

Why we need to exactly satisfy essential BC when using energy methods

$\Pi(y_e) \leq \Pi(y)$

y is a trial solution

$\Pi(y) = \int_0^L \frac{1}{2} EI y''^2 dx - \int y q dx$

internal energy external work

trial function y_1

y_e

from here a trial solution should have 2 derivatives

Example $q=0$

$\Pi(y_1) = \int_0^L \frac{1}{2} EI y_1''^2 dx - \int y_1 q dx$

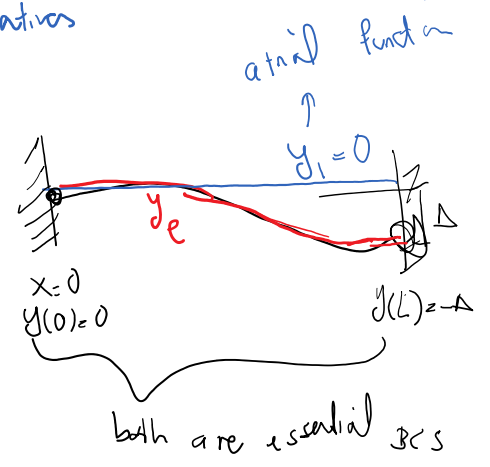
$= 0$

Exact solution

$\Pi(y_e) = \int_0^L \frac{1}{2} EI y_e''^2 dx > 0$

$\Pi(y_e) > \Pi(y_1)$

exact solution



whereas we should have had $\Pi(y_e) \leq \Pi(y) \forall y$

For energy method, we need to exactly satisfy all essential BCs. Unlike different forms of WRS, here we have not choice. We must satisfy them for trial functions.

-> "Essential"

Let's look at energy statement again:

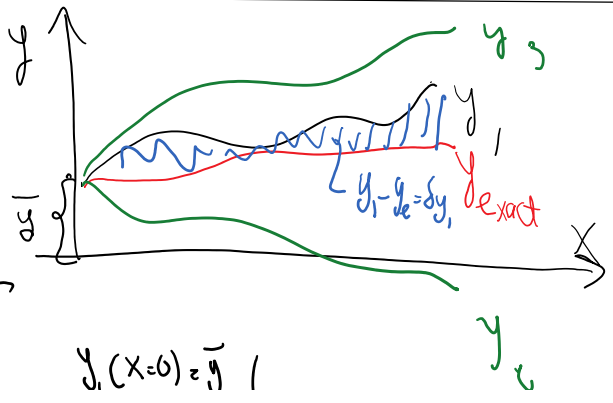
essential BC $y(0) = \bar{y}$

Find y_e such that $\Pi(y_e) = \{ f(x), f(0) = \bar{y} \}$

$\Pi(y_e) \leq \Pi(y)$ (with a star in a circle)

where

$r_1 < r_2$ $r_1 < r_2$



where

$$\Pi(y) = \int_0^L \frac{1}{2} EI y''^2 dx - \int_0^L q y dx$$

another way to write (*) is by looking at the increment of y_e

Find y_e such that for all

increments $\delta y \in W = \{f \in C^1(\Omega) \mid f(0) = 0\}$

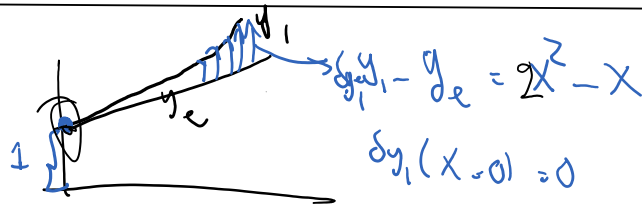
$$\Pi(y_e) \leq \Pi(y_e + \delta y)$$

$$\left. \begin{aligned} y_1(x=0) &= \bar{y} \\ y_2(x=0) &= \bar{y} \end{aligned} \right\} \rightarrow$$

$$\delta y_1(0) = (y_1 - y_2)(0) = \bar{y} - \bar{y} = 0$$

So the increment of a trial solution w.r.t. the exact solution is zero at all essential BCs (similar to the weight function condition in weak statement)

$$\begin{aligned} \bar{y} &= 1 \\ y_e &= 1 + x \\ y_1 &= 1 + 2x^2 \end{aligned}$$



Energy Method for Solid Mechanics

The total energy in solid mechanics is,

$$\Pi = (V - W) - T = \text{Total energy} \quad (85a)$$

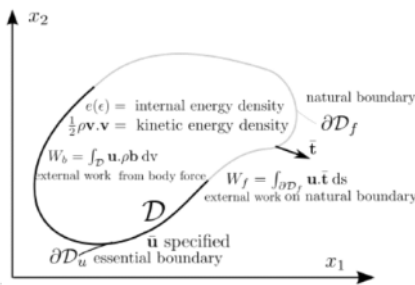
$$T = \int_D \frac{1}{2} \rho v \cdot v \, dv = \text{Kinetic energy} \quad (85b)$$

$$V = \int_D e(\epsilon) \, dv = \text{Internal energy} \quad (85c)$$

$$W = W_b + W_f = \text{External work} \quad (85d)$$

$$W_b = \int_D \mathbf{u} \cdot \rho \mathbf{b} \, dv \quad (85e)$$

$$W_f = \int_{\partial D_f} \mathbf{u} \cdot \bar{\mathbf{t}} \, ds \quad (85f)$$



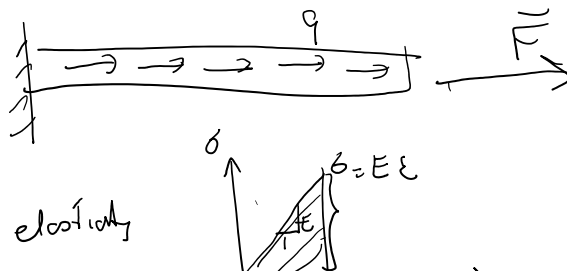
- For static problems $T = 0$.
- Internal energy density, $e(\epsilon) = \frac{1}{2} \epsilon : \sigma(\epsilon) = \frac{1}{2} C_{ijkl} \epsilon_{ij} \epsilon_{kl}$ for linear solid.
- Natural boundary forces are naturally incorporated into the energy (W_f).
- Essential boundary conditions are incorporated into function space:

$$\mathbf{u} \in \mathcal{V} = \{ \mathbf{v} \mid \mathbf{v} \in C^1(D) : \forall \mathbf{x} \in \partial D_u \, \mathbf{v}(\mathbf{x}) = \bar{\mathbf{u}}(\mathbf{x}) \}, \text{ is a solution if } \forall \bar{\mathbf{u}} \in \mathcal{V}, \Pi(\mathbf{u}) \leq \Pi(\bar{\mathbf{u}}). \quad (86)$$

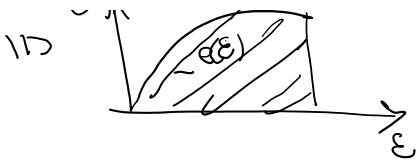
Derive the energy statement for a bar problem:

$$V = \int_V e(\epsilon) \, dv$$

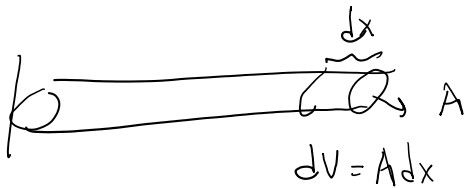
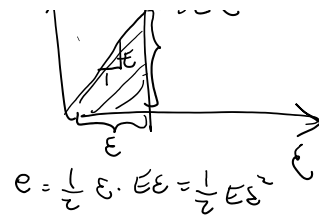
internal energy per unit volume



1D linear elasticity



1D linear elasticity



$e(\epsilon) = \frac{1}{2} E \epsilon^2$, $\epsilon = \frac{du}{dx}$
 $\int_{\text{volume}} e(\epsilon) dV = \int_0^L \frac{1}{2} E u'{}^2 \cdot A dx = \int_0^L \frac{1}{2} EA u'{}^2 dx$ (1)

External work

$W = W_b + W_f$
 ↓
 work

(2)



$W_b = \int_0^L u \cdot dF_r = \int_0^L u (q dx) = \int_0^L u q dx$ (3)

$W_f = u \bar{F}|_{x=L}$ (4)

$\Pi(u) = \int_0^L \frac{1}{2} EA u'{}^2 dx - \int_0^L q u dx - u \bar{F}|_{x=L}$

Let's find the exact solution for this problem

Assume u is the exact solution

$\Pi(u) \leq \Pi(u + \delta u)$ $\forall \delta u \in \mathcal{W} = \{ f \in C^1(\Omega, L) \mid f(0) = 0 \}$
 \Downarrow
 $\Pi(u + \delta u) - \Pi(u) \geq 0$
 we must have this condition

$\Pi(u+\delta u) - \Pi(u) \geq 0$ we must have this condition

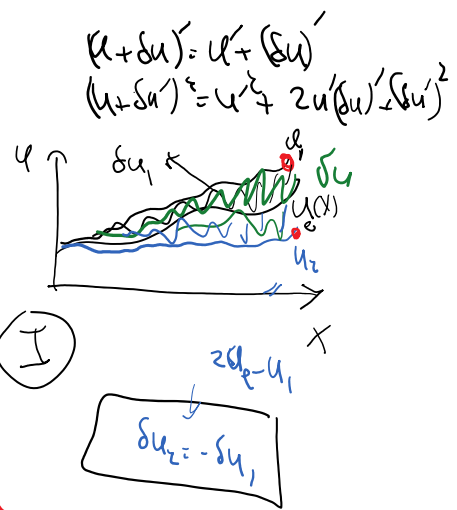
$$\Pi(u+\delta u) - \Pi(u) = \int_0^L \frac{1}{2} EA [(u+\delta u)']^2 dx - \int_0^L (u+\delta u) q dx - (u+\delta u) \bar{F} \Big|_{x=L}$$

$$= \int_0^L \frac{1}{2} EA u'^2 dx - \int_0^L u q dx - u \bar{F} \Big|_{x=L} + \int_0^L EA (2u' \delta u') dx + \int_0^L \frac{1}{2} EA (\delta u')^2 dx - \int_0^L q \delta u dx - \int_0^L \delta u q dx - u \bar{F} \Big|_{x=L} - \delta u \bar{F} \Big|_{x=L}$$

$\int_0^L \frac{1}{2} EA u'^2 dx$ (crossed out)
 $-\int_0^L \frac{1}{2} EA u'^2 dx$ (crossed out)
 $\int_0^L EA (2u' \delta u') dx$ (underlined)
 $\int_0^L \frac{1}{2} EA (\delta u')^2 dx$ (underlined)
 $-\int_0^L q \delta u dx$ (crossed out)
 $-\int_0^L \delta u q dx$ (crossed out)
 $-u \bar{F} \Big|_{x=L}$ (crossed out)
 $-\delta u \bar{F} \Big|_{x=L}$ (crossed out)

Note this for \otimes

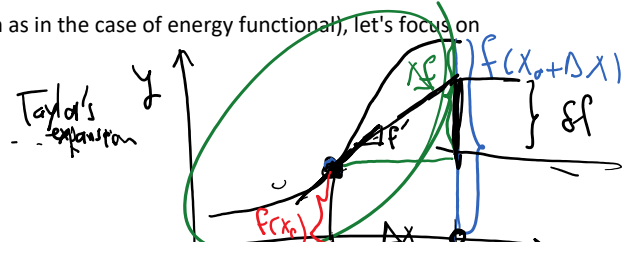
$\Delta \Pi = \Pi(u+\delta u) - \Pi(u) =$
 $\int_0^L EA u' \delta u' dx - \int_0^L q \delta u dx - \delta u \bar{F} \Big|_{x=L}$ (1st order in δu)
 $+ \int_0^L \frac{1}{2} EA (\delta u')^2 dx \geq 0$ (2nd order in increments)
 $\delta^2 \Pi$ (circled)
 if $\delta^2 \Pi > 0$ we have a local minimum



Let's discuss what we can conclude from the inequality in equation (I)

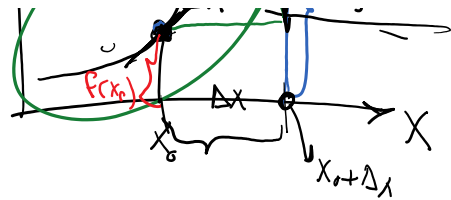
Instead of functionals (a scalar function whose argument is a function as in the case of energy functional), let's focus on functions from real numbers to real numbers

Δf total variation = $f(x_0 + \Delta x) - f(x_0)$
 $f(x_0) = f_0, f'(x_0) = f'_0, f''(x_0) = f''_0$



total variation

$$f(x_0) + \Delta x f'(x_0) + \frac{1}{2} \Delta x^2 f''(x_0) + \dots - f(x_0)$$



$$\Delta f = \underbrace{f'(x_0) \Delta x}_{\text{first variation } \delta f} + \underbrace{\frac{1}{2} f''(x_0) \Delta x^2}_{\delta^2 f} + \frac{1}{6} f'''(x_0) \Delta x^3 + \dots$$

if $\Delta x \rightarrow 0$ we can ignore $\delta^2 f, \delta^3 f, \dots$ relative to δf

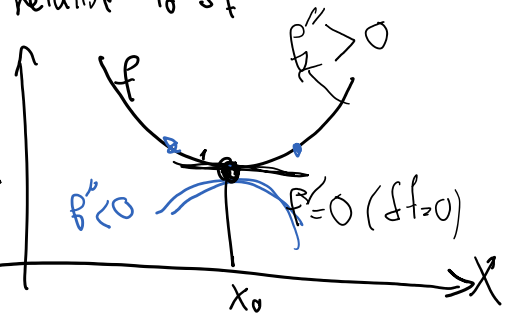
if f has a minimum @ x_0 what can we say about $\delta f, \delta^2 f$

$$\begin{matrix} \delta f \\ = \\ 0 \end{matrix}$$

Necessary condition for min/max

> 0 local min
 < 0 local max

$= 0$... need to go to H.O.T



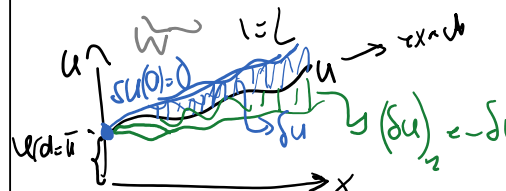
if we have $\delta f = 0$ & $\delta^2 f > 0$ we have a local minimum

$$\Delta \Pi = \Pi(u + \delta u) - \Pi(u) = \int_0^L EA u' \delta u dx - \int_0^L q \delta u dx - \delta u | F |_{x=L} + \int_0^L \frac{1}{2} EA [\delta u']^2 dx \geq 0$$

1st order in δu

2nd order in increments

$\delta^2 \Pi$ if $\delta^2 \Pi > 0$ we have



$\delta^2 \Pi$ if $\delta^2 \Pi > 0$ we have a local

For 0 to be the exact soln, we should have

$$\Pi(u) \leq \Pi(u + \delta u) \rightarrow \delta \Pi = \Pi(u + \delta u) - \Pi(u) \geq 0$$

$$\delta \Pi = \int_0^L (\delta u)' EA u' dx - \int_0^L \delta u \cdot q dx - \delta u|_{x=L} \bar{F} = 0$$

$$\delta^2 \Pi = \int_0^L \frac{1}{2} EA (\delta u')^2 dx > 0$$

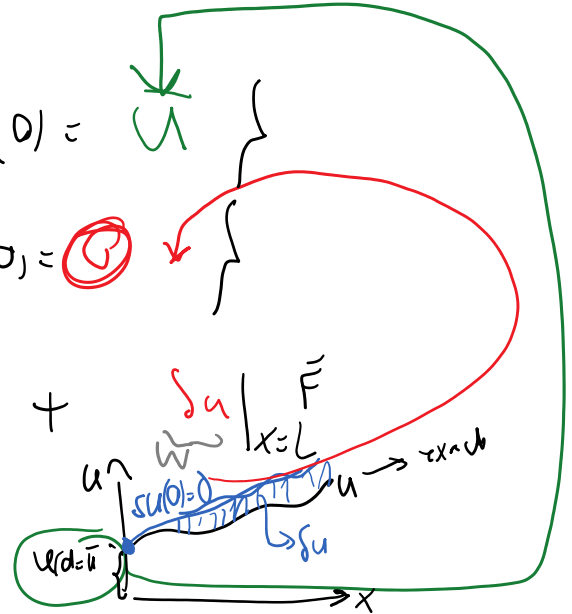
$\delta u \neq 0$

for exact soln all we need to do is to satisfy

$$\delta \Pi = 0$$

Find $u \in \mathcal{V} = \{ f \in C^1 \mid f(0) = \bar{u} \}$
 $\exists \forall \delta u \in \mathcal{W} = \{ f \in C^1 \mid f(0) = 0 \}$

$$\delta \Pi = 0 \rightarrow \int_0^L (\delta u)' EA u' dx = \int_0^L \delta u q dx + \delta u|_{x=L} \bar{F}$$



Weak Statement ☺