

From last time

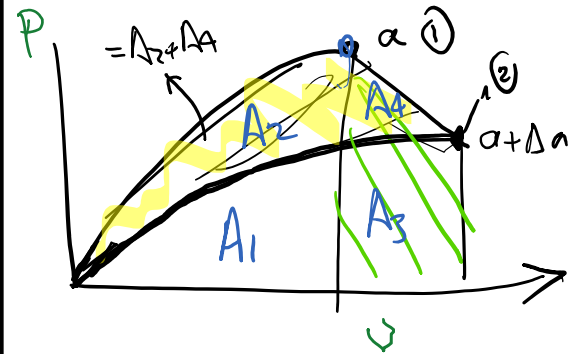


Summary

$$G = \frac{1}{B} \left(-\frac{dU_e}{da} + P \frac{dU}{da} \right)$$

$$= \frac{1}{B} \frac{dU_e}{da} \quad \text{fixed load}$$

$$\approx -\frac{1}{B} \frac{dU_e}{da} \quad \text{grip}$$

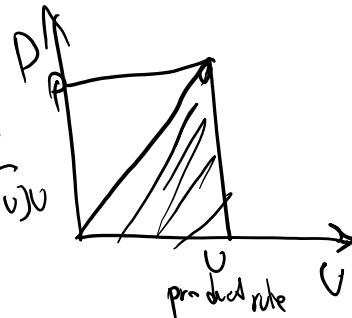


Does (★) match the linear results if the material is linear.

Linear case:

$$U_e = U_e = \frac{1}{2} P U = \frac{1}{2} (K U) U$$

$$U_e = \frac{1}{2} K U^2$$



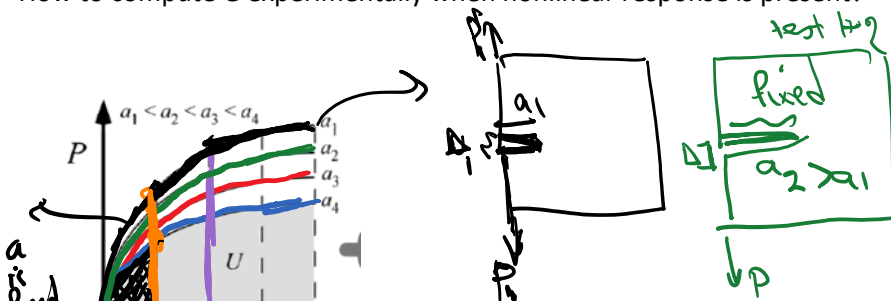
$$G = \frac{1}{B} \left(-\frac{dU_e}{da} + P \frac{dU}{da} \right) = \frac{1}{B} \left(-\frac{d \left(\frac{1}{2} K U^2 \right)}{da} + P \frac{dU}{da} \right)$$

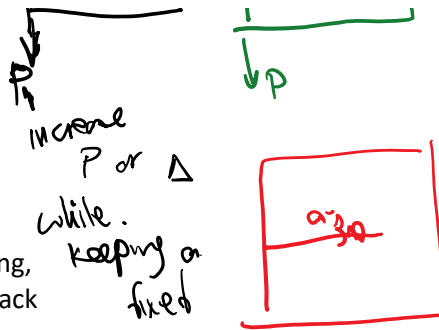
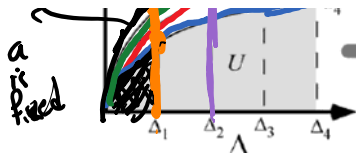
$$= \frac{1}{B} \left(-\frac{1}{2} \frac{dK}{da} \cdot U^2 - \frac{K U \frac{dU}{da}}{P} + P \frac{dU}{da} \right)$$

$$G = -\frac{U^2}{2B} \frac{dK}{da}$$

This matches our result from linear analysis

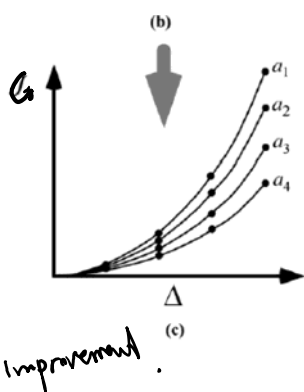
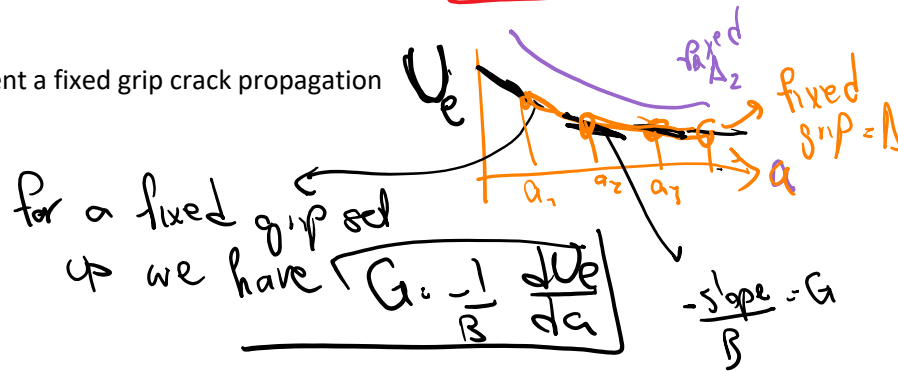
How to compute G experimentally when nonlinear response is present?





In none of these experiments a is increasing, that is none of the experiments involve crack propagation.

Still, we want to represent a fixed grip crack propagation set up.



The issue with this experiment is that we need SEVERAL $P-U$ curves for fixed a 's.

This is cumbersome

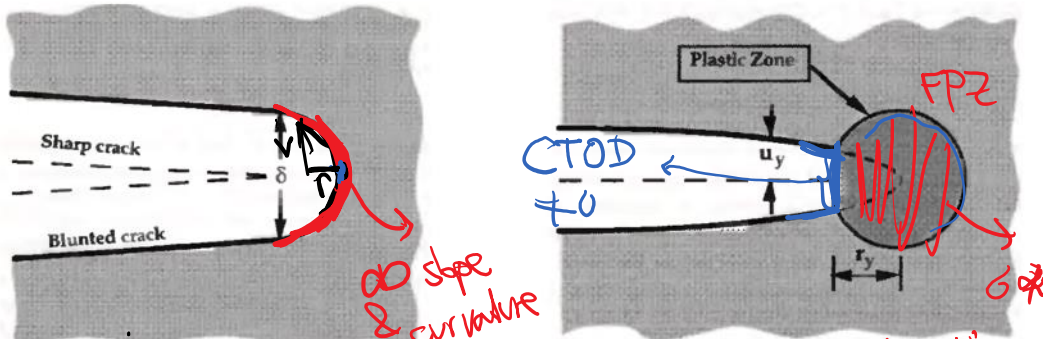
- Rice proposes a method to obtain J with only one test for certain geometries

cf. Anderson 3.2.5 for details

Crack Tip Opening Displacement (CTOD) versus J-integral

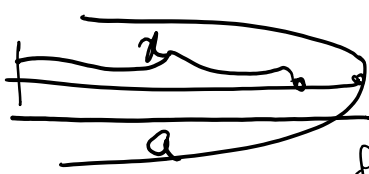
5.4. Crack tip opening displacement (CTOD), relations with J and G

LEFM theory



the asymptotic expansion for v :

$$v = \frac{\kappa + 1}{4\mu} \sigma \sqrt{a^2 - x^2}$$

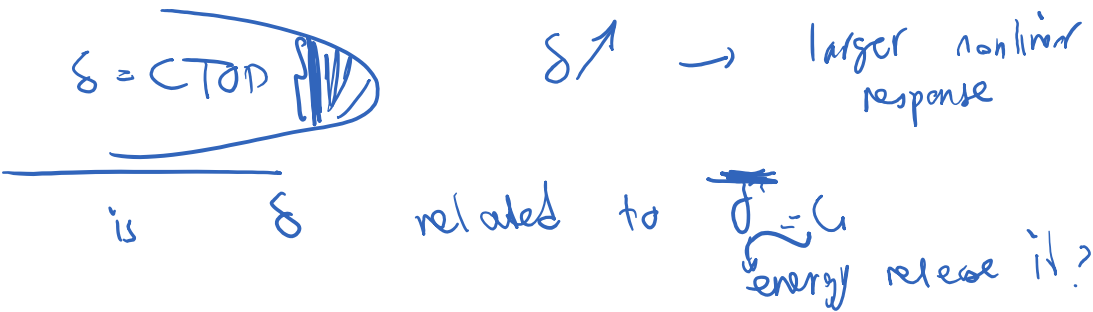


∞ slope & curvature that we have for LCFM $\rightarrow \delta \rightarrow \sigma$ as $\sigma \rightarrow 0$

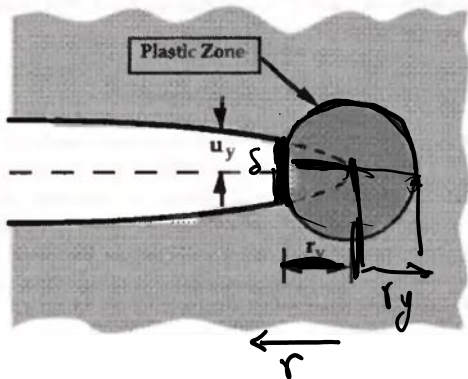
exact displacement (NOT the asymptotic) for a mid-crack specimen

In reality we have yielding & nonlinear response in FPZ

Parallel to Rice's work in the USA, Wells in the UK looked at CTOD as a measure of nonlinear material response in the FPZ



Estimates for CTOD:



without stress redistribution see, if we can get a scale for δ

$$u_y = \frac{(\kappa + 1) K_I}{2\mu} \sqrt{\frac{r}{2\pi}}$$



$$r \rightarrow r_y \rightarrow u_y = \frac{\delta}{2}$$

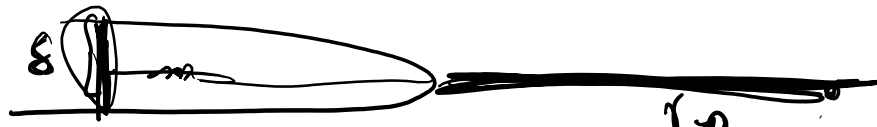
$$\delta/2 = \frac{(\kappa + 1)}{2\mu} \sqrt{\frac{r_y}{2\pi}} \rightarrow$$

$$r_{y,z} = \left(\frac{K_I}{\sigma_y} \right)^2 \rightarrow$$

$$\delta = \frac{K+1}{\sqrt{2\pi}} \frac{K_I}{\mu \sigma_y} \left. \begin{array}{l} K = \frac{3+\nu}{1+\nu} \text{ p. strain} \\ \mu = \frac{E}{2C(1+\nu)} \end{array} \right\} \rightarrow$$

$$\delta \approx \frac{4}{\sqrt{\pi}} \frac{K_I^2}{\sigma_y E}$$

estimate for CTOD



$\delta \propto \frac{K_I^2}{E \sigma_y}$

FPZ size

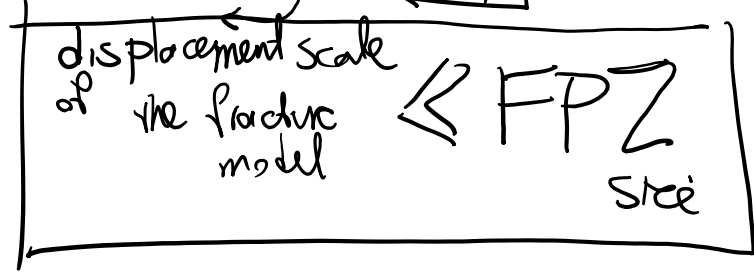
$r_p = \frac{K_I^2}{\sigma_y^2}$

$$\frac{\delta}{r_p} = \frac{\sigma_y}{E}$$

$\rightarrow 0.01 \sim 0.01$

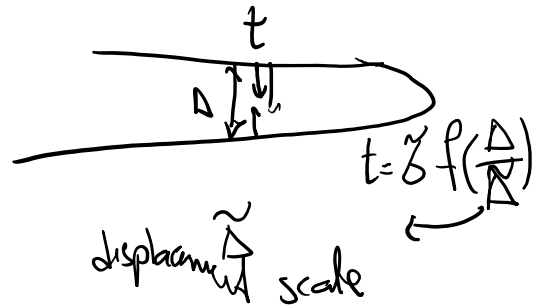
$$\delta \ll r_p$$

for LEFM $r_s = \frac{K_I^2}{\sigma^2}$
for LEFM $\delta \propto \sigma_y$



In many instances, LEFM, PFM, Traction-separation laws, frictional laws, ... FPZ size is much larger than the displacement scale of the model and typically the ratio is proportional to E/σ_y

Cohesive models



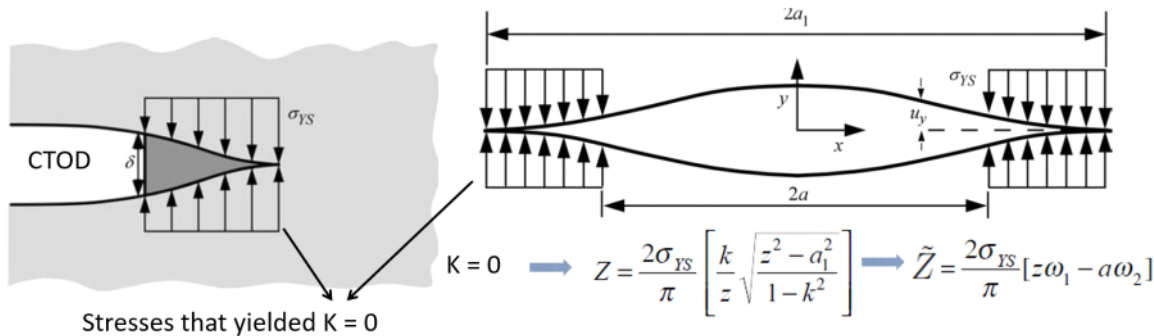
$$FPZ \propto \sqrt{\frac{E}{\sigma_y}}$$

$$FPZ \propto \tilde{\Delta} \left(\frac{E}{\sigma} \right)$$

displacement $\tilde{\Delta}$ scale \leftarrow

The above estimate for CTOD was very crude and did not consider the actual solution of the problem with nonlinear response.

Crack Tip Opening Displacement: Strip yield model



$$\rightarrow u_y = \frac{2}{E} \text{Im} \tilde{Z} = \frac{4\sigma_{YS}}{\pi E} \left[a \coth^{-1} \left(\frac{1}{a_1} \sqrt{\frac{a_1^2 - z^2}{1 - k^2}} \right) - z \coth^{-1} \left(\frac{k}{z} \sqrt{\frac{a_1^2 - z^2}{1 - k^2}} \right) \right]$$

$$\rightarrow u_y = \frac{2}{E} \text{Im} \tilde{Z} = \frac{4\sigma_{YS}}{\pi E} \left[a \coth^{-1} \left(\frac{1}{a_1} \sqrt{\frac{a_1^2 - z^2}{1 - k^2}} \right) - z \coth^{-1} \left(\frac{k}{z} \sqrt{\frac{a_1^2 - z^2}{1 - k^2}} \right) \right]$$

$$z = a \rightarrow \delta = 2u_y = \frac{8\sigma_{YS}a}{\pi E} \ln \left(\frac{1}{k} \right) = \frac{8\sigma_{YS}a}{\pi E} \left[\frac{1}{2} \left(\frac{\pi \sigma}{2 \sigma_{YS}} \right)^2 + \frac{1}{12} \left(\frac{\pi \sigma}{2 \sigma_{YS}} \right)^4 + \dots \right] \rightarrow$$

$$\text{For } \sigma/\sigma_{YS} \rightarrow 0 \quad \delta = \frac{K_I^2}{\sigma_{YS} E}$$

this matches our crude earlier estimate
 $\delta \approx \left(\frac{4}{\pi} \right) \frac{K_I^2}{\sigma_{YS} E}$

Is there a relation between J and CTOD?

Both can measure the extent of nonlinear response, but are they related?

$$\left. \begin{aligned} \delta &= \frac{K_I^2}{E \sigma_{YS}} \\ J = G &= \frac{K_I^2}{F \approx E'} \end{aligned} \right\} \rightarrow \delta = \frac{J}{\sigma_y} \rightarrow$$

$$\delta = \frac{u_d}{E \approx E'}$$

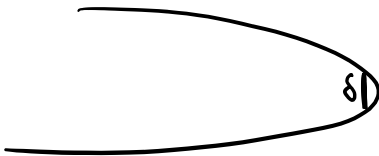
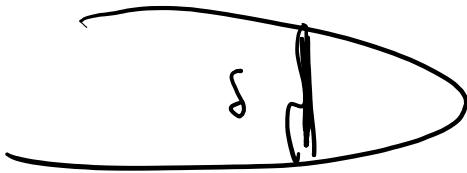
they are related

$$J = \delta \sigma_y$$

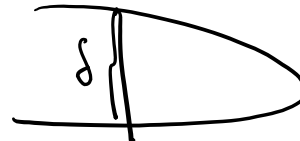
$$\delta = \frac{J}{\sigma_y}$$

$J \uparrow$ (higher loads) \rightarrow higher energy release rate

OR if $\sigma_y \downarrow$ (more ductility) $\rightarrow \delta \uparrow$



LEFM good
 σ_y high quasi-brittle material



ductile
 σ_y low

CTOD-J relation

- When SSY is satisfied $G = J$ so we expect:

$$G = m \sigma_y \delta \Rightarrow J = m \sigma_y \delta$$

- In fact this equation is valid well beyond validity of LEFM and SSY

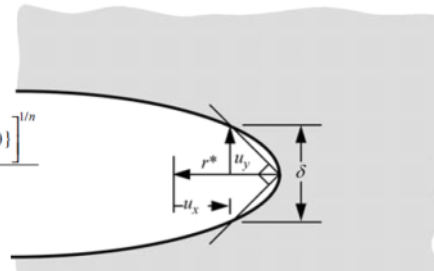
- E.g. for HRR solution Shih showed that:

$$u_i = \frac{\alpha \sigma_o}{E} \left(\frac{EJ}{\alpha \sigma_o^2 I_n r} \right)^{\frac{n}{n+1}} r \tilde{u}_i(\theta, n)$$

$$d_n = \frac{2 \tilde{u}_y(\pi, n) \left[\frac{\alpha \sigma_o}{E} \{ \tilde{u}_x(\pi, n) + \tilde{u}_y(\pi, n) \} \right]^{1/n}}{I_n}$$

- δ is obtained by 90 degree method:
 Deformed position corresponding to $r^* = r$ and $\varphi = -\pi$ forms 45 degree w.r.t crack tip)

$$\frac{\delta}{2} = u_y(r^*, \pi) = r^* - u_x(r^*, \pi)$$

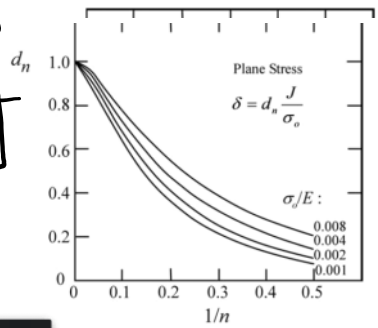


$\varphi = -\pi$ (forms 45 degree w.r.t crack tip)

$$\frac{\delta}{2} = u_y(r^*, \pi) = r^* - u_x(r^*, \pi)$$

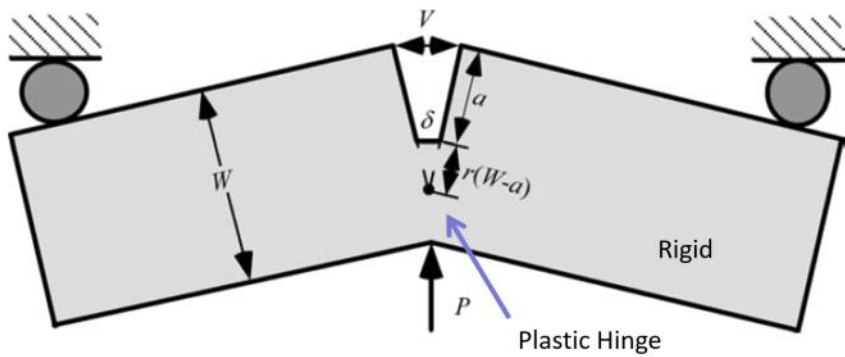
$$r^* = \left(\frac{\alpha \sigma_o}{E} \right)^{1/n} \{ \tilde{u}_x(\pi, n) + \tilde{u}_y(\pi, n) \}^{\frac{n+1}{n}} \frac{J}{\sigma_o I_n} \Rightarrow J = m \sigma_o \delta$$

for $m = \frac{1}{d_n}, d_n = \frac{2 \tilde{u}_y(\pi, n) \left[\frac{\alpha \sigma_o}{E} \{ \tilde{u}_x(\pi, n) + \tilde{u}_y(\pi, n) \} \right]^{1/n}}{I_n}$



Experimentally, one can measure CTOD and from there from the relation above compute J (this is another way to experimentally measure J)

For the description of the experiment, please refer to Anderson:



$$\frac{\delta_t}{CMOD} = \frac{r(W-a)}{r(W-a)+a} \quad \text{similarity of triangles}$$

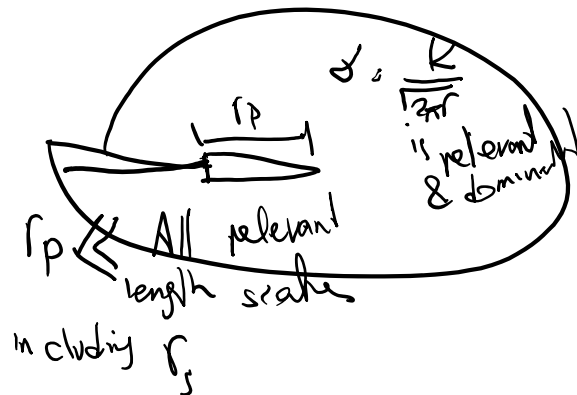
r : rotational factor [-], between 0 and 1

For high elastic deformation contribution, elastic corrections should be added

When can we use LFEM, PFM, ...?

SSY is satisfied \rightarrow LEFM

necessary condition (but not sufficient)



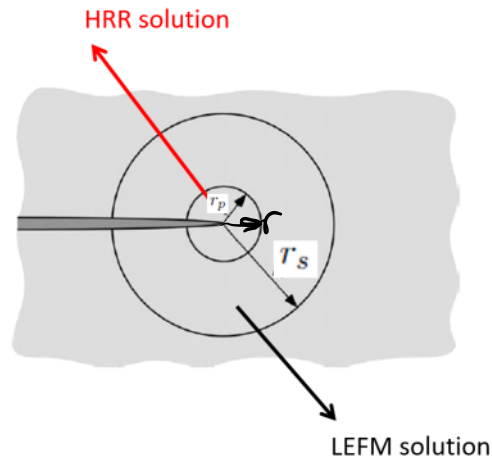
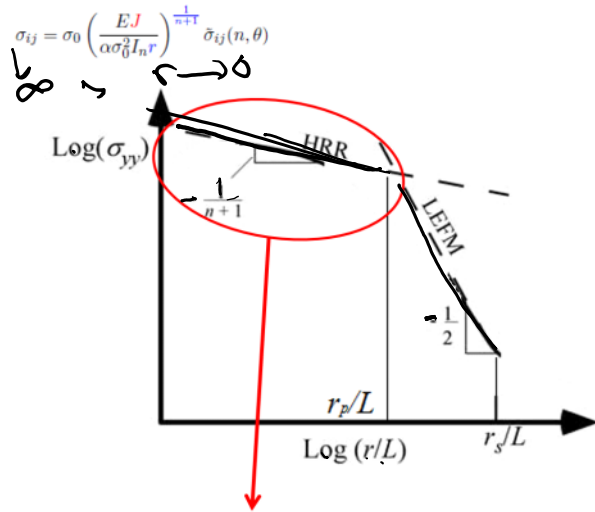
$$\frac{r_p}{\left(\frac{K_I}{\sigma_y} \right)^2} \ll \frac{r_s}{\left(\frac{K_I}{\sigma} \right)^2} \rightarrow \left(\frac{\sigma}{\sigma_y} \right)^2 \ll 1 \quad \text{necessary condition}$$

if SSY is violated we need to use PFM

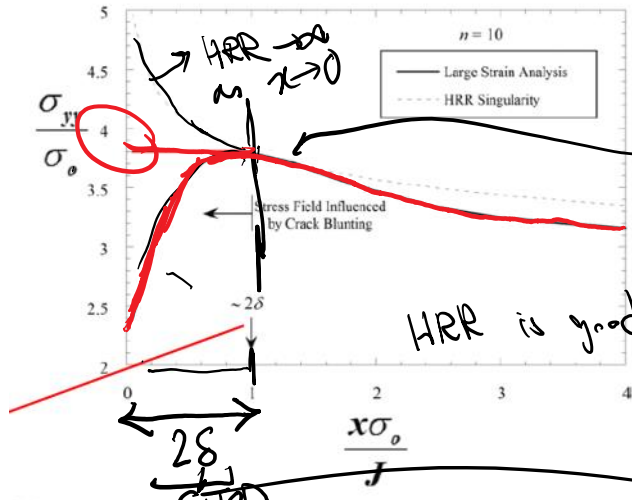
$G = J$ still valid (as long as there is not much plastic unloading)
 $\sigma = \sigma^{\alpha} r^{\beta}$ stress solution
 HRR - solution

So, when PFM HRR solution is no longer relevant?

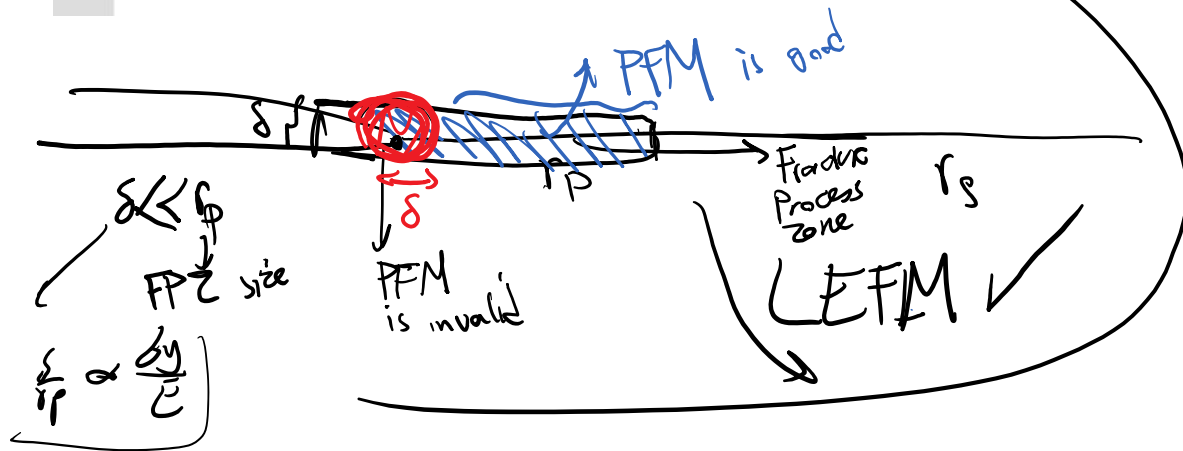
HRR solution: Stress singularity



Stress is still singular but with a weaker power of singularity!



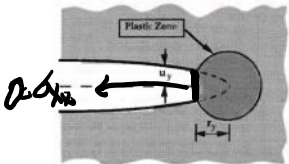
McMeeking and Parks, ASTM STP 668, ASTM 1979



Limitations of HRR analysis

- Small strain: $\epsilon = \frac{1}{2}(\nabla u + \nabla^T u)$ (accurate for $\epsilon \lesssim 0.1$)
- Small deformation theory (e.g., not using PK stresses, etc)
- Elastic HRR model instead of plastic model
- Crack tip blunting: $\Rightarrow \sigma_{xx} = 0$

otherwise need to use large-strain definition
 $G = \frac{1}{2}(C^T - I)$
 $C = \nabla u + I$



From SSY to LSY

Large strain radius $r_n \propto \delta$ (CTOD): $\delta = O\left(\frac{K^2}{E\sigma_y}\right)$

plastic radius: $r_p = O\left(\frac{K^2}{\sigma_y^2}\right)$

K-dominant radius: $r_s = O\left(\frac{K^2}{\sigma^2}\right)$

$\bar{\sigma}$: applied stress

SSY (Small Scale Yielding)

$r_n \ll r_p \ll r_s \Rightarrow$

$\frac{r_p}{r_s} \propto \left(\frac{\bar{\sigma}}{\sigma_y}\right)^2 \ll 1$

Elastic plastic condition

$r_n \ll r_p \approx r_s \Rightarrow$

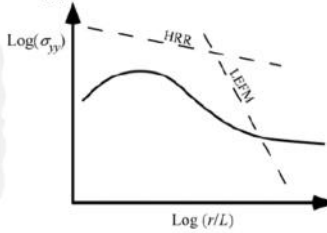
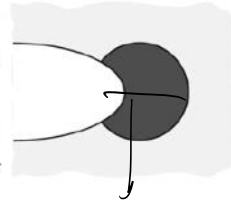
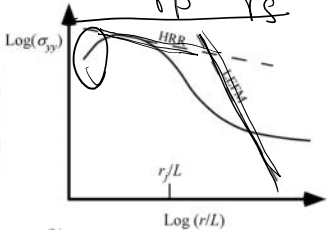
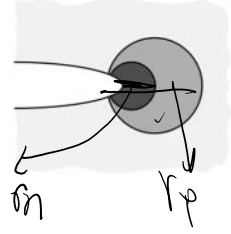
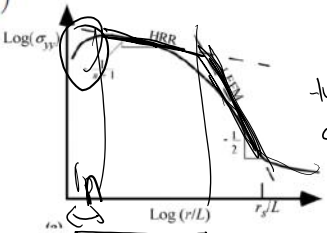
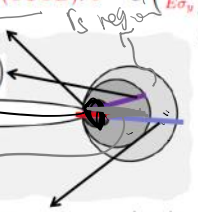
$\frac{r_n}{r_p} \ll 1, \frac{r_p}{r_s} \propto \left(\frac{\bar{\sigma}}{\sigma_y}\right)^2 \approx 1$

LSY (Large Scale Yielding)

$r_n \approx r_p$

Note that $\frac{\delta}{r_p} \propto \left(\frac{\sigma_y}{E}\right)$

- Large Strain Region
- J-Dominated Zone
- K-Dominated Zone
- No Single-Parameter Characterization



r_n takes over Fracture process zone

LEFM: SSY satisfied and generally have

$\bar{\sigma} \ll \sigma_y$

Relevant parameters: G (energy) K (stress)

PFM (or NFM): SSY is gradually violated and

$\bar{\sigma} \approx \sigma_y$

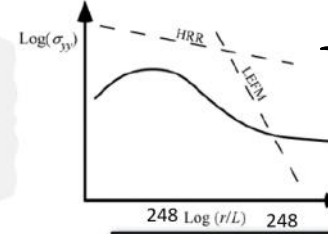
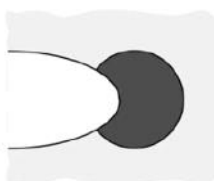
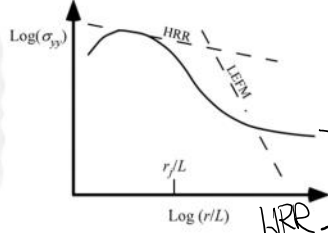
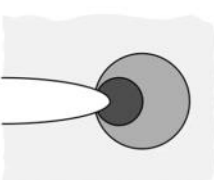
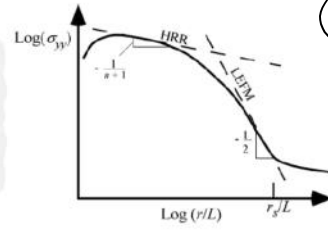
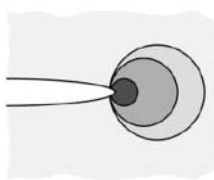
Relevant parameters: J (energy & used for stress)

LSY condition:

No single parameter can characterize fracture!

J + other parameters (e.g. T stress, Q-J, etc)

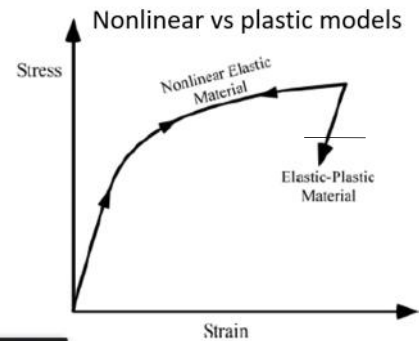
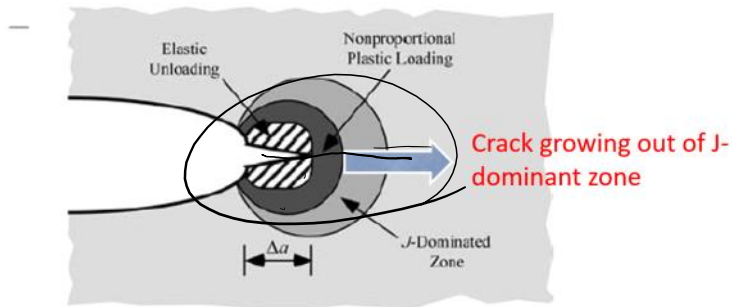
- Large Strain Region
- J-Dominated Zone
- K-Dominated Zone
- No Single-Parameter Characterization



LSY: When a single parameter (G, K, J, CTOD) is not enough?

- Under considerable plastic deformation and crack propagation when unloading and non-proportional zones grow out of J dominant zone with crack propagation. Reasons are:

– Unloading In J integral analysis plastic model was replaced by a



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