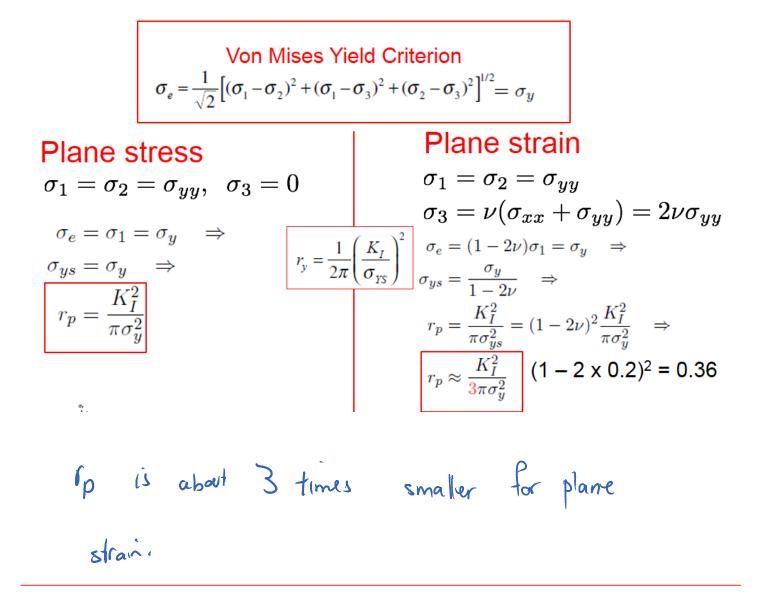
2016/10/11 Tuesday, October 11, 2016 8:39 AM

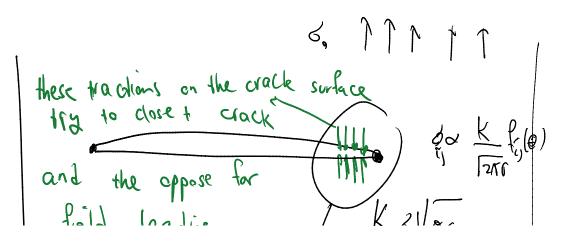
> From last time $K = \text{Jominant radius} r_s \propto \left(\frac{k}{\delta}\right)^k$ applied stress plostic/nonlinear zone size K α P applied loved l yreld stress

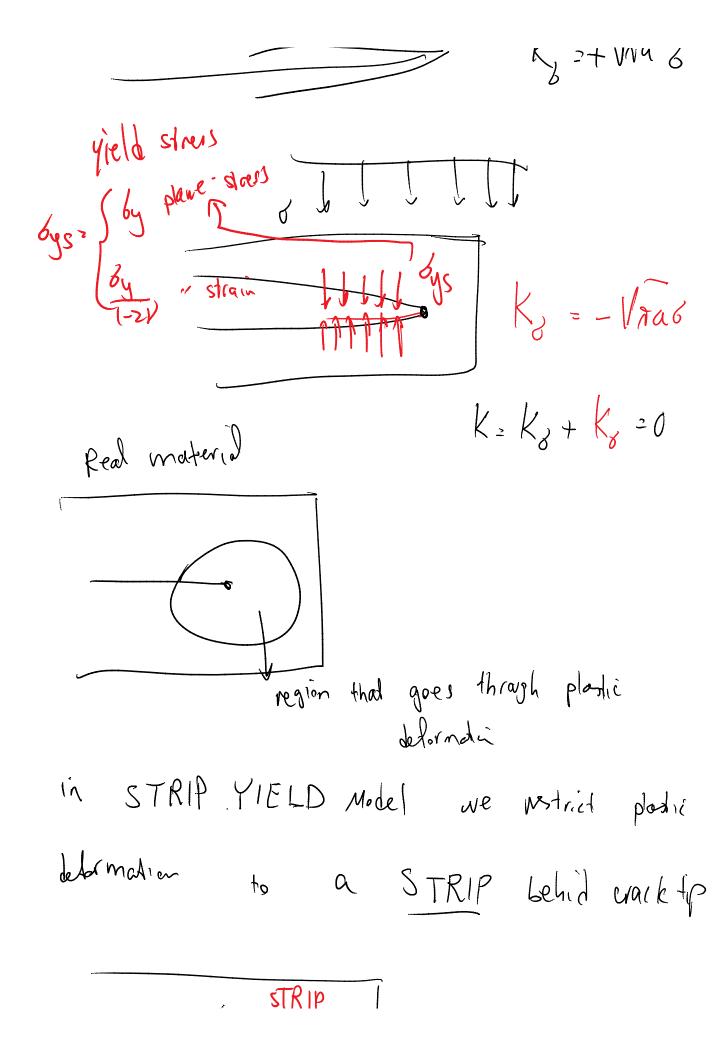
From last time

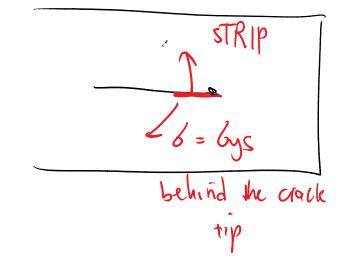


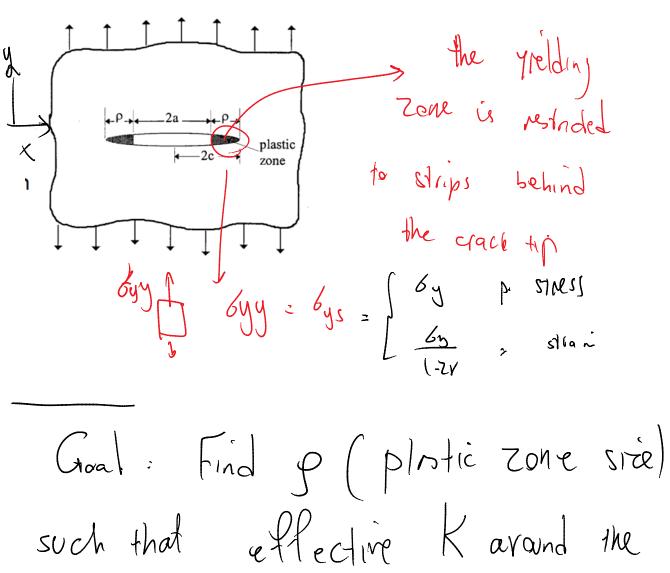
More accurate estimate of r_p:

3. Strip Yield Model

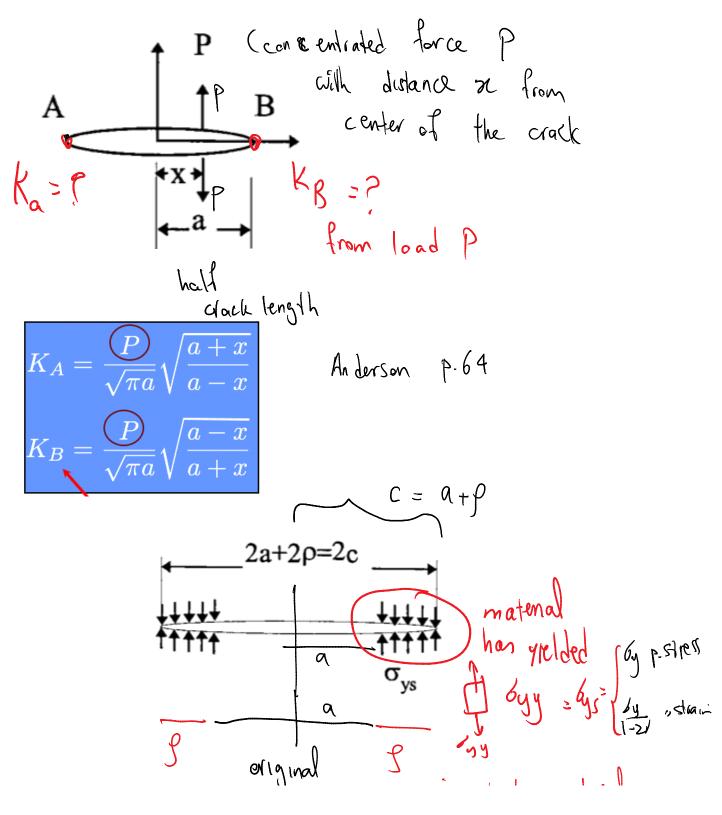


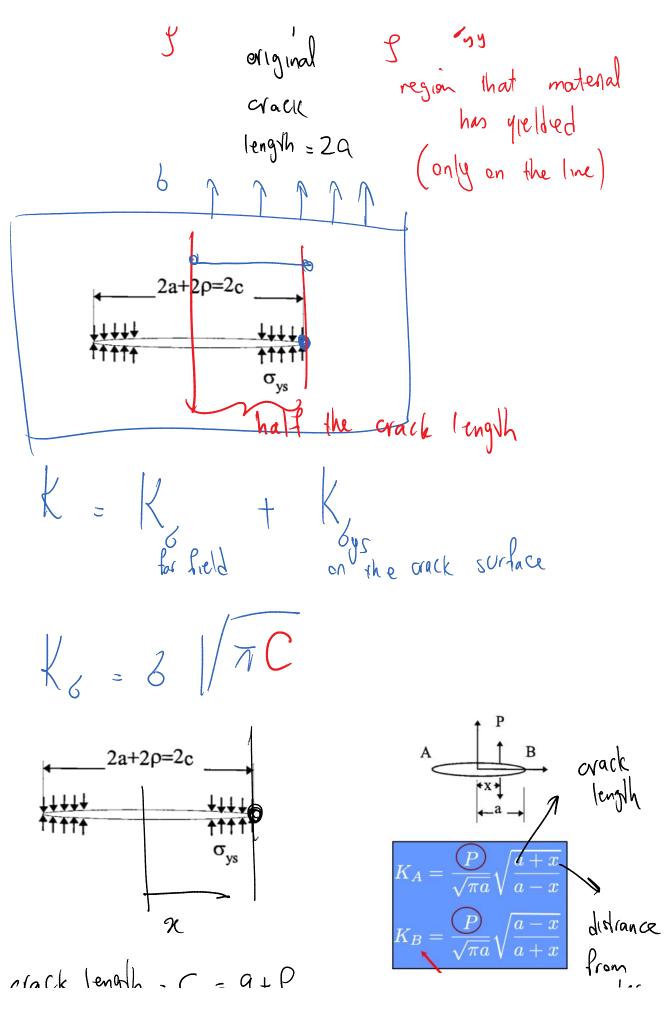




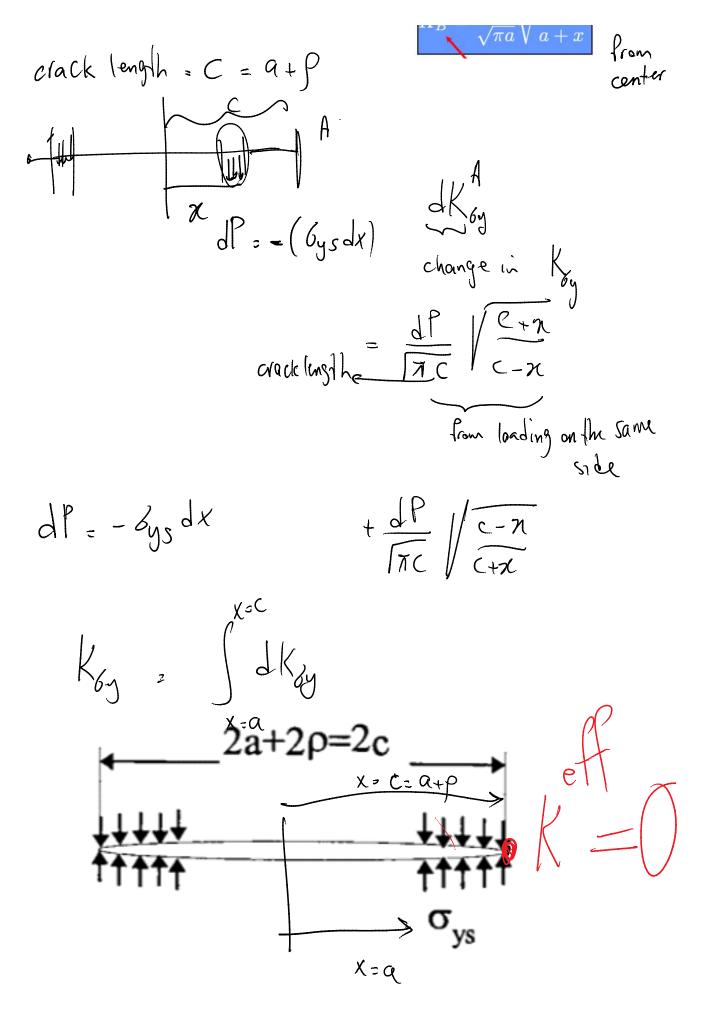


Mark tip is zero



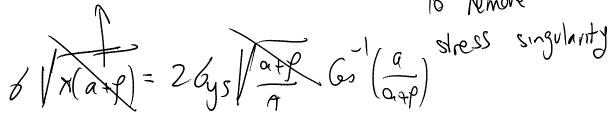


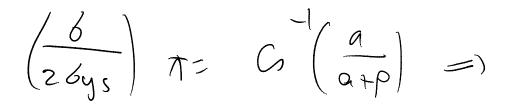
ME524 Page 7



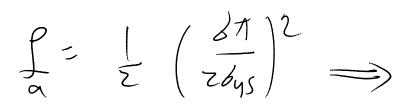
$$K_{I}^{\sigma_{ys}} = -\frac{\sigma_{ys}}{\sqrt{\pi c}} \int_{a}^{c} \left(\sqrt{\frac{c-x}{c+x}} + \sqrt{\frac{c+x}{c-x}} \right) dx$$

$$K_{I}^{\sigma_{ys}} = -2\sigma_{ys} \sqrt{\frac{a+\rho}{\pi}} \cos^{-1} \left(\frac{a}{a+\rho} \right)$$





$$=\left|-\frac{9}{a}\right| - \left|-\frac{1}{2}\left(\frac{6\cdot T}{26ys}\right)^{2}\right| = \right|$$



$$\begin{array}{ccc} f = \frac{1}{2} \left(\frac{1}{2} \frac{1}{2} \frac{1}{3} \right) \end{array}$$

K
For

Gryinal

Gryinal

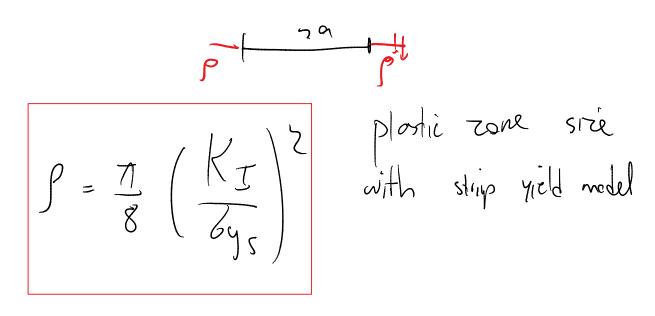
Grack

$$\begin{array}{c} f = \frac{1}{2} \left(\frac{1}{2} \frac{1}{3} \frac{1}{3} \right) \end{array}$$

Gryinal

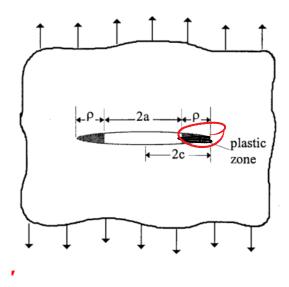
Gry S

$$\begin{array}{c} f = \frac{1}{2} \left(\frac{1}{2} \frac{1}{3} \frac{1}{3} \right) \end{array}$$



$$\boxed{\rho = \frac{\pi^2 \sigma^2 a}{8\sigma_{ys}^2} = \left(\frac{\pi}{8} \left(\frac{K_I}{\sigma_{ys}}\right)^2\right)^2 \text{close to} \\ 0.392 \\ 182 \\ 0.392 \\ 182$$

Strip yield model is exact:



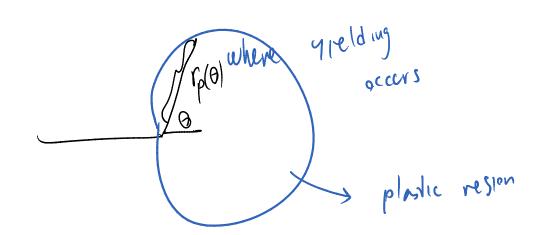
BUT the limitation is that the yield region is restricted to a strip ahead of the crack tip

What is the next level of accuracy? Some model even more realistic than strip yield model.

Next level is allowing the yielding not only ahead of the crack in a strip but in the entire 2D region

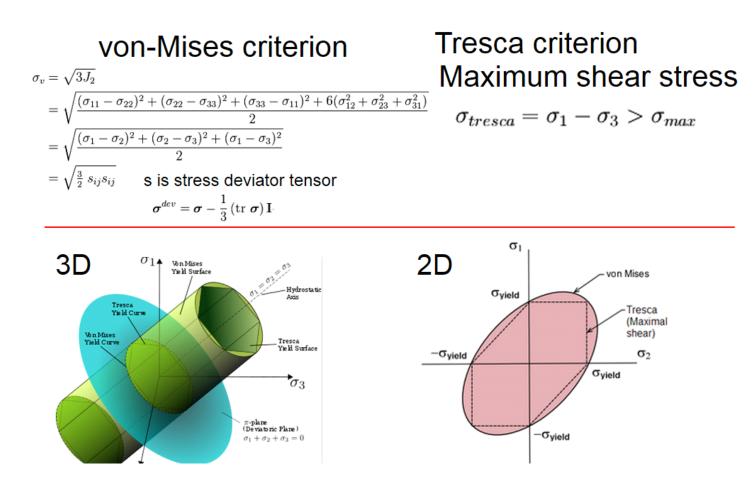
 \sim

strip yield model rp(0) where yrelding occers 0 plastic reston

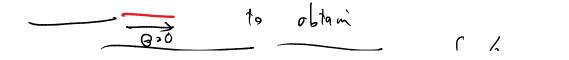


5.2.2 Plastic zone shape: 2D models

- plane stress versus plane strain plastic zones

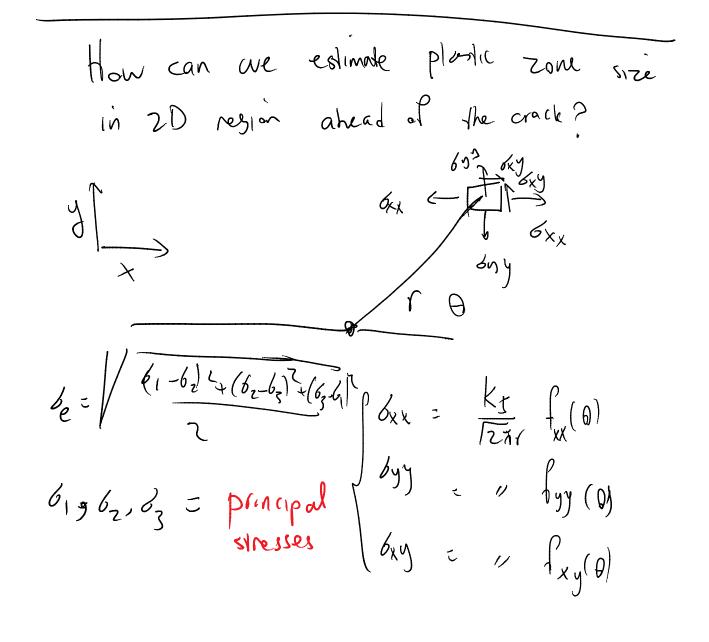


Before we used von Mises yield condition ONLY ahead of the crack

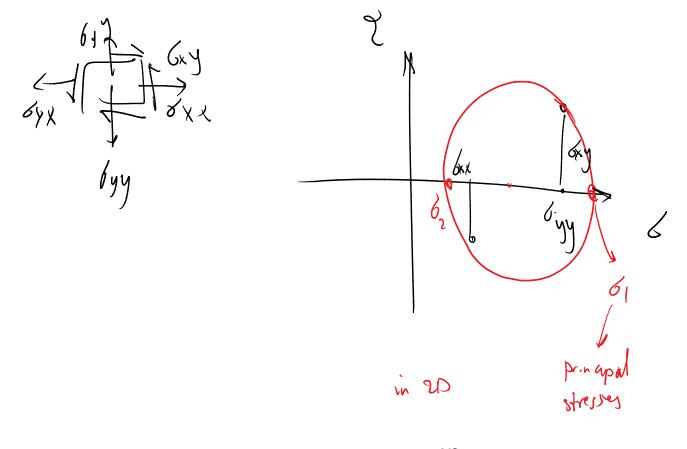


$$b_{q} = \int \frac{(6_{1} - 6_{2})^{2} + (6_{2} - 6_{3})^{2} + (6_{3} - 6_{1})^{2}}{2} = \begin{cases} 6_{yy} & p. \text{ shorts} \\ (1 - 2\lambda) 6_{yy} & p. \text{ shorts} \\ (1 - 2\lambda) 6_{yy} & p. \text{ shorts} \end{cases}$$

$$for \quad \Theta = 0$$



$$\sigma_{xx} = \frac{K_I}{\sqrt{2\pi r^*}} \cos\left(\frac{\theta}{2}\right) \left[1 - \sin\left(\frac{\theta}{2}\right) \sin\left(\frac{3\theta}{2}\right)\right]$$
Slide 124
$$\sigma_{yy} = \frac{K_I}{\sqrt{2\pi r^*}} \cos\left(\frac{\theta}{2}\right) \left[1 + \sin\left(\frac{\theta}{2}\right) \sin\left(\frac{3\theta}{2}\right)\right]$$
$$\tau_{xy} = \frac{K_I}{\sqrt{2\pi r^*}} \cos\left(\frac{\theta}{2}\right) \sin\left(\frac{\theta}{2}\right)$$



$$\sigma_1, \sigma_2 = \frac{\sigma_{xx} + \sigma_{yy}}{2} \pm \left[\left(\frac{\sigma_{xx} - \sigma_{yy}}{2} \right)^2 + \tau_{xy^2} \right]^{1/2}$$

$$b_3 = \begin{cases} 0 & \text{plane stress} \\ V(b_1 + b_2) & \text{plane strain} \\ (c_3 = \frac{1}{E}(b_3 - V(b_1 + b_2)) = 0 \end{cases}$$

$$\sigma_{1} = \frac{K_{I}}{\sqrt{2\pi r}} \cos \frac{\theta}{2} \left(1 + \sin \frac{\theta}{2} \right)$$

$$\sigma_{2} = \frac{K_{I}}{\sqrt{2\pi r}} \cos \frac{\theta}{2} \left(1 - \sin \frac{\theta}{2} \right)$$

$$\sigma_{3} = \begin{cases} 0 & \text{plane stress} \\ \frac{2\nu K_{I}}{\sqrt{2\pi r}} \cos \frac{\theta}{2} & \text{plane strain} \end{cases}$$

$$\begin{array}{c} k_{1X}, b_{XY}, b_{YY} (T, 0) & known \\ Mohr's circle (O p. stress) \\ & S \\ & S \\ & S \\ & S \\ & & \\ &$$

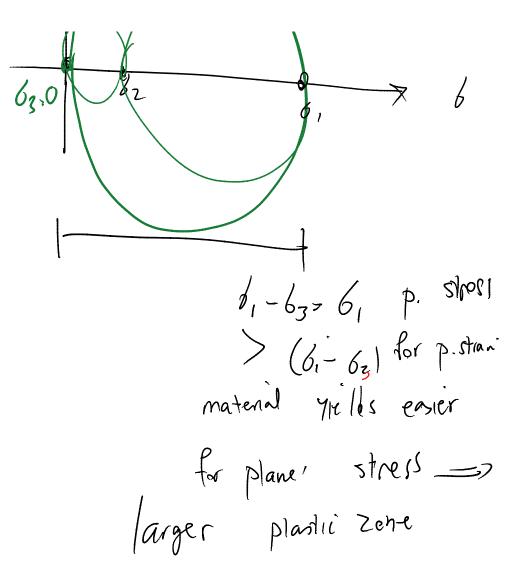
$$\begin{aligned} \mathcal{L}_{e} = \int \frac{1}{(d_1 - d_2)^2 + (d_2 - d_3)^2 + (d_3 - d_3)^2}{2} \\ \hline \mathcal{L}_{e} = \int \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_3 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_3 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_3 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_3 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_3 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_3 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_3 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{1}{(d_1 - d_3)^2 + (d_2 - d_3)^2} \\ \hline \mathcal{L}_{e} = \frac{1}{2} \frac{$$

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$$\frac{\left[r_{y}(\theta) = \frac{1}{4\pi} \left(\frac{K_{I}}{\sigma_{ys}}\right)^{2} \left[1 + \cos\theta + \frac{3}{2}\sin^{2}\theta\right]}{\left[1 - 2\mu\right]^{2} \left(1 - 2\mu\right)^{2} \left(1 + \cos\theta\right) + \frac{3}{2}\sin^{2}\theta} \right] \text{ plane stress}} \quad \text{for down as a } \alpha$$

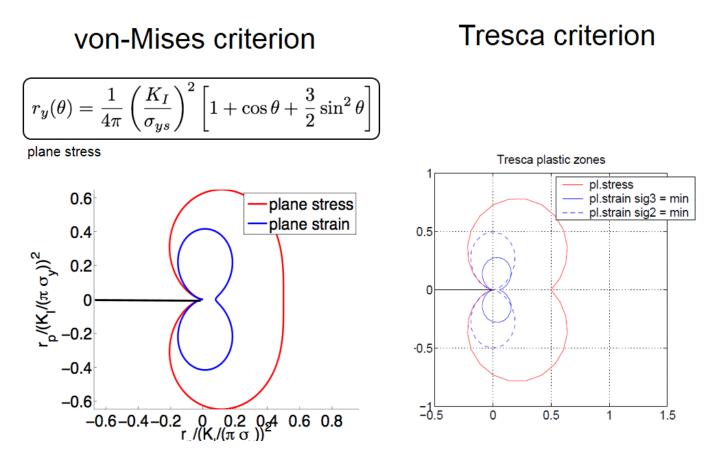
$$r_{y}(\theta) = \frac{1}{4\pi} \left(\frac{K_{I}}{\sigma_{ys}}\right)^{2} \left[(1 - 2\mu)^{2} \left(\frac{1}{4\pi} + \cos\theta\right) + \frac{3}{2}\sin^{2}\theta}\right] \text{ plane stress}} \quad \text{for down at } p$$

$$(abt used from for an interval of the stress of$$



Shapes:

Plastic zone shape



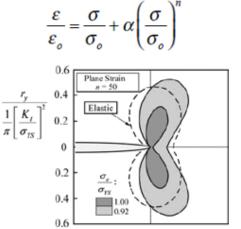
This is a crude approximation (better than 1D versions) but still the stress solution we have does not satisfy balance of linear momentum!

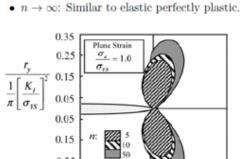
How do we get the correct solution?

We need to solve the problem numerically because it's a difficult problem

Stress redistributed for 2D

Dodds, 1991, FEM solutions Ramberg-Osgood material model





0.25

0.35

n: High strain-hardening.

Effect of definition of yield (some level of ambiguity)

Effect of strain-hardening: Higher hardening (lower n) => smaller zone