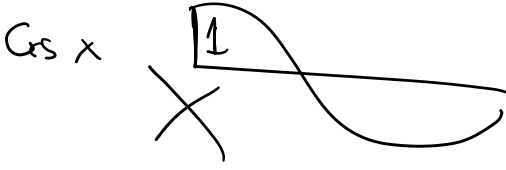
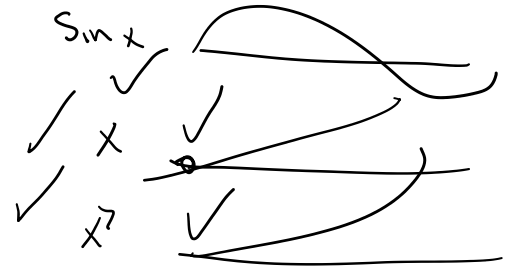
1st order

$$wEAu|_{x=L} - wEAu|_{x=0} - \int_0^L wEAu' dx + \int_0^L wq dx + w\bar{F}|_{x=L} - wEAu|_{x=L} = 0$$

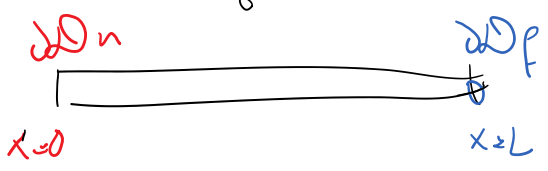


When we go from WRS to weak statement we only limit ourselves to weight functions that are zero on essential BC.

good for WRS
bad for weak statement



~~$$-w(x=0)EAu(x=0) - \int_0^L w'EAu' dx + \int_0^L wq dx + w(L)\bar{F}|_{x=L} = 0$$~~



$$\int_0^L wEAu' dx = \int_0^L wq dx + w\bar{F}|_{x=L}$$

Weak statement

$$\int_0^L w \left(\frac{d}{dx} (EAu') \right) dx + w(\bar{F} - EAu')|_{x=L} = 0 \quad \text{WRS}$$

2 derivatives

WRS

$$\Omega = [0, L]$$



Find $u(x) \in \mathcal{V} = \left\{ f \in C^2(\Omega) \mid \begin{array}{l} \forall x \in \Omega \quad u(x) = \bar{u} \\ \text{where } \bar{u} \text{ is the } \text{AD problem} \\ u(0) = \bar{u} \end{array} \right\}$

space of solutions

such that for all $w(x) \in \mathcal{W} = \left\{ f \in C^1(\Omega) \mid \text{---} \right\}$

we have

$$\int_0^L w \left(\frac{d}{dx} (EAu') + q \right) dx + w(\bar{F} - EAu')|_{x=L} = 0$$

2 derivatives



Because we did not add weight time residual essential boundary to the WRS

~~WRS~~

WRS

Weak statement

Find $u(x) \in \mathcal{V} = \left\{ f \in C^1(\Omega) \mid f(0) = \bar{u} \right\}$

such that for all weight functions $w \in \mathcal{W} = \left\{ f \in C^1(\Omega) \mid f(0) = 0 \right\}$

$$\int_0^L w' EAu' dx = \int_0^L w q dx + w \bar{F} |_{x=L}$$

Key points for the weak statement:

- Derivative orders are balanced for weight and solution
- Both w and u satisfy the essential BC.
 - o For solution the actual essential BC
 - o For the weight, the homogenous (e.g. 0) version of that.

2D examples:

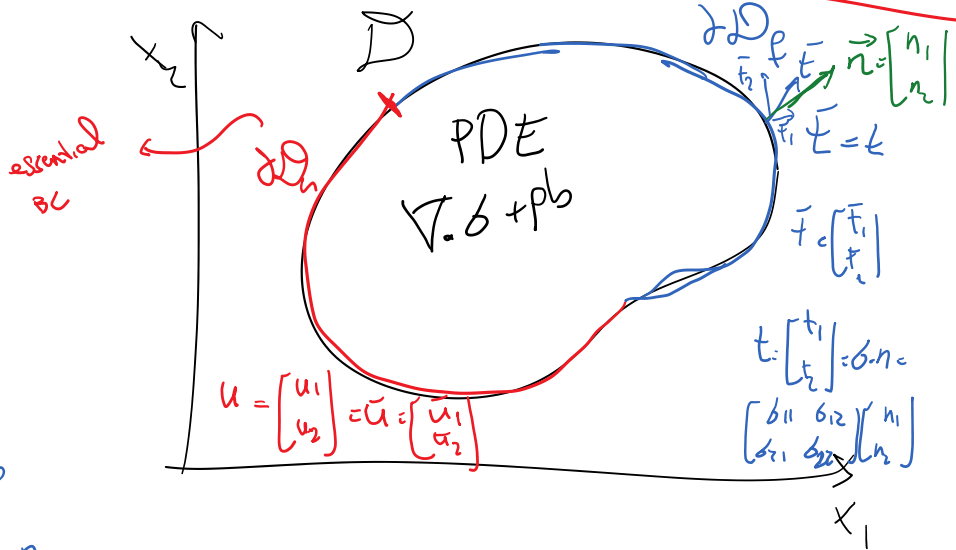
1. Elastostatics (in class)
2. Heat conduction (HW2)

$$R_i: \forall x \in \Omega$$

$$R_i = \nabla \cdot \sigma + pb$$

$$R_f: \forall x \in \partial \Omega_f$$

$$R_f = \bar{t} - \sigma \cdot n$$



$$R_i: \forall x \in \partial \Omega_u \quad R_i = \bar{u} - u$$

Weighted residual statement

Find $u \in \mathcal{V} = \{ f \in C^2(\Omega) \mid \forall x \in \partial \Omega_u \quad u(x) = \bar{u} \}$ we strongly satisfy essential BC

such that $\forall w \in \mathcal{W} = \{ f \in C^0(\bar{\Omega}) \mid \}$

$$\int_{\Omega} w R_i \, dV + \int_{\partial \Omega_f} w R_f \, dS = 0$$

$\int_{\Omega} w R_i \, dV$ \downarrow $\nabla \cdot \sigma + pb$
2 der of u

~~$\int_{\partial \Omega_u} w R_i \, dS = 0$~~

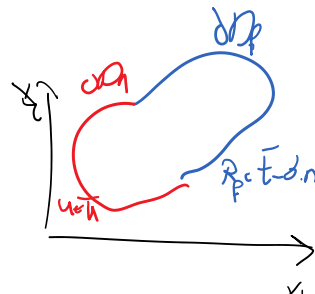
$\partial \Omega_u$ same function of weight

\downarrow we cross this term as we satisfy it strongly

The process of deriving the weak statement

$$\int_{\Omega} w (\nabla \cdot \sigma + pb) \, dV + \int_{\partial \Omega_f} w (\bar{t} - \sigma \cdot n) \, dS = 0$$

$\int_{\Omega} w (\nabla \cdot \sigma + pb) \, dV$ \downarrow $(\sigma u)''$



$\delta \mathcal{D}$
 similar to (1D) $\int \omega (\delta u)' + \rho$ + $\omega (\bar{F} - EAu')$

$\bar{F} - F = \bar{F} - EAu'$

$\int \omega (\nabla \cdot \sigma) dV$

\mathcal{D} 1/2 derivative for u

2D version of IBP

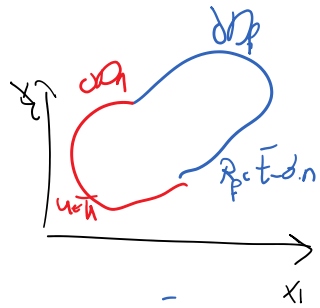
$\int_{\mathcal{D}} \omega (\nabla \cdot \sigma) dV = \int_{\partial \mathcal{D}} \omega \sigma \cdot n dS - \int_{\mathcal{D}} \nabla \omega \cdot \sigma dV$

WRS $\int_{\mathcal{D}} \omega (\nabla \cdot \sigma) dV + \int_{\mathcal{D}} \omega (pb) dV + \int_{\partial \mathcal{D}_f} \omega (\bar{F} - \sigma \cdot n) dS$

$\left(\int_{\partial \mathcal{D}} \omega (\sigma \cdot n) dS - \int_{\mathcal{D}} \nabla \omega \cdot \sigma dV \right) + \int_{\mathcal{D}} \omega pb dV$

$+ \int_{\partial \mathcal{D}_f} \omega \bar{F} dS - \int_{\partial \mathcal{D}_f} \omega \sigma \cdot n dS = 0$

$\int_{\partial \mathcal{D}_n} \omega (\sigma \cdot n) dS - \int_{\partial \mathcal{D}_f} \omega \sigma \cdot n dS$



$= \int_{\partial \mathcal{D}_n} \omega \sigma \cdot n dS$

$\omega = 0$

$\text{on } \partial \mathcal{D}_n$