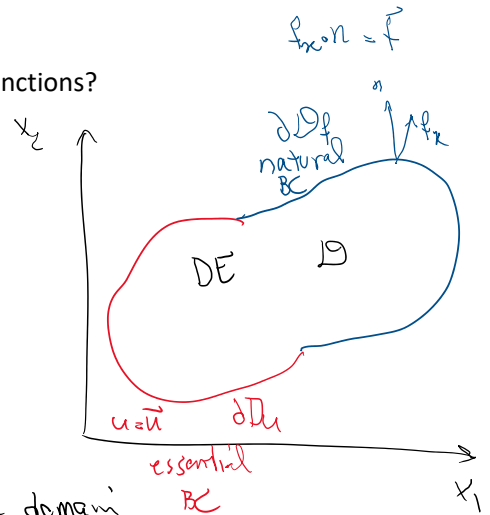


Why Least Square method is a Weighted Residual method with specific weight functions?

$$R_i(u) = L_M(u) - r \quad \text{source term}$$

error in satisfying the differential equation

$$R_f(u) = \bar{F} - f_{x \cdot n} = \bar{F} - L_f(u)$$



L_M & L_f are differential operators for inside the domain & natural boundary

Examples

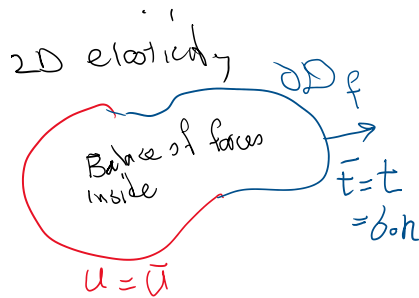


$$R_i = (EAU')' + q = L_M(u) - r$$

$$L_M = (EA(x)')', r = -q$$

$$R_f = \bar{F} - F|_L = \bar{F} - (EAU')|_{x=L}$$

$$L_f = EA(x)'$$



$$R_i(u) = \nabla \cdot \sigma + pb = \nabla \cdot (C\varepsilon) + pb = \nabla \cdot C \left(\frac{\nabla u + \nabla u^t}{2} \right) + pb = L_M(u) - r$$

$r = -pb$

$$L_M = \nabla \cdot C \left(\frac{\nabla + \nabla^t}{2} \right)$$

$$R_f = \bar{T} - t = \bar{T} - b \cdot n = \bar{T} - (C\varepsilon) \cdot n = \bar{T} - C \left(\frac{\nabla u + \nabla u^t}{2} \right) \cdot n$$

$$L_f = C \left(\frac{\nabla + \nabla^t}{2} \right)$$

2D heat conduction

$$R_i(T) = \nabla \cdot q - \bar{Q} = \nabla \cdot (-k\nabla T) - \bar{Q}$$

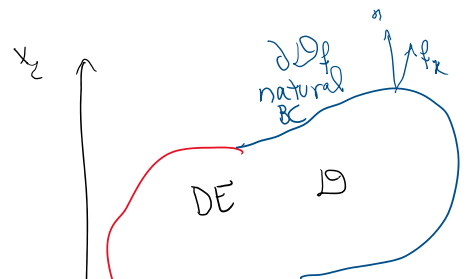
$$L_M = \nabla \cdot (-k\nabla)$$

$$R_f = \bar{q} - q \cdot n = \bar{q} + (k\nabla T) \cdot n$$

$$L_f = -k\nabla(\cdot)$$

$$R_i(u) = L_M(u) - r \quad \text{source term}$$

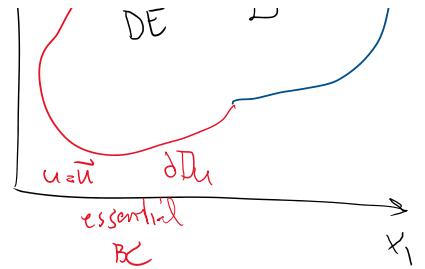
error in satisfying the differential equation



is satisfying the differential equation

(1)

$$R_p(u) = \bar{f} - f_{k02} = \bar{f} - \mathcal{L}_p(u)$$



$$R^2 = \int_D R_i^2(u) dV + \int_{\partial D_p} R_p^2(u) dS \quad (2)$$

$$u = \phi_p + \sum_{i=1}^n a_i \phi_i$$

discretization

(3)

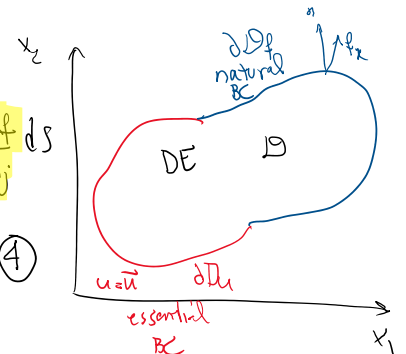
we have n unknowns a_1, \dots, a_n

$R^2(a_1, \dots, a_n)$, we want to minimize $R^2(a_1, \dots, a_n)$

$$\nabla R^2 = \begin{pmatrix} \frac{\partial R^2}{\partial a_1} \\ \vdots \\ \frac{\partial R^2}{\partial a_n} \end{pmatrix} = \begin{pmatrix} 0 \\ \vdots \\ 0 \end{pmatrix}$$

$$\frac{\partial R^2}{\partial a_j} = \frac{\partial}{\partial a_j} \int_D R_i^2(a_1, \dots, a_n) dV + \int_{\partial D_p} R_p^2(a_1, \dots, a_n) dS$$

$$= \int_D \frac{\partial R_i^2}{\partial a_j} dV + \int_{\partial D_p} \frac{\partial R_p^2}{\partial a_j} dS = \int_D 2R_i \frac{\partial R_i}{\partial a_j} dV + \int_{\partial D_p} 2R_p \frac{\partial R_p}{\partial a_j} dS$$



we need to plug in values of $R_i, \frac{\partial R_i}{\partial a_j}, R_p$ & $\frac{\partial R_p}{\partial a_j}$ in (4)

(4)

$$(3) : u = \phi_p + \sum_k a_k \phi_k$$

$$\Rightarrow R_i(a_1, \dots, a_n) = R_i(u) = \mathcal{L}_M \left(\phi_p + \sum_{k=1}^n a_k \phi_k \right) - \bar{f} = \mathcal{L}_M(\phi_p) + \sum_{k=1}^n a_k \mathcal{L}_M(\phi_k) - \bar{f}$$

Example 1D bar $\mathcal{L}_M = (EA(x))'$

$$(EA(\phi_p + \sum_{k=1}^n a_k \phi_k))' = (EA\phi_p)' + (EA\phi_k)'$$

(because \mathcal{L}_M is linear we can open it $\mathcal{L}_M(a_1 f_1 + a_2 f_2) = a_1 \mathcal{L}_M(f_1) + a_2 \mathcal{L}_M(f_2)$)

(because L_M is linear we can open it $L_M(a_1 f_1 + a_2 f_2) = a_1 L_M(f_1) + a_2 L_M(f_2)$)
 $\Rightarrow R_i = L_M(\Phi) + a_1 L_M(\Phi) + \dots + a_n L_M(\Phi_n)$ or (5 i)

$$\frac{\partial R_i}{\partial a_j} = L_M(\Phi_j)$$

(5 ii)

Similarly $R_f(u) = R_f(a_1, \dots, a_n) = \bar{F} - L_f(\Phi_p + \sum_k a_k \Phi_k)$ or

$$= \bar{F} - L_f(\Phi) - \sum_k a_k L_f(\Phi_k)$$

(5 iii) L_f is linear

$$\frac{\partial R_f(u)}{\partial a_k} = -L_f(\Phi_k)$$

(5 iv)

(4) $\int_D \frac{\partial R_i}{\partial a_j} dV + \int_{\partial D_f} \frac{\partial R_f}{\partial a_j} dS = \int_D 2R_i \frac{\partial R_i}{\partial a_j} dV + \int_{\partial D_f} 2R_f \frac{\partial R_f}{\partial a_j} dS$

(5 ii) $w_j = L_M(\Phi_j)$ (5 iii) $(a_f)_j = -L_f(\Phi_j)$

R2:

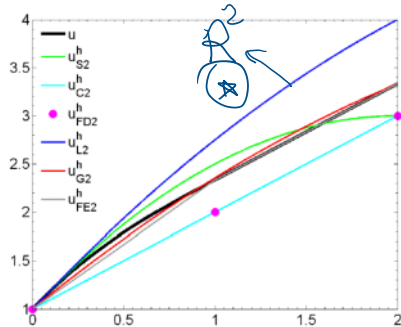
+ It can be added to a formulation to satisfy certain condition weakly, e.g. interpenetration.

+ Always get symmetric matrices

- Minimizing the residual does not mean we have the best solution (u_h closest to the exact solution)

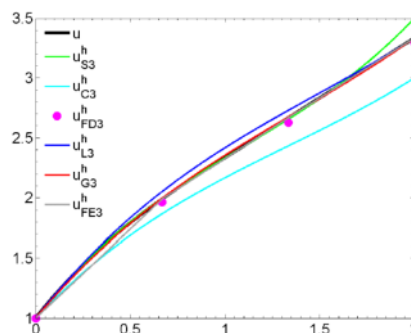
- Both solution and the weights get very high derivative orders (M)

Bar example, $n = 2$, Comparison of solutions



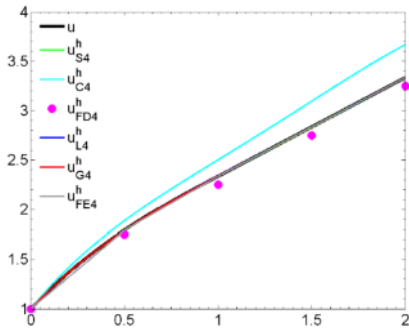
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Bar example, $n = 3$, Comparison of solutions



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Bar example, $n = 4$, Comparison of solutions



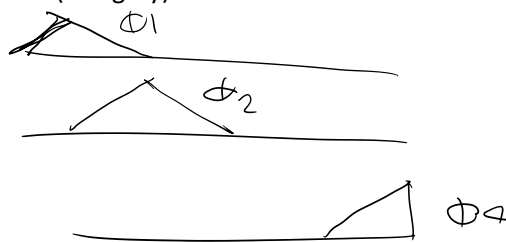
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⊛ R^2 is not the closest to the exact solution $(u^h - u^{exact})$ because we minimize the error in

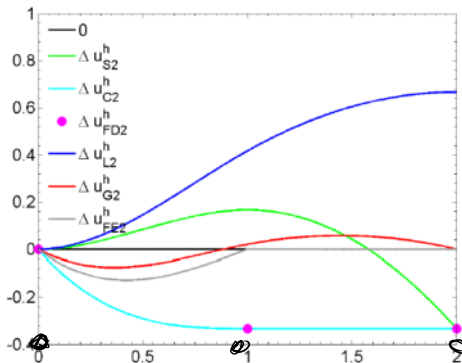
differential equation not w.r.t. the exact solution.

We see that the Galerkin methods (Spectral - red) and (FE - gray) have the best solutions.

u, u, u

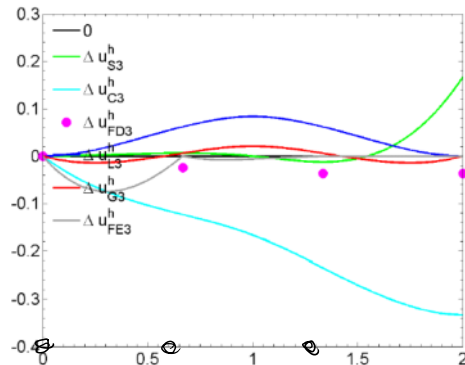


Bar example, $n = 2$, Comparison of solutions



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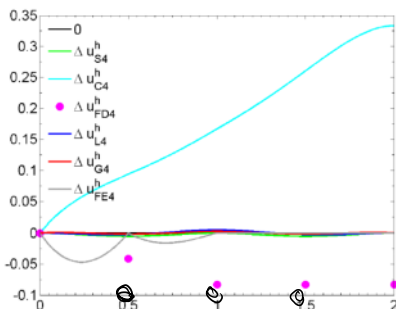
Bar example, $n = 3$, Comparison of solutions



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For FEM the errors are equal to zero ⊛

Bar example, $n = 4$, Comparison of solutions



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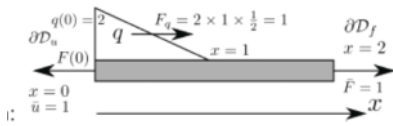
Why we have larger errors on the left side?

Unfortunately, this property is only for 1D and in 2D and 3D we don't have this very nice property

The exact solution can be summarized as,

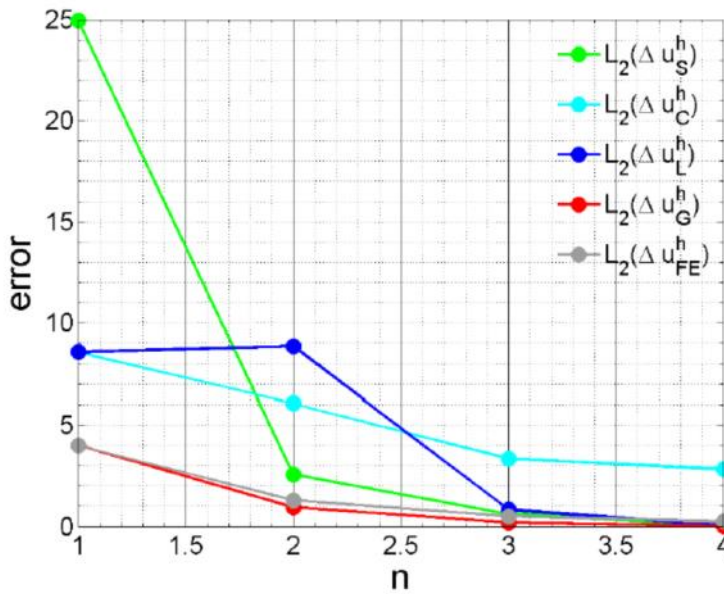
$$u(x) = \begin{cases} \frac{x^3}{3} - x^2 + 2x + 1 & 0 \leq x \leq 1 \\ x + \frac{4}{3} & 1 < x \leq 2 \end{cases} \quad (179)$$

Because of the source term we have a higher order solution there and harder to capture that numerically



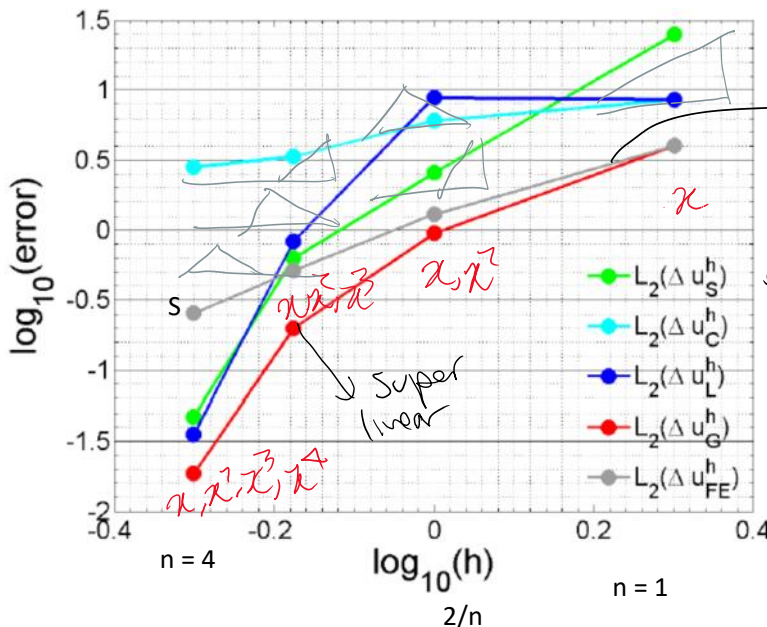
Error convergence:

Bar example, Error Convergence



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Bar example, Error Convergence



linear
 $(\log(\text{error}) = C + \log(h))$
 $\text{error} = C h^\alpha$
 $\text{error} = C h^{f(p)} = 2-p+1$

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n = 4

$$\log_{10}(h) \\ 2/n$$

n = 1

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